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# Can Monte Carlo methods be used to simulate active-matter systems?

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(3) Department of Physics, Massachusetts Institute of Technology, Cambridge, Massachusetts 02139, USA.

# What is Active Matter?

Brownian motion

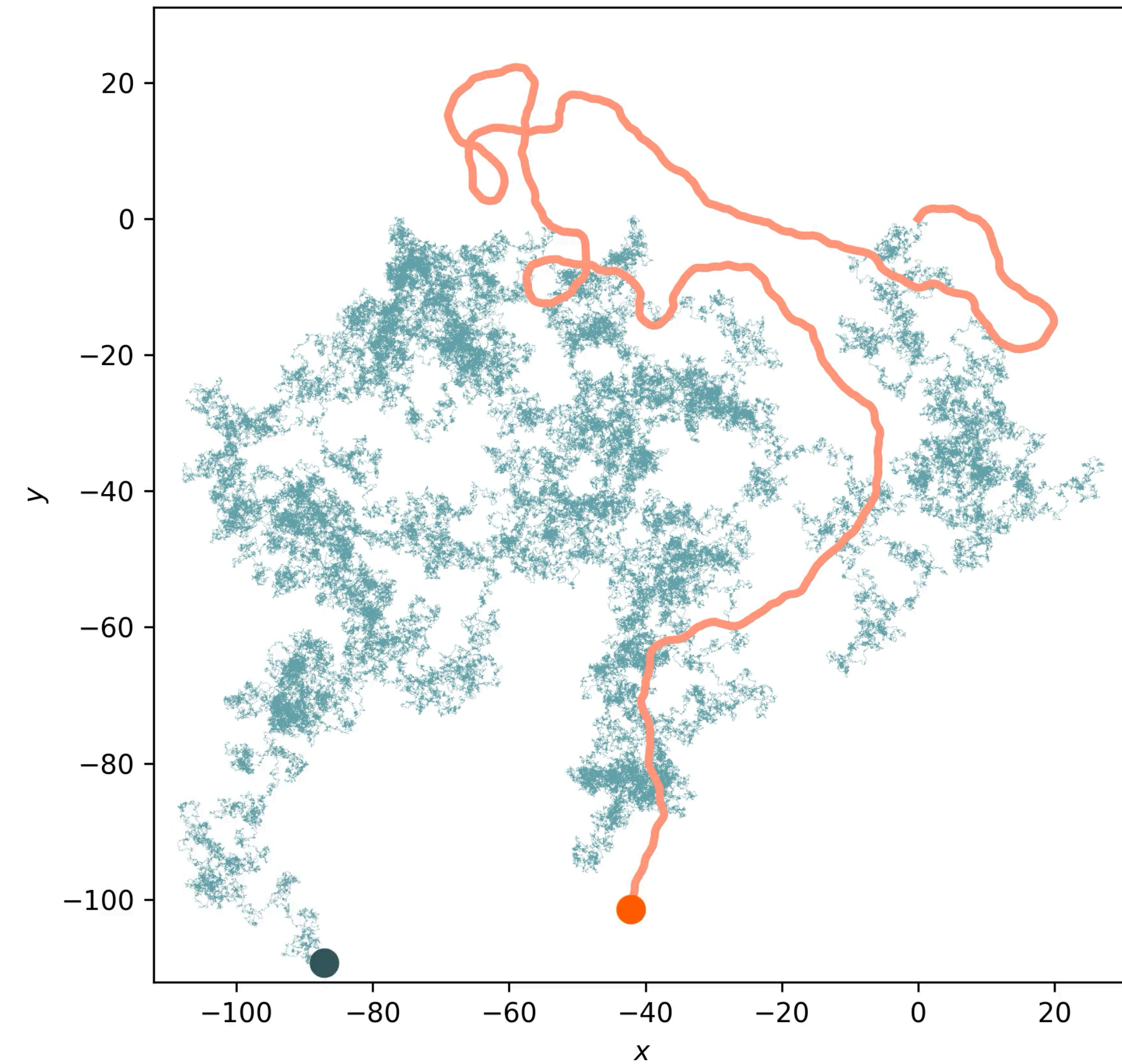
Active Brownian particle

$$\mathbf{v}(t) := \frac{d\mathbf{r}(t)}{dt} = \sqrt{2D} \boldsymbol{\eta}(t)$$

$$\langle \eta_x(t) \eta_y(t') \rangle = \delta_{x,y} \delta(t - t')$$

Detailed Balance

NO Detailed Balance





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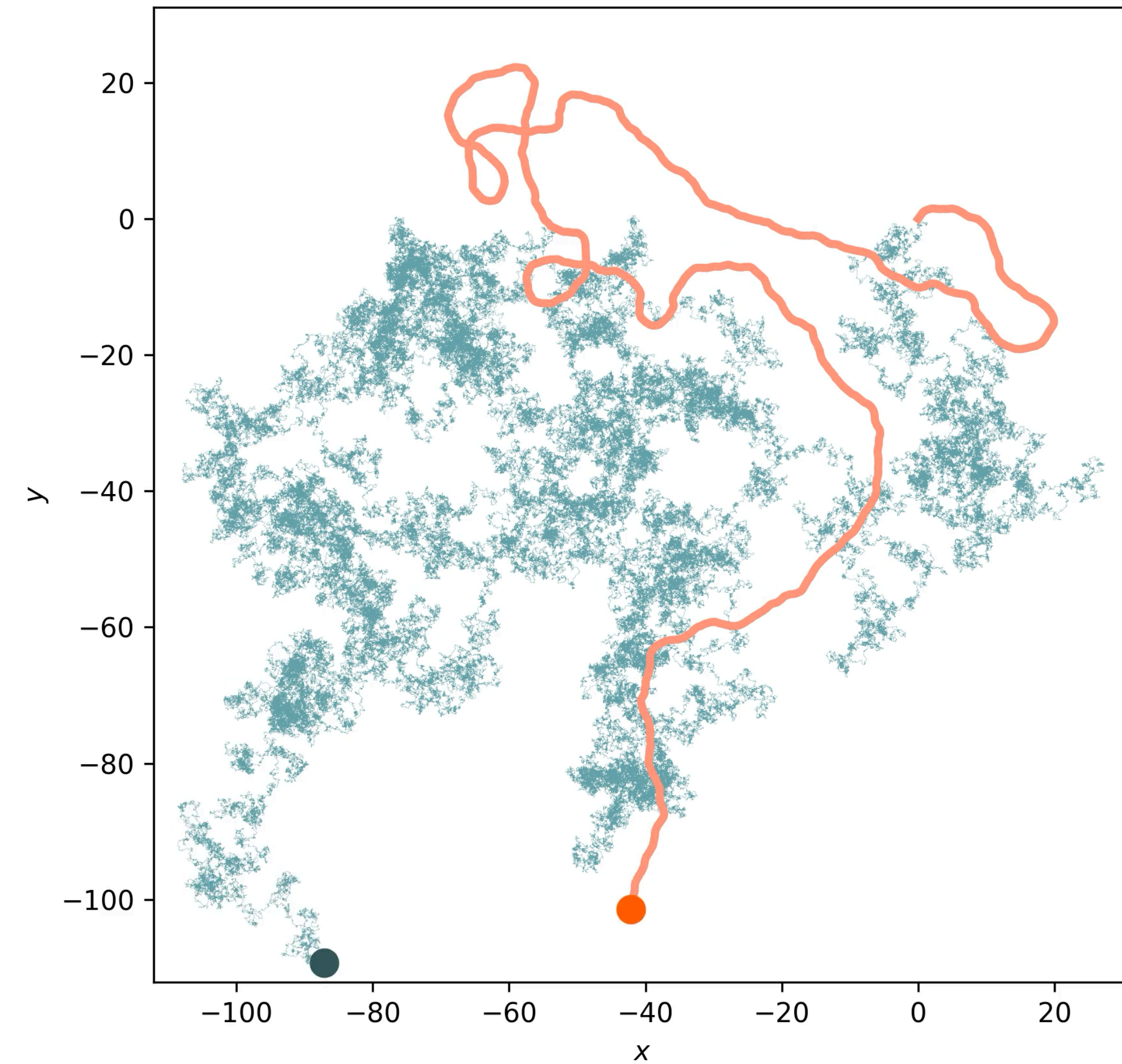
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**Detailed Balance**

