

Optimal friction for Langevin sampling

Martin Chak

joint work w/ Nikolas Kantas, Tony Lelièvre, Grigorios Pavliotis

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Given

- a smooth potential function $V : \mathbb{R}^n \rightarrow \mathbb{R}$ and its gradient $\nabla V : \mathbb{R}^n \rightarrow \mathbb{R}^n$, which can be evaluated at any point in \mathbb{R}^n at a cost, and
- an everywhere differentiable observable function $L^2(e^{-V}) \ni f : \mathbb{R}^n \rightarrow \mathbb{R}$,

approximate the integral

$$\pi(f) := \int_{\mathbb{R}^n} f \left(\frac{e^{-V}}{\int_{\mathbb{R}^n} e^{-V}} \right).$$

MCMC Solution (for large n): Use stochastic dynamics, e.g. (discretizations of)

$$\begin{aligned} \text{overdamped Langevin:} \quad & dq_t = -\nabla V(q_t)dt + \sqrt{2}dW_t, \\ \text{or (underdamped) Langevin:} \quad & dq_t = p_t dt, \\ & dp_t = -\nabla V(q_t)dt - \Gamma p_t dt + \sqrt{2\Gamma}dW_t \end{aligned}$$

for any positive symmetric definite $\Gamma \in \mathbb{R}^{n \times n}$ and set $\pi_T(f) := \frac{1}{T} \int_0^T f(q_t)dt$.
Overdamped Langevin dynamics has the invariant measure $\pi(dq) \propto e^{-V(q)}dq$ and
underdamped Langevin has the invariant measure $\tilde{\pi}(dqdp) \propto e^{-V(q) - \frac{p^2}{2}} dqdp$.

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Underdamped Langevin $dq_t = p_t dt$, $dp_t = -\nabla V(q_t)dt - \Gamma p_t dt + \sqrt{2\Gamma}dW_t$

In order to obtain some measure of ‘speed’ at which

$\pi_T(f) := \frac{1}{T} \int_0^T f(q_t)dt \rightarrow \int_{\mathbb{R}^n} f \left(\frac{e^{-V}}{\int_{\mathbb{R}^n} e^{-V}} \right) = \pi(f)$ occurs as $T \rightarrow \infty$, use the fact that $\frac{1}{\sqrt{T}} \int_0^T (f(q_t) - \pi(f))dt \rightarrow \mathcal{N}(0, \sigma^2)$ in distribution as $T \rightarrow \infty$ for some $\sigma^2 > 0$:

Claim: $\sigma^2 = 2 \int_{\mathbb{R}^n} \nabla_p \phi^\top \Gamma \nabla_p \phi d\tilde{\pi}$, where ϕ solves

$$-\mathcal{L}\phi := (-p^\top \nabla_q + \nabla V(q)^\top \nabla_p + \Gamma p^\top \nabla_p - \nabla_p^\top \Gamma \nabla_p)\phi = f - \pi(f).$$

‘Proof’: Suppose there exists a well-behaved solution $\phi : \mathbb{R}^{2n} \rightarrow \mathbb{R}$ and use Itô’s lemma to obtain

$$\begin{aligned} \phi(q_t, p_t) &= \phi(q_0, p_0) + \int_0^t \mathcal{L}\phi(q_s, p_s)ds + \int_0^t \nabla_p \phi(q_s, p_s)^\top \sqrt{2\Gamma}dW_s \\ &= \phi(q_0, p_0) - \int_0^t (f(q_s) - \pi(f))ds + \int_0^t \nabla_p \phi(q_s, p_s)^\top \sqrt{2\Gamma}dW_s \end{aligned}$$

so that, using Itô’s isometry and ergodicity, the variance of $\frac{1}{\sqrt{t}} \int_0^t (f(q_s) - \pi(f))ds$ as $t \rightarrow \infty$ is

$$\sigma^2 = \lim_{t \rightarrow \infty} \frac{2}{t} \mathbb{E} \int_0^t \nabla_p \phi^\top(q_s, p_s) \Gamma \nabla_p \phi(q_s, p_s)ds = 2 \int_{\mathbb{R}^n} \nabla_p \phi^\top \Gamma \nabla_p \phi d\tilde{\pi}.$$

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Directional derivative of σ^2 with respect to Γ

To summarise: $\frac{1}{\sqrt{T}} \int_0^T (f(q_t) - \pi(f)) dt \rightarrow \mathcal{N}\left(0, 2 \int_{\mathbb{R}^n} \nabla_p \phi^\top \Gamma \nabla_p \phi d\tilde{\pi}\right)$ in distribution, where $-\mathcal{L}\phi := (-p^\top \nabla_q + \nabla V(q)^\top \nabla_p + \Gamma p^\top \nabla_p - \nabla_p^\top \Gamma \nabla_p) \phi = f - \pi(f)$.

Theorem (Main result for underdamped Langevin dynamics)

$d\sigma^2 \cdot \delta\Gamma = -2 \int \nabla_p \phi^\top \delta\Gamma \nabla_p \tilde{\phi} d\tilde{\pi}$, where $\tilde{\phi}(q, p) = \phi(q, -p)$.

The direction $\delta\Gamma = \int \nabla_p \phi \otimes \nabla_p \tilde{\phi} d\tilde{\pi}$ guarantees a decrease in asymptotic variance!

Monte Carlo expression for solution ϕ

Claim: $\phi(q, p) = \int_0^\infty \mathbb{E}^{(q,p)}[f(q_t) - \pi(f)] dt$. ($\mathbb{E}^{(q,p)}$ denotes $(q_0, p_0) = (q, p)$)

'Proof': By the Feynman-Kac representation formula,

$$\begin{aligned} -\mathcal{L} \int_0^\infty \mathbb{E}^{(q,p)}[f(q_t) - \pi(f)] dt &= - \int_0^\infty \mathcal{L} \mathbb{E}^{(q,p)}[f(q_t) - \pi(f)] dt \\ &= - \int_0^\infty \partial_t \mathbb{E}^{(q,p)}[f(q_t) - \pi(f)] dt = f(q) - \pi(f). \end{aligned}$$

Monte Carlo expression for solution $\nabla_p \phi$

$$\nabla_p \phi(q, p) = \nabla_p \int_0^\infty \mathbb{E}[f(q_t) - \pi(f)] dt = \int_0^\infty \mathbb{E}[\nabla f(q_t)^\top \partial_{p_0} q_t] dt,$$

with initial condition $(q_0, p_0) = (q, p)$ for dynamics (q_t, p_t) .

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Simulate (q_t, p_t) and $(\partial_{p_0} q_t, \partial_{p_0} p_t)$ and increment Γ by

$$\delta\Gamma = \int \nabla_p \phi \otimes \nabla_p \tilde{\phi} d\tilde{\pi}, \quad \nabla_p \phi(q, p) = \int_0^\infty \mathbb{E}[\nabla f(q_t)^\top \partial_{p_0} q_t] dt,$$

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1D quadratic potential $V(q) = \frac{5}{2}q^2$:

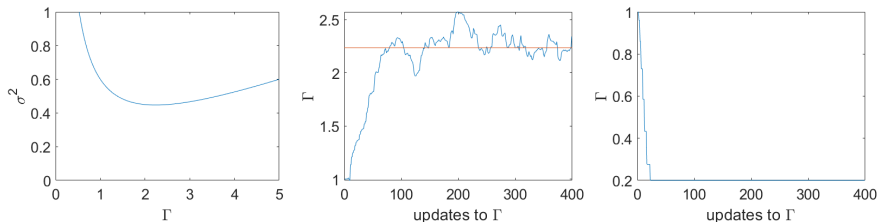


Figure: Left: Known relationship between σ^2 and Γ for $f(q) = \frac{1}{2}q^2$. Middle: changes to Γ for $f(q) = \frac{1}{2}q^2$. Right: for $f(q) = q$.

Simulate (q_t, p_t) and $(\partial_{p_0} q_t, \partial_{p_0} p_t)$ and increment Γ by

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where $(q_0, p_0) = (q, p)$.

Bayesian inference for binary regression

- 2359 data points (no. of summands in potential)
- 642 features (no. of dimensions in the SDE)

Problem: to find the posterior mean, that is, $f_i(q) = q_i$.

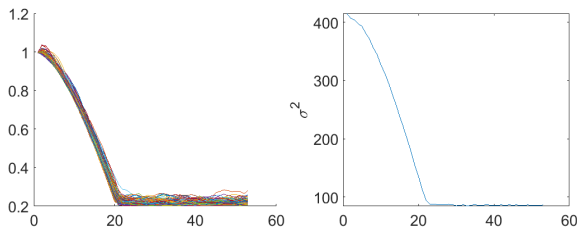


Figure: Left: Diagonal values of Γ over iterations. Right: Sum over i of estimated asymptotic variances for $f_i(q)$.

Optimal importance sampling for overdamped Langevin dynamics

joint work w/ Tony Lelièvre, Gabriel Stoltz, Urbain Vaes,

with acknowledgements to A. Duncan, G. Pavliotis

Optimal importance sampling for i.i.d. samples

Let $(X_n)_{n \in \mathbb{N}}$ be i.i.d. r.v.'s drawn from $\pi_U(dx) \propto e^{-V(x)-U(x)} dx$.