## Active Brownian particle

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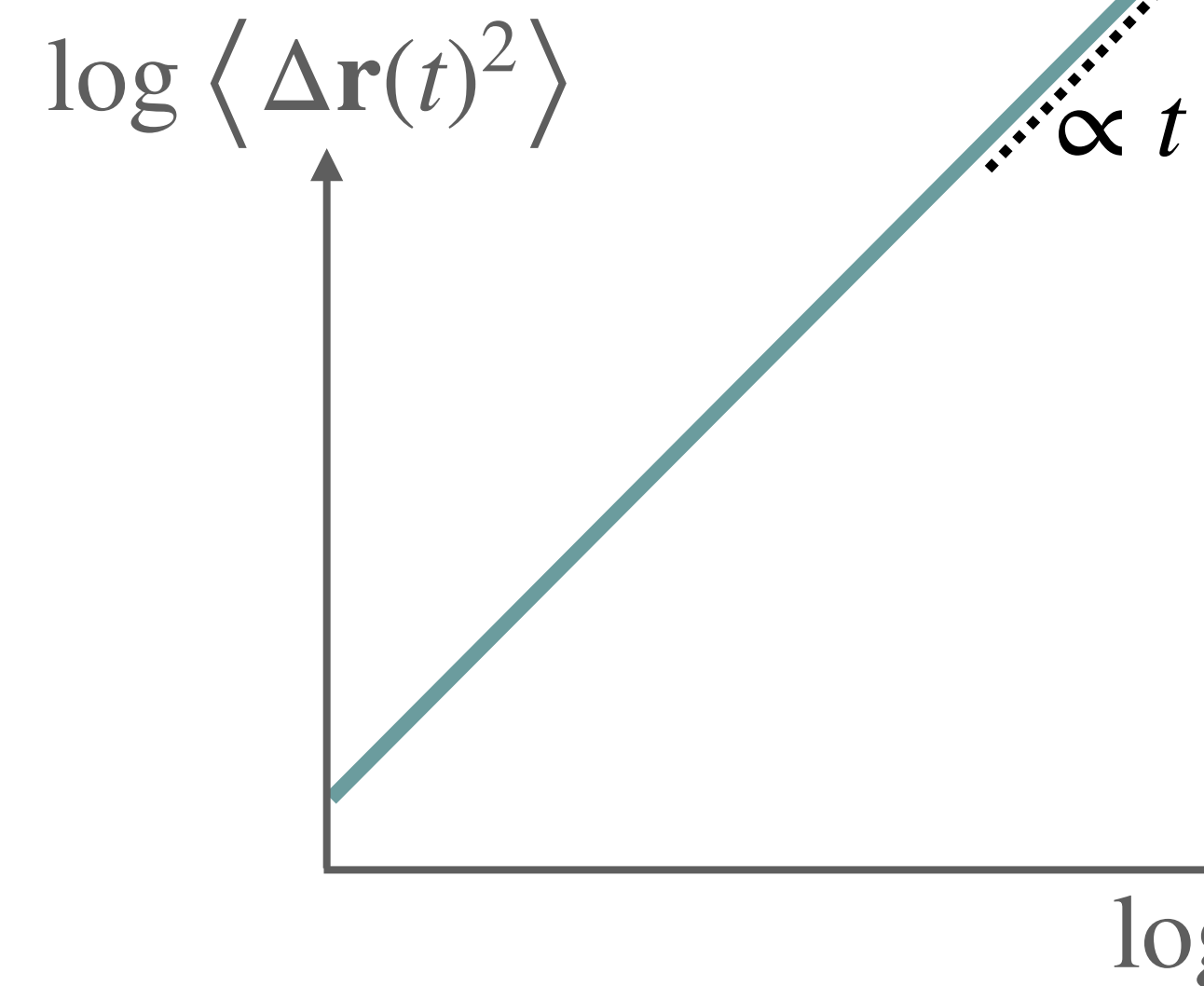
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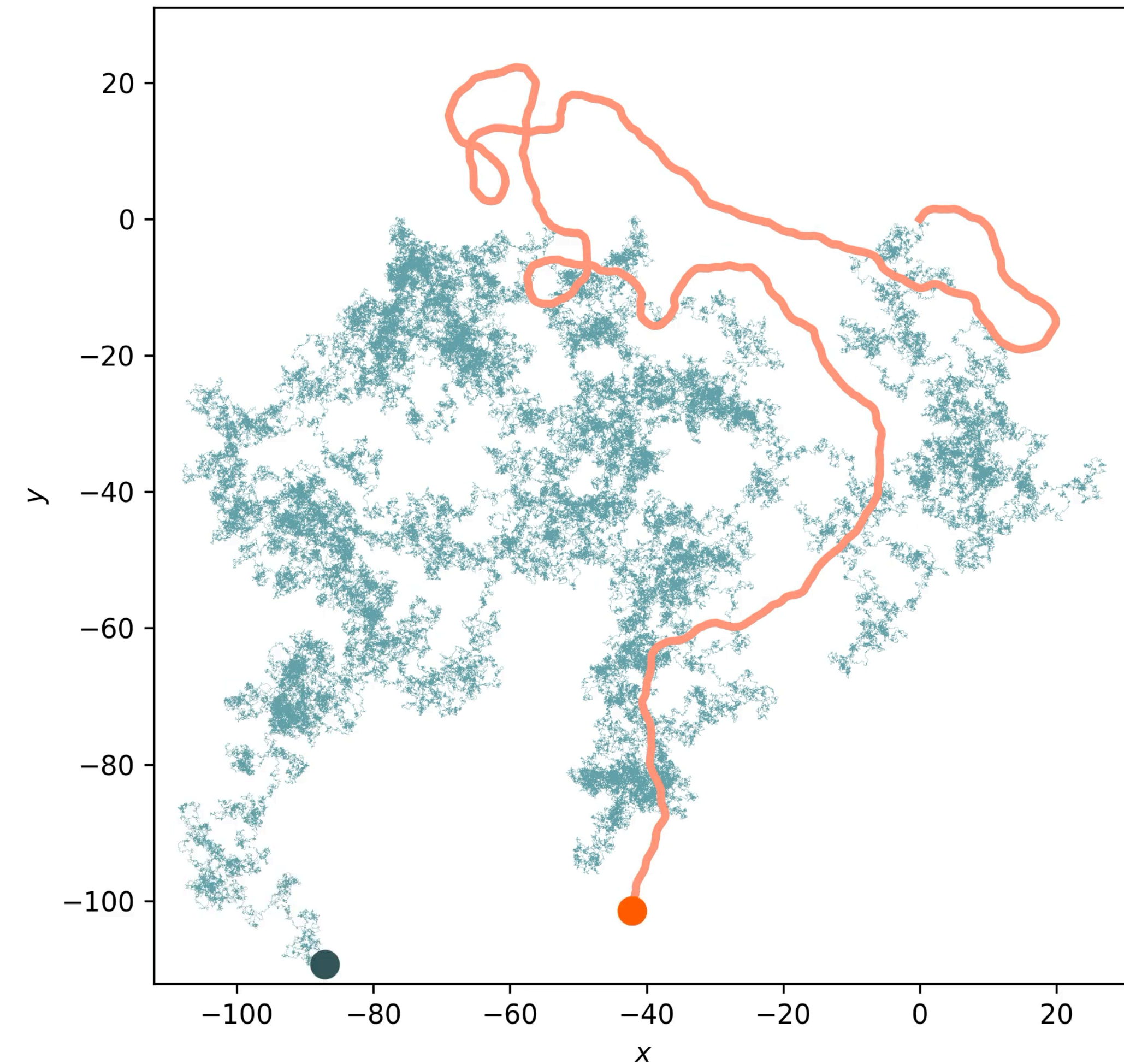
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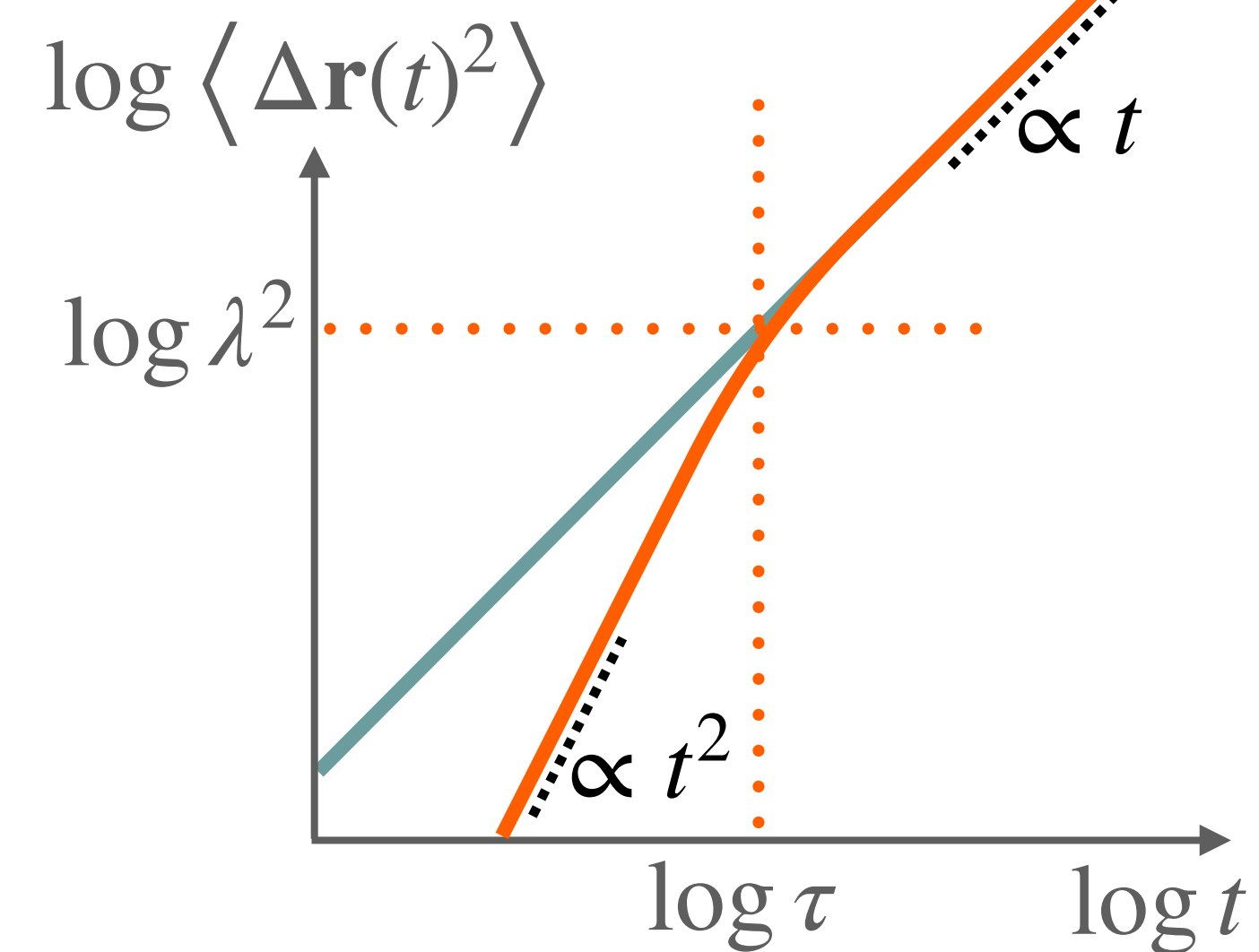
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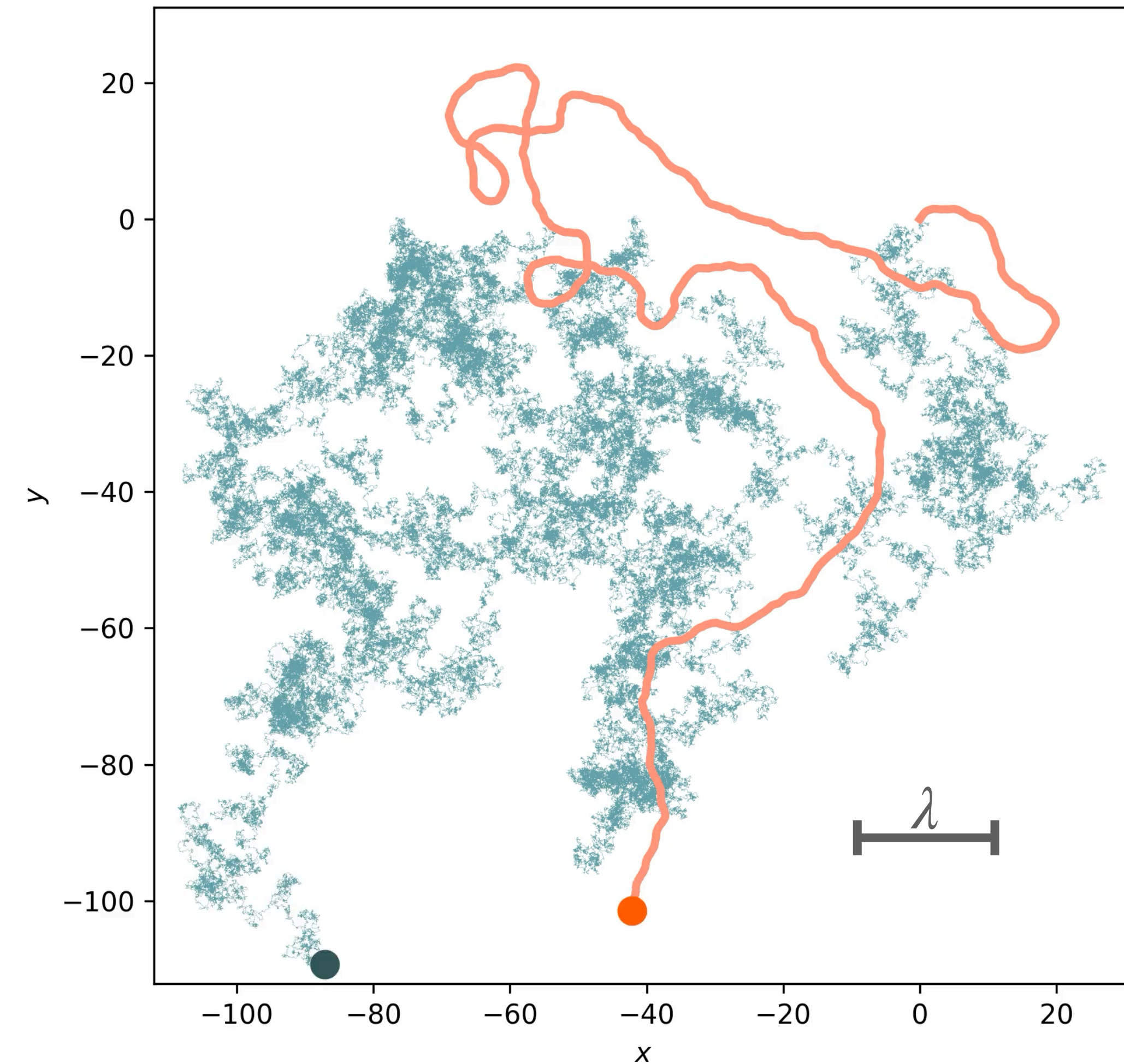
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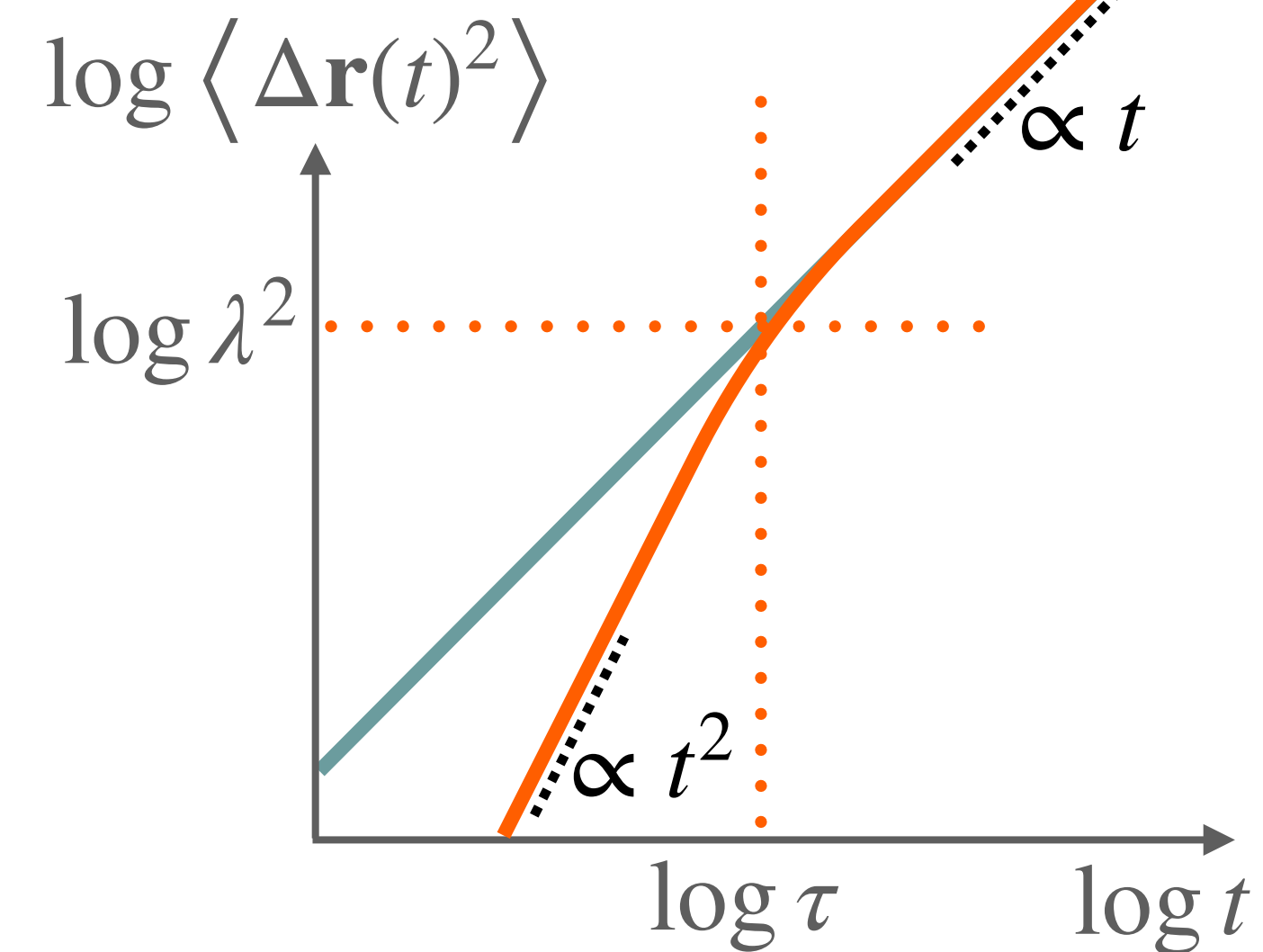
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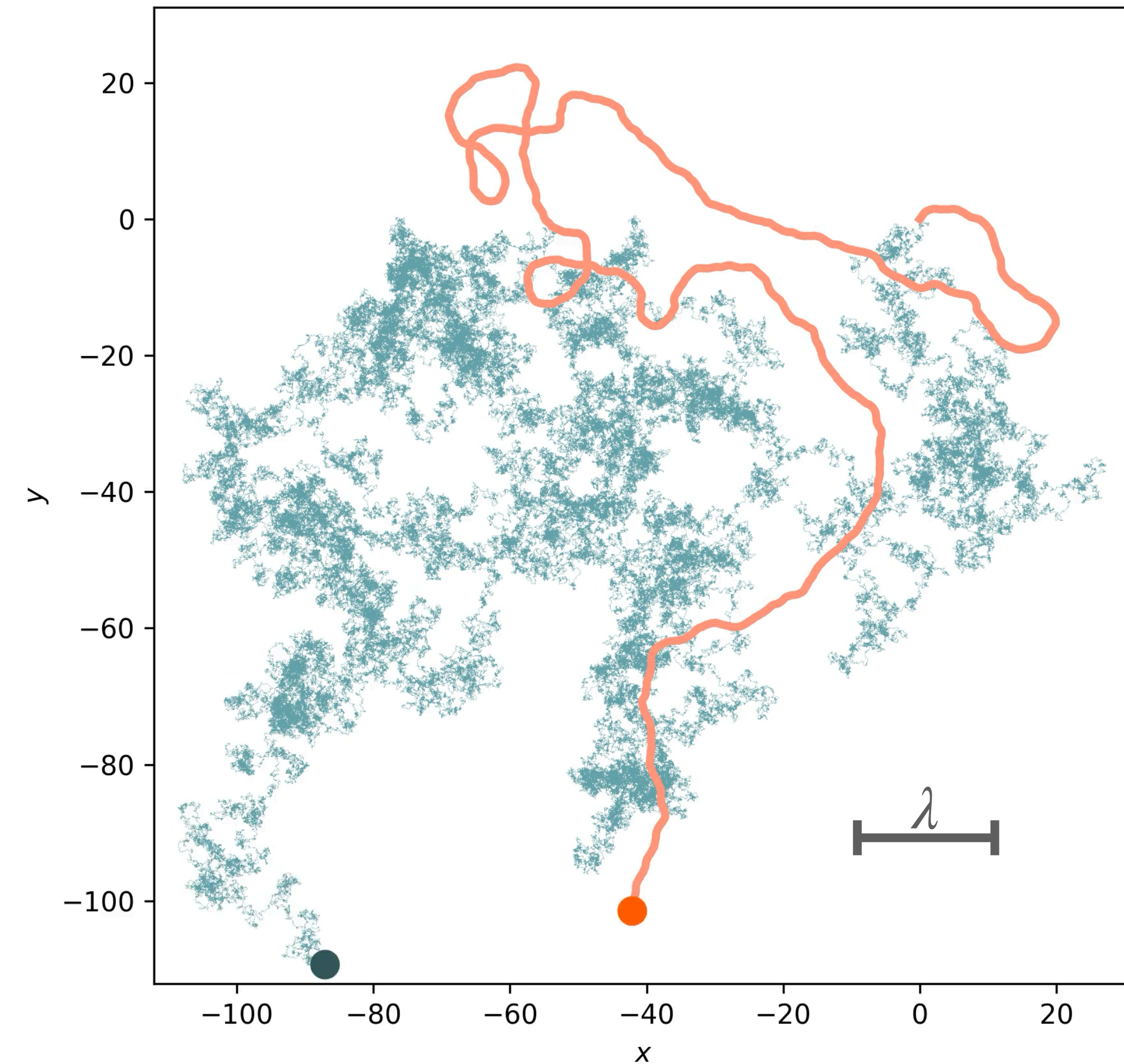
## Detailed Balance



## Active Brownian particle

A diagram of an active Brownian particle, represented by an orange circle. An arrow labeled  $\mathbf{n}$  points from the center of the circle at an angle  $\theta$  relative to a horizontal line. To the right of the diagram is the equation  $\frac{d\mathbf{r}(t)}{dt} = v_0 \mathbf{n}(\theta_t)$ .

## NO Detailed Balance





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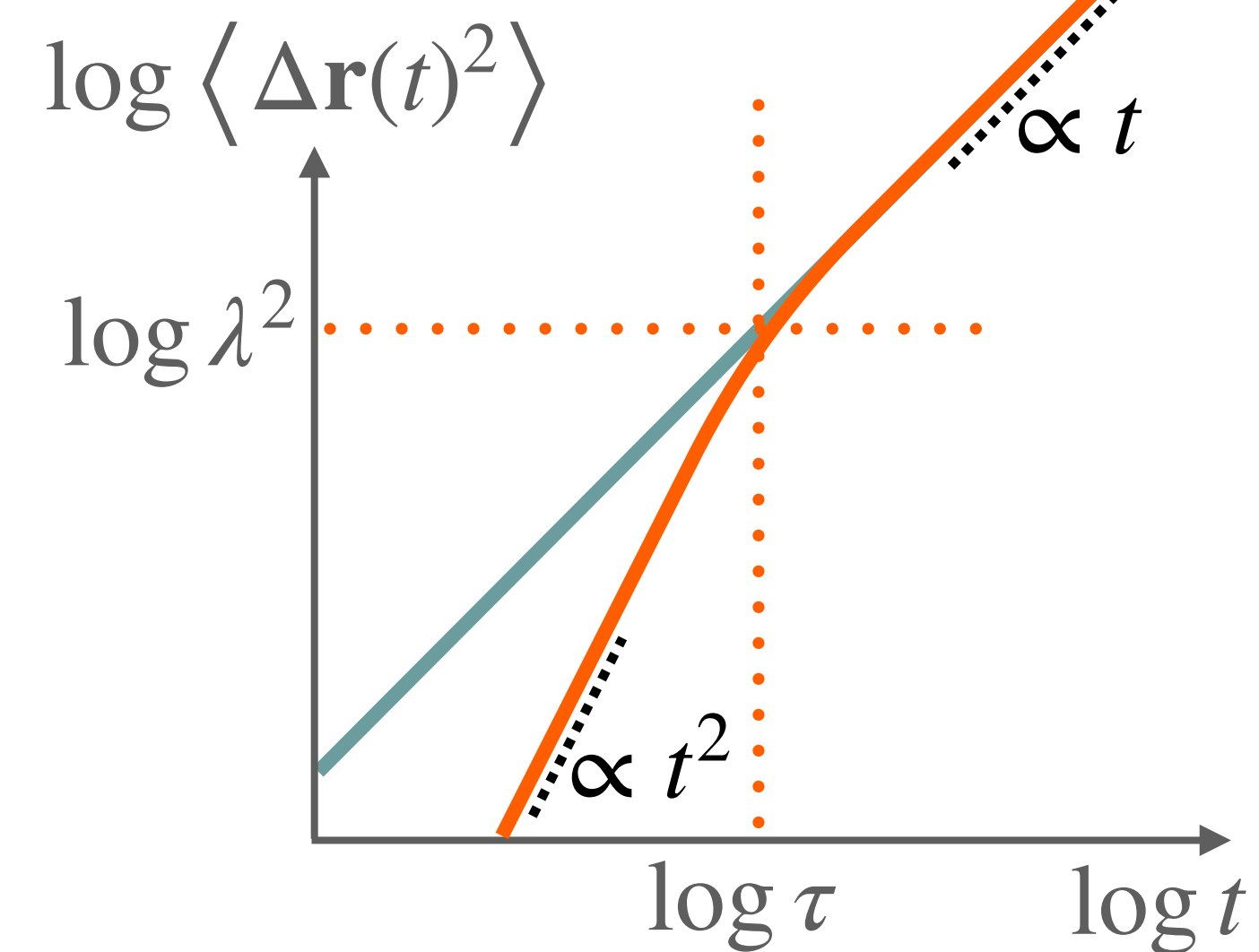
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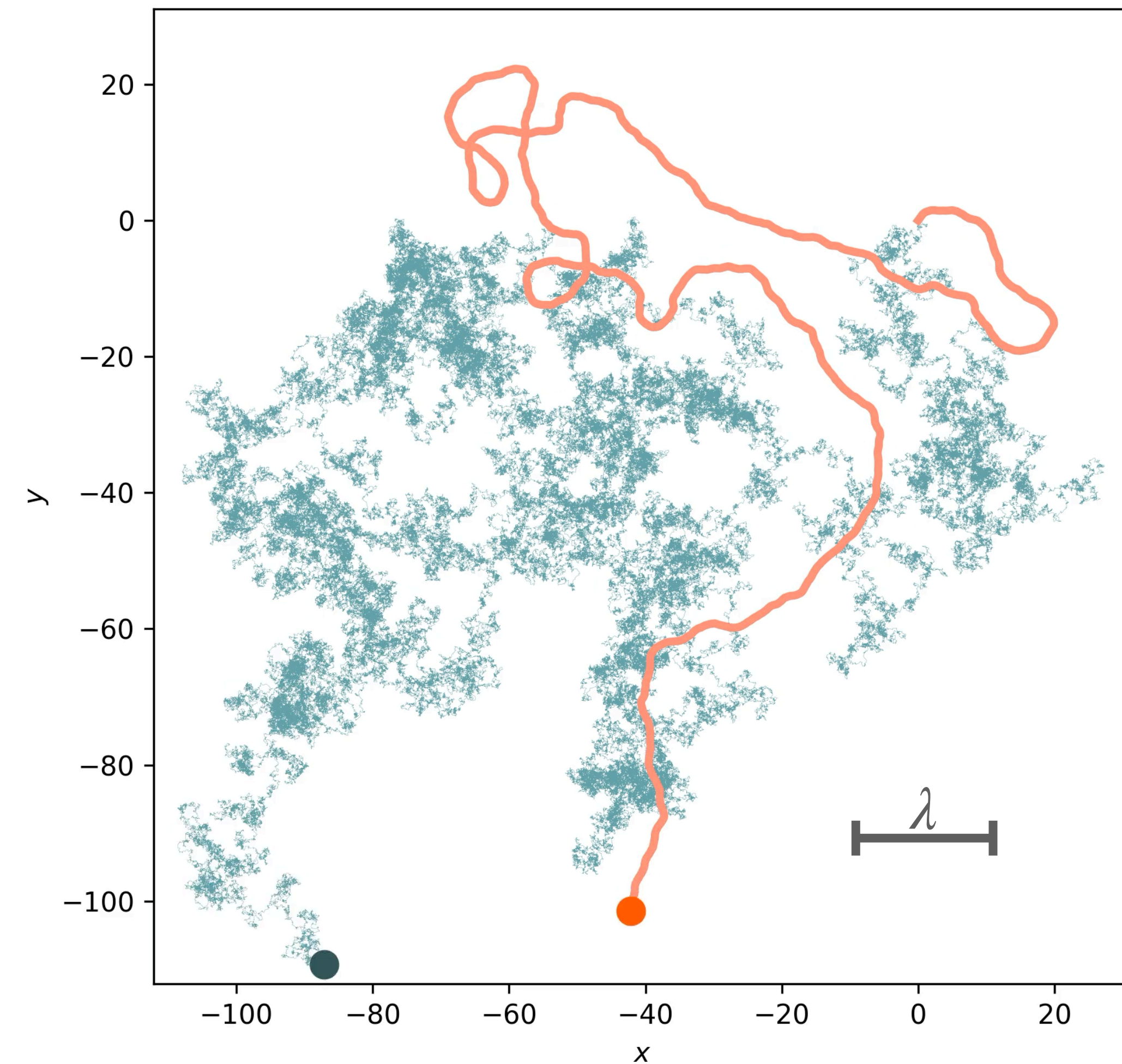
## Active Brownian particle

A diagram of an active Brownian particle, represented by an orange circle. An arrow labeled  $\mathbf{n}$  points from the center of the circle, and an angle  $\theta$  is indicated between the arrow and a horizontal line. To the right of the diagram is the equation:

$$\frac{d\mathbf{r}(t)}{dt} = v_0 \mathbf{n}(\theta_t)$$

$$\frac{d\theta(t)}{dt} = \sqrt{2\tau^{-1}} \xi(t)$$

## NO Detailed Balance





# What is Active Matter?

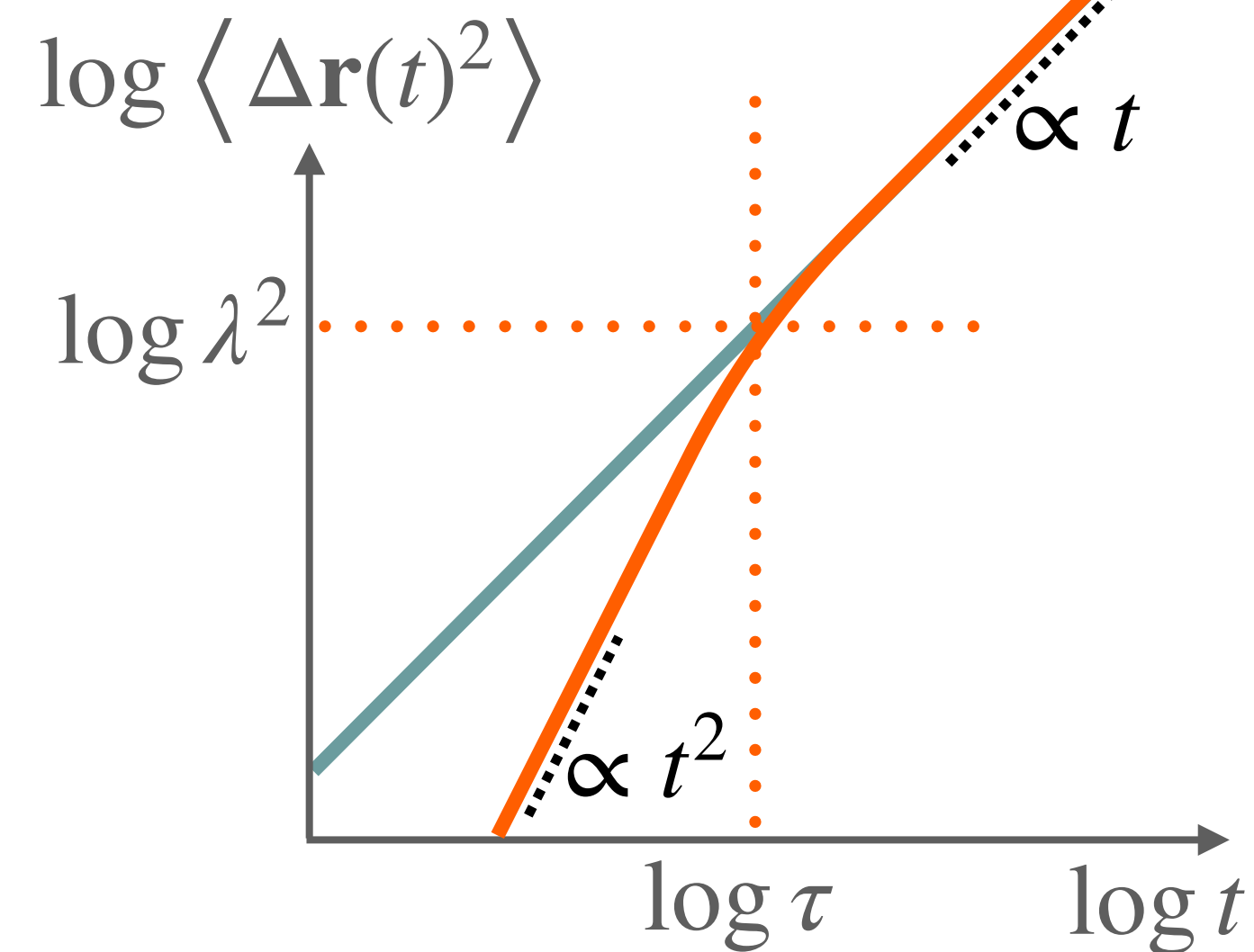
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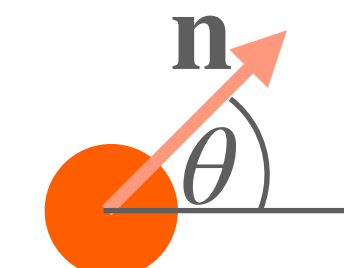
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## Detailed Balance



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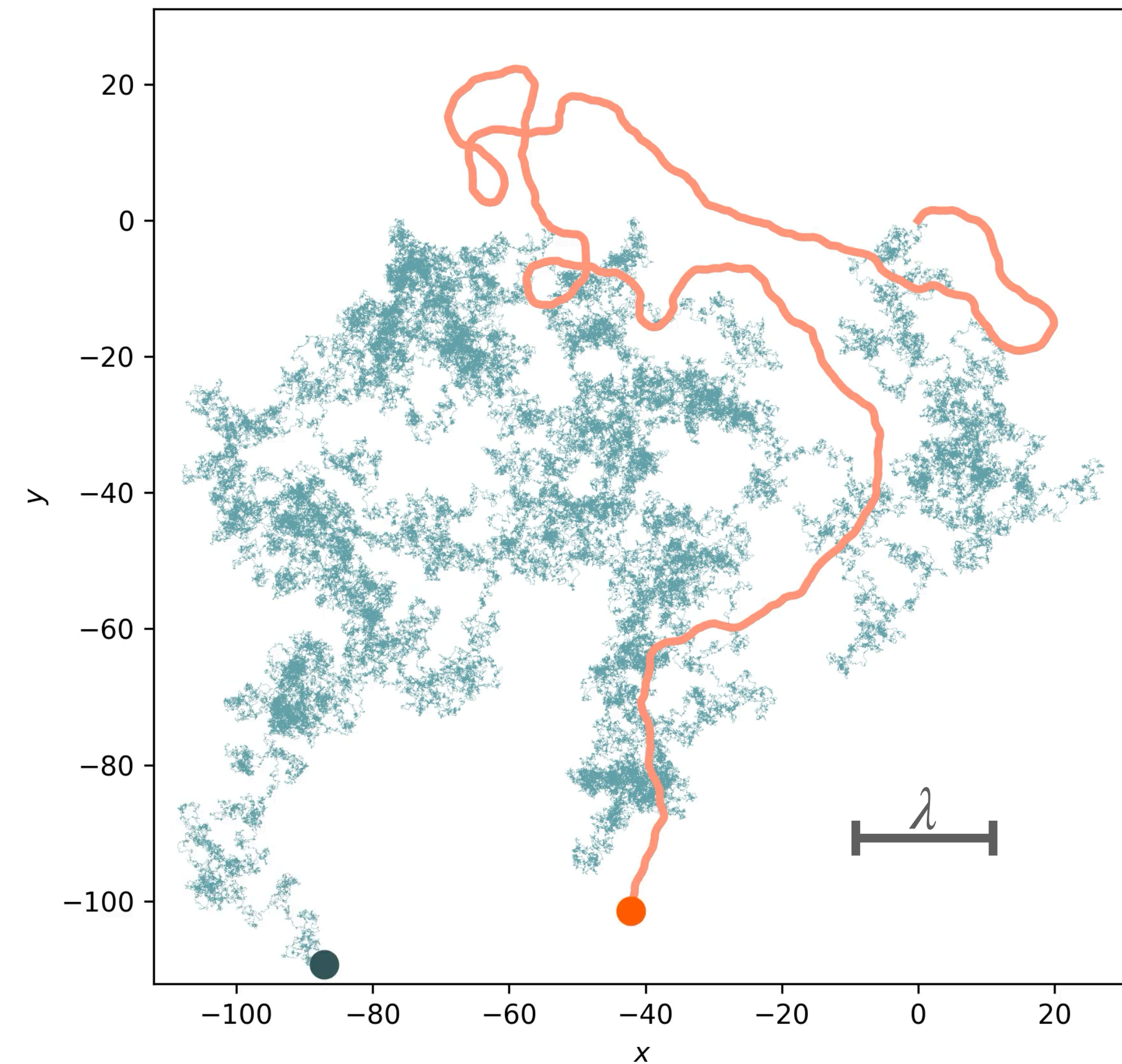
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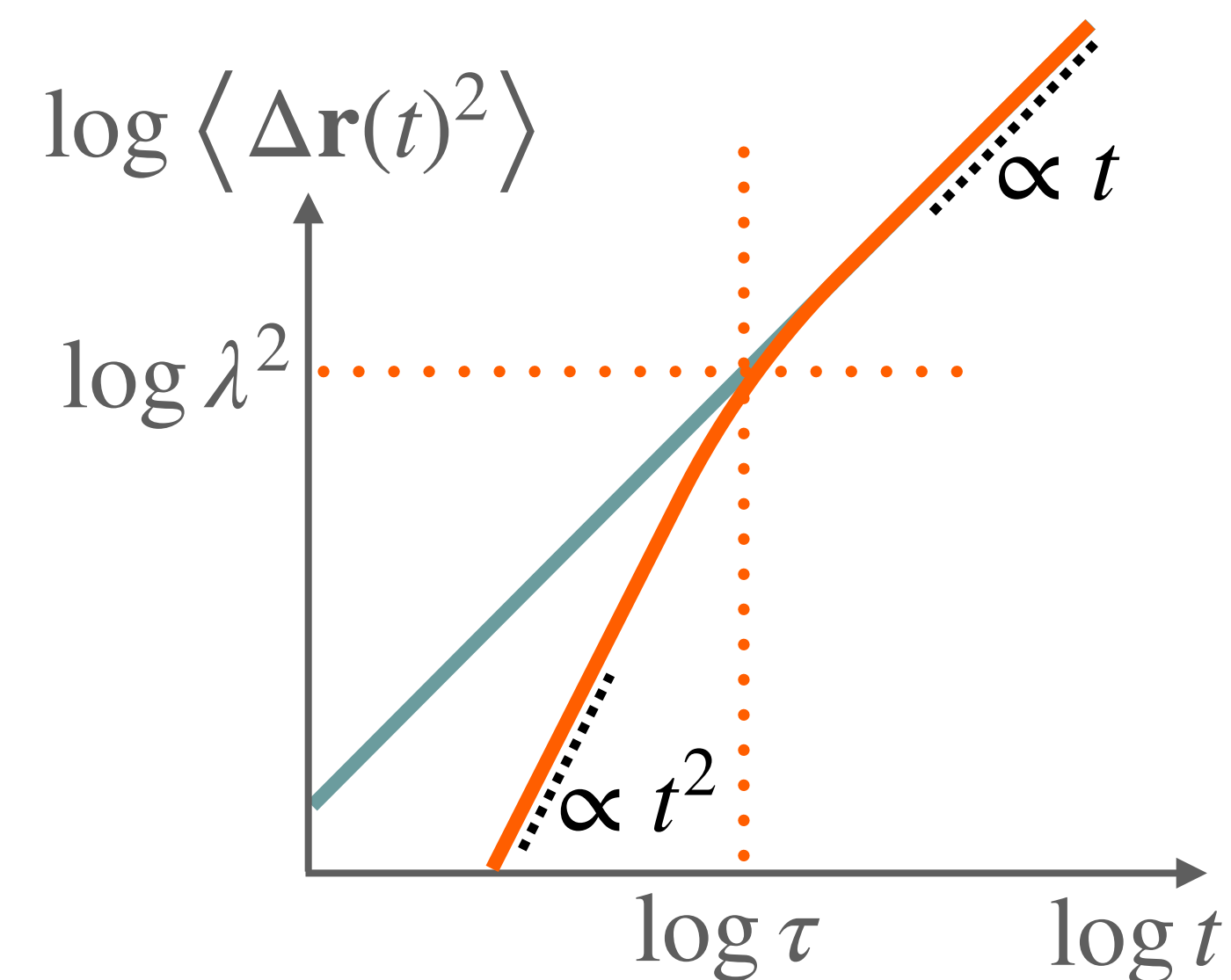
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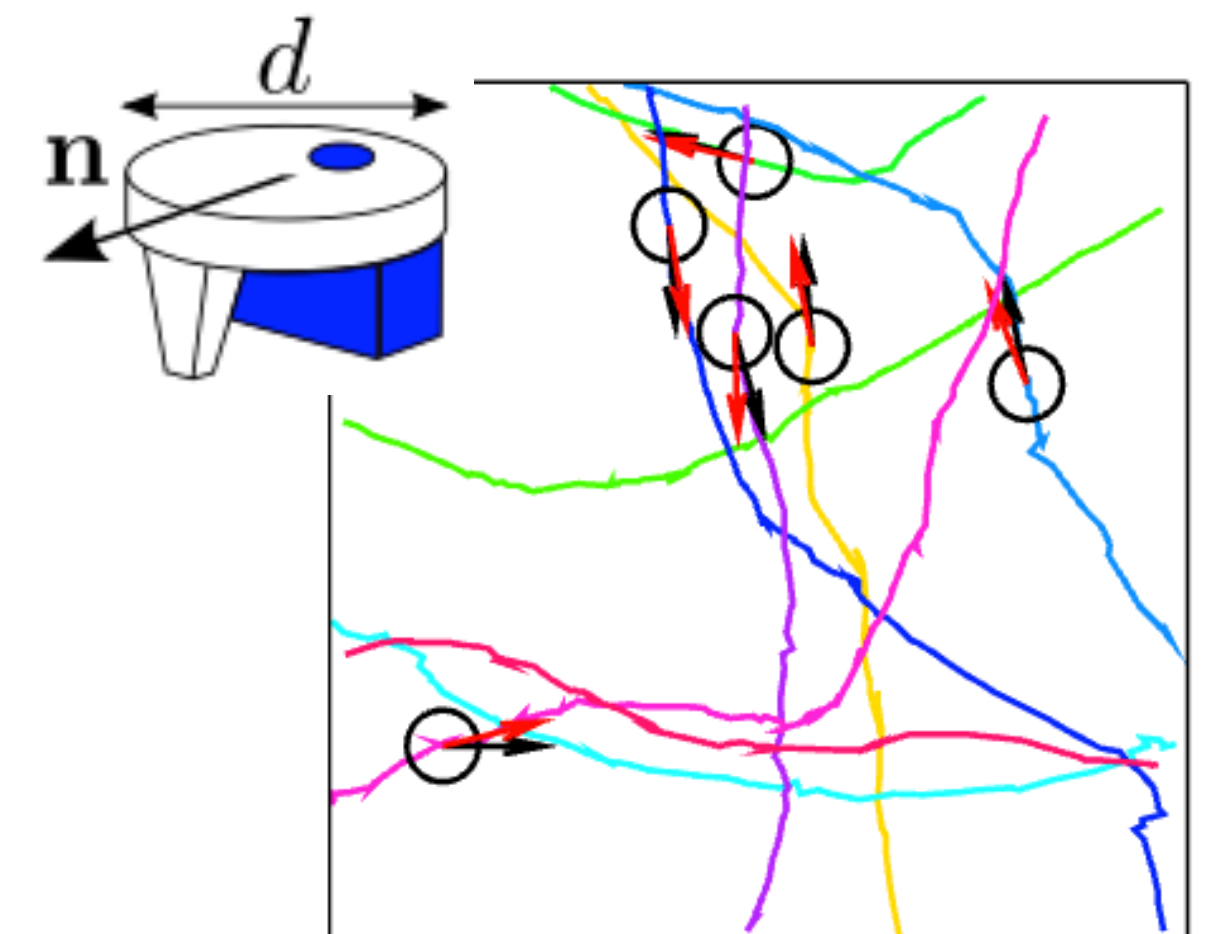
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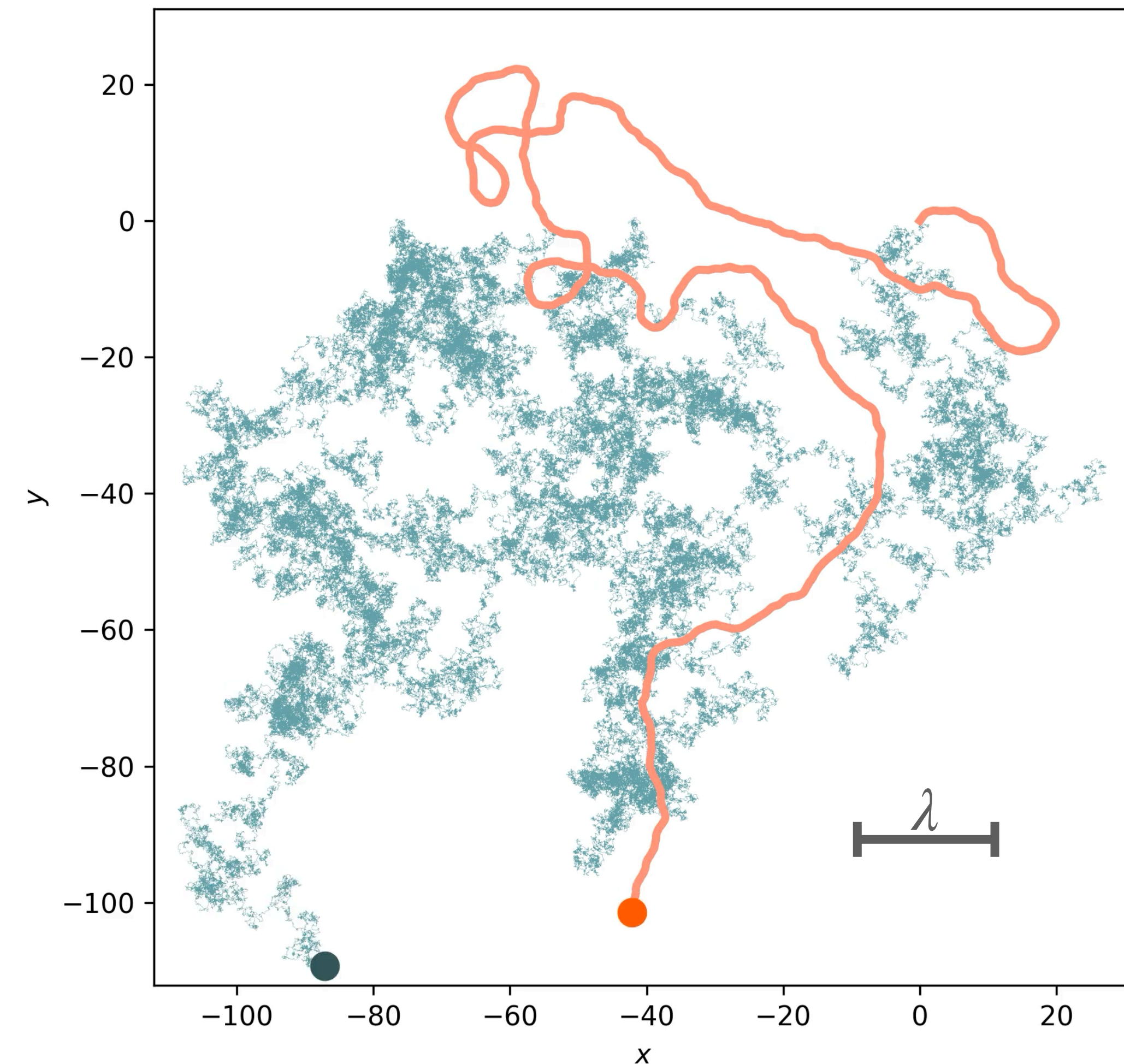
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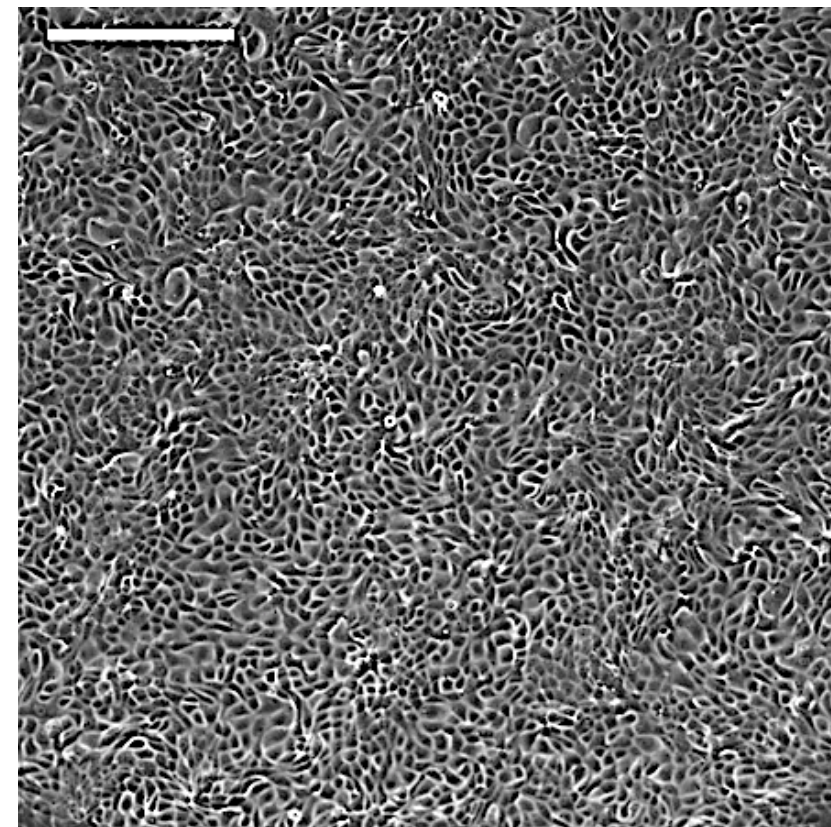
J. Deseigne, O. Dauchot, H. Chaté  
PRL 105, 098001 (2010).





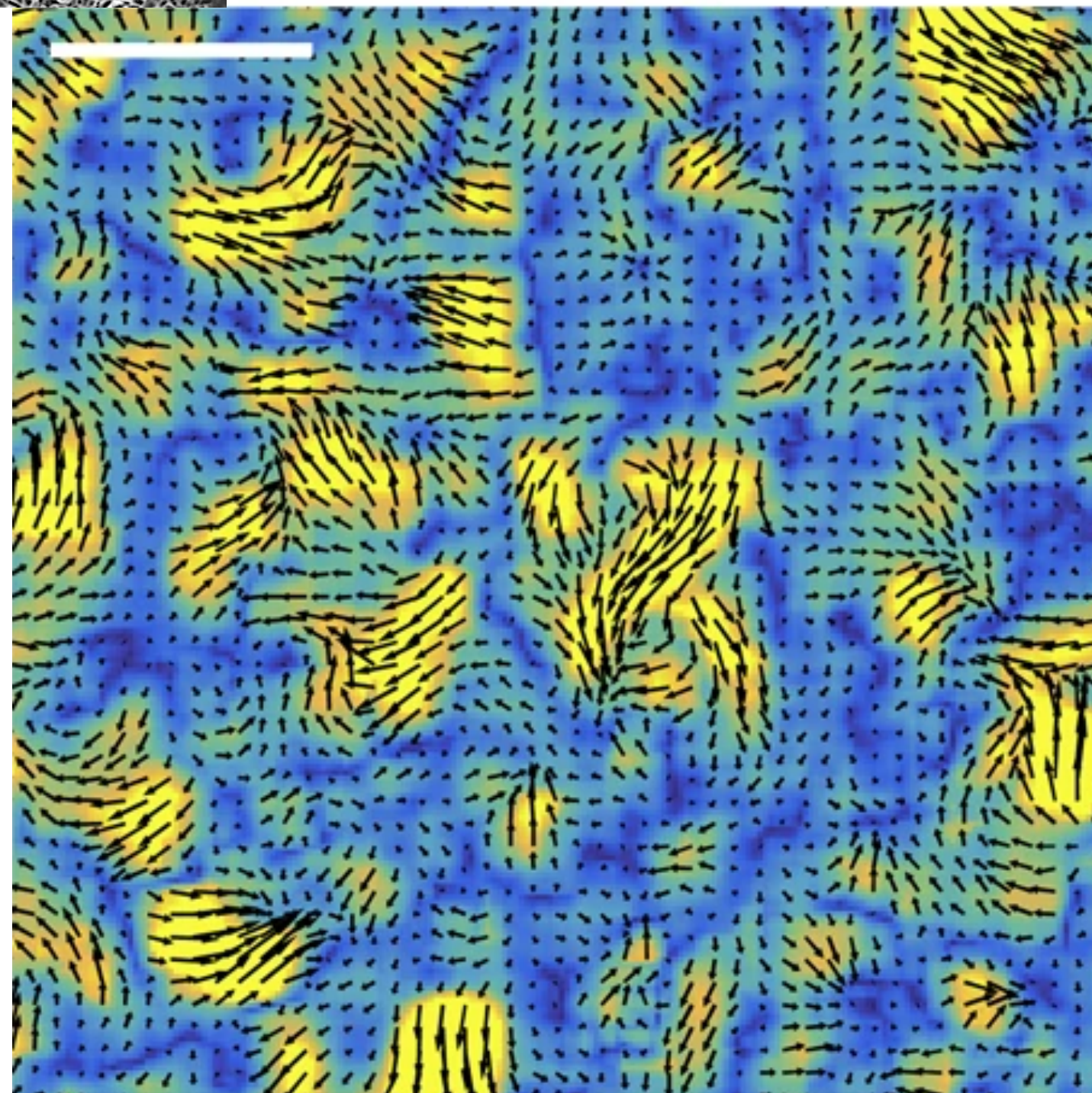
# Collective phenomena in active matter

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Complex chaotic flows  
in monolayers of cells

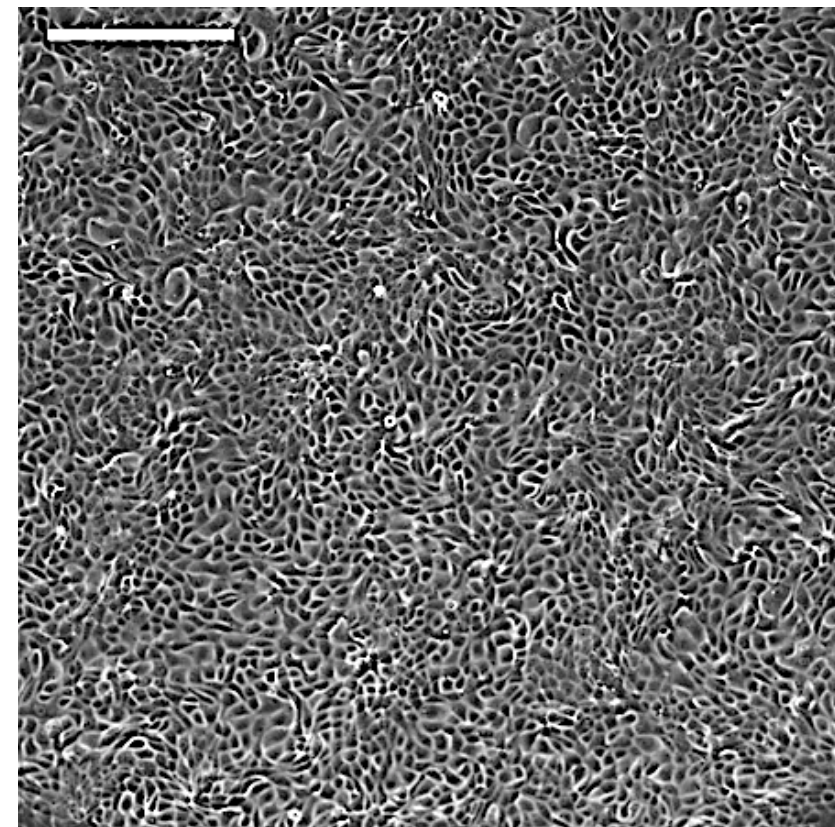
0 15 30  $V$  ( $\mu\text{m}\cdot\text{h}^{-1}$ )



SZ Lin et al. Communications Physics (2021).

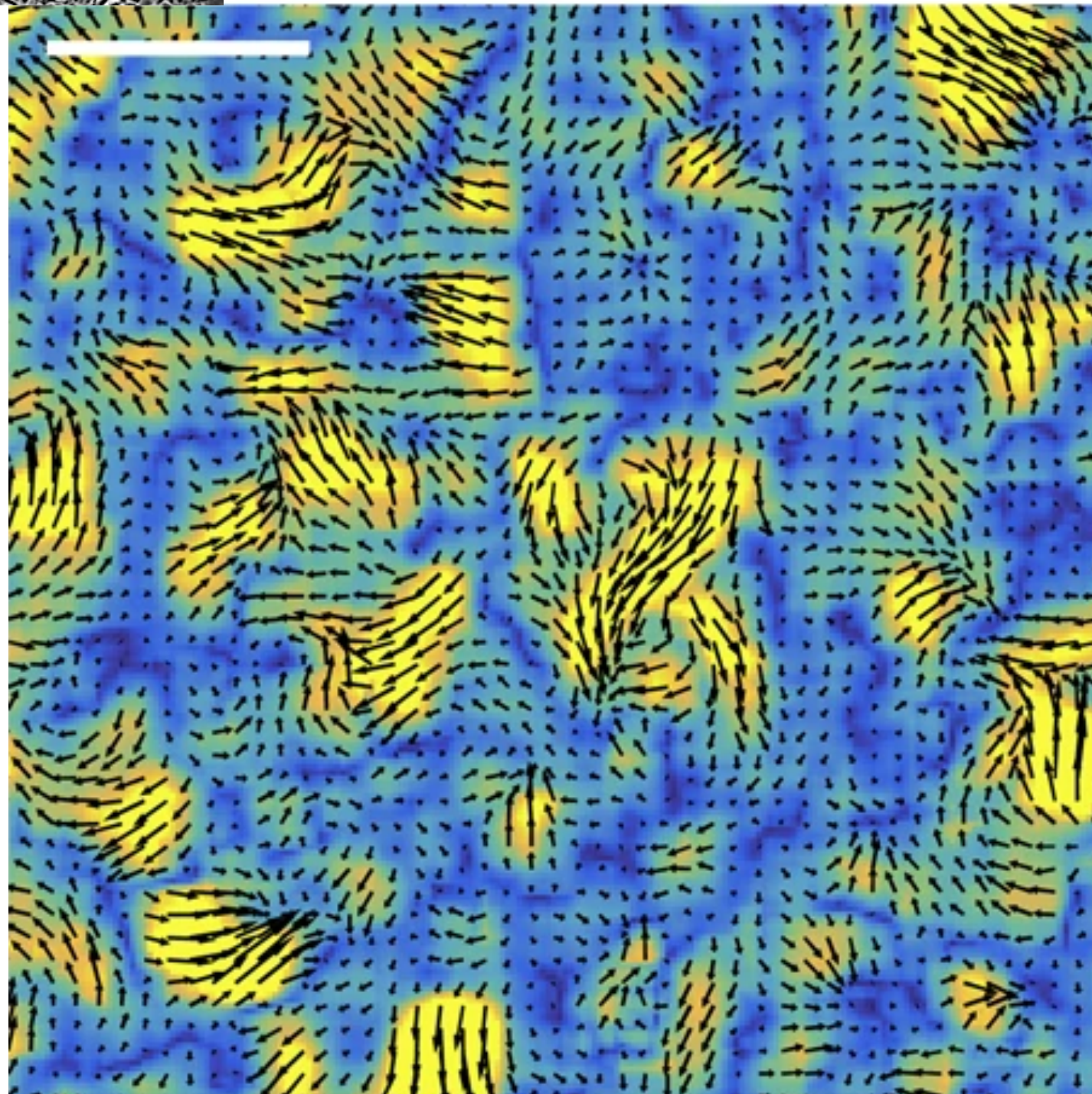


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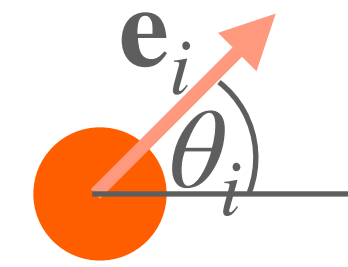
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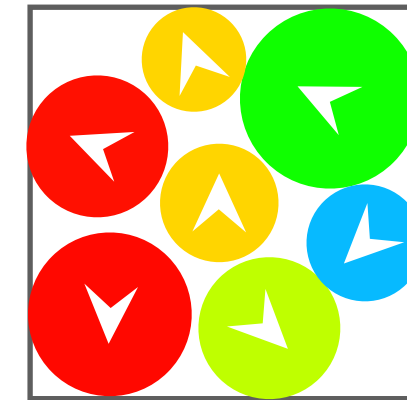


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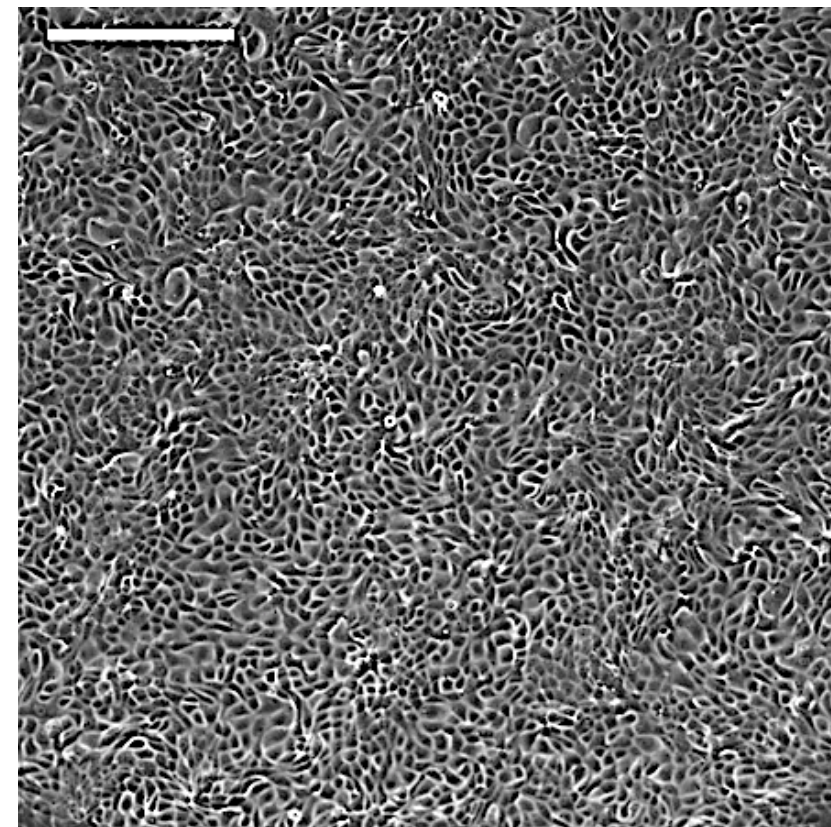


$$\frac{d\mathbf{r}_i(t)}{dt} = v_0 \mathbf{n}(\theta_i[t]) + \sum \mathbf{F}_{ij} \quad \frac{d\theta_i(t)}{dt} = \sqrt{2\tau^{-1}} \xi_i(t)$$



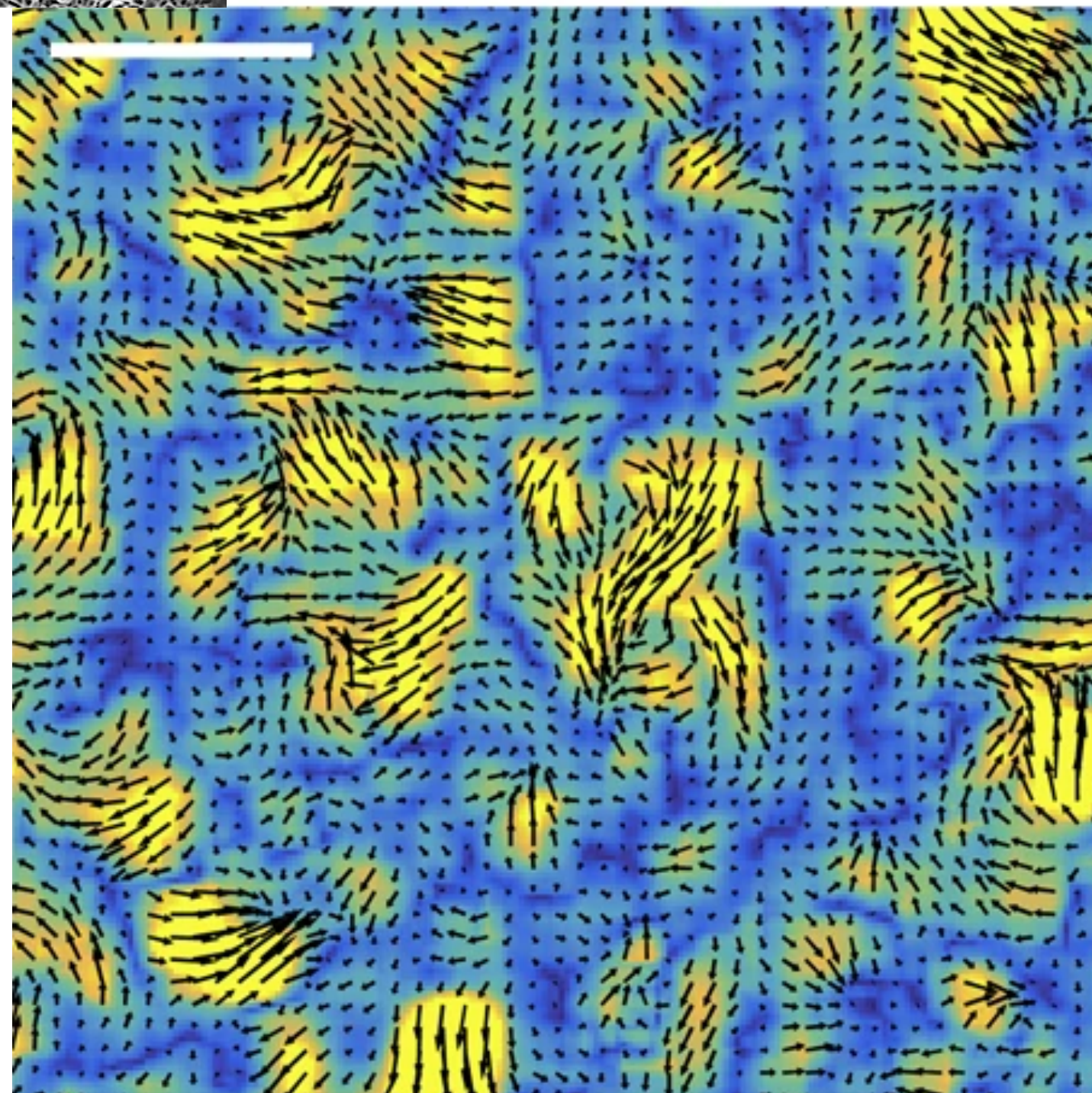


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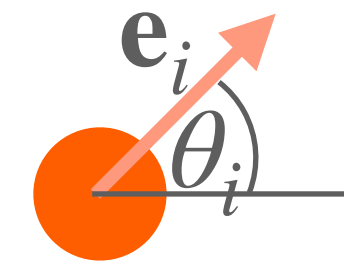
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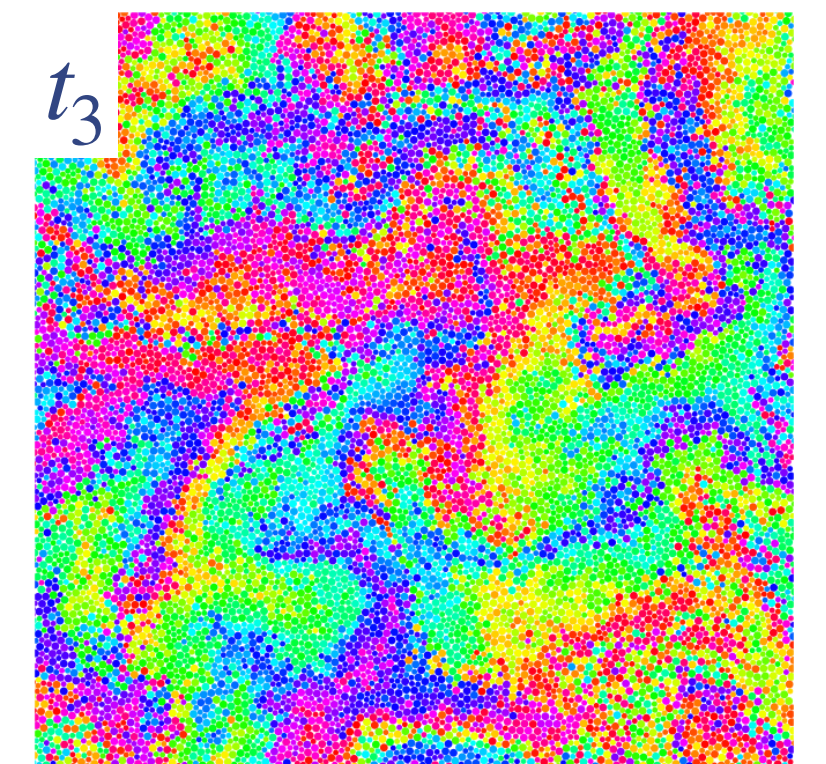