To estimate $I = \int f d\pi$, use that

$$\frac{\frac{1}{N} \sum_{i=1}^N f(X_n) e^{U(X_n)}}{\frac{1}{N} \sum_{n=1}^N e^{U(X_n)}} \rightarrow \frac{\int f e^U d\pi_U}{\int e^U d\pi_U} = I$$

almost surely as $N \rightarrow \infty$. The CLT and Slutsky's theorem gives

$$\begin{aligned} \sqrt{N} \left(\frac{\frac{1}{N} \sum_{i=1}^N f(X_n) e^{U(X_n)}}{\frac{1}{N} \sum_{n=1}^N e^{U(X_n)}} - I \right) &= \frac{\frac{1}{\sqrt{N}} \sum_{n=1}^N (f(X_n) - I) e^{U(X_n)}}{\frac{1}{N} \sum_{n=1}^N e^{U(X_n)}} \\ &\rightarrow \mathcal{N} \left(0, \frac{Z_U^2}{Z^2} \int |(f - I) e^U|^2 d\pi_U \right) \end{aligned}$$

in distribution as $N \rightarrow \infty$, where Z_U and Z are the normalization constants for $\pi_U(dx) \propto e^{-V(x)-U(x)} dx$ and $\pi(dx) \propto e^{-V(x)} dx$ respectively.

By Cauchy-Schwarz,

$$\left(\int |f - I| e^{-V} \right)^2 \leq \left(\int |f - I|^2 e^{U-V} \right) \left(\int e^{-U-V} \right) = Z_U^2 \int |(f - I) e^U|^2 d\pi_U,$$

with equality for $U = -\log|f - I|$.

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Overdamped Langevin $dq_t = -\nabla V(q_t)dt - \nabla U(q_t)dt + \sqrt{2}dW_t$

By the same ideas as before, there is an explicit CLT variance

$$\sqrt{T} \frac{\frac{1}{T} \int_0^T (f(q_t) - \pi(f)) e^{U(q_t)} dt}{\frac{1}{T} \int_0^T e^{U(q_t)} dt} \rightarrow \mathcal{N}\left(0, \frac{2Z_U^2}{Z^2} \int |\nabla \phi|^2 d\pi_U\right),$$

where ϕ solves $-\mathcal{L}\phi = f - \pi(f)$.

Theorem (Explicit optimal U on \mathbb{R} (when $U = U : \mathbb{R} \rightarrow \mathbb{R}$))

Under mild assumptions, the minimum asymptotic variance is achieved by

$$U(x) = -V(x) - \log \left| \int_{-\infty}^x (f(y) - \pi(f)) e^{-V(y)} dy \right|.$$

Theorem (Directional derivative of $\sigma^2 = \frac{2Z_U^2}{Z^2} \int |\nabla \phi|^2 d\pi_U$ with respect to U)

$$d\sigma^2 \cdot \delta U = \frac{2Z_U^2}{Z^2} \int \delta U \left(|\nabla \phi|^2 - \int |\nabla \phi|^2 d\pi_U \right) d\pi_U.$$

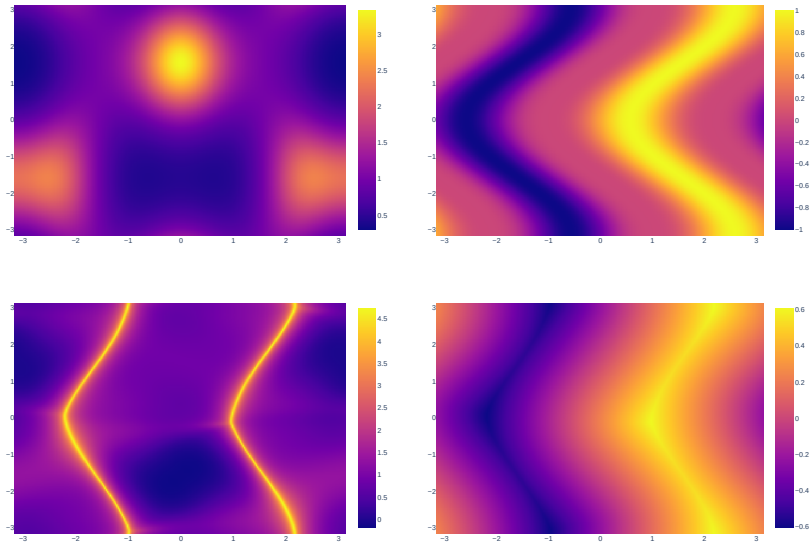


Figure: Unperturbed potential V (top left), observable f (top right), optimal potential $V + U$ (bottom left), and corresponding solution to the Poisson equation ϕ_U (bottom right).

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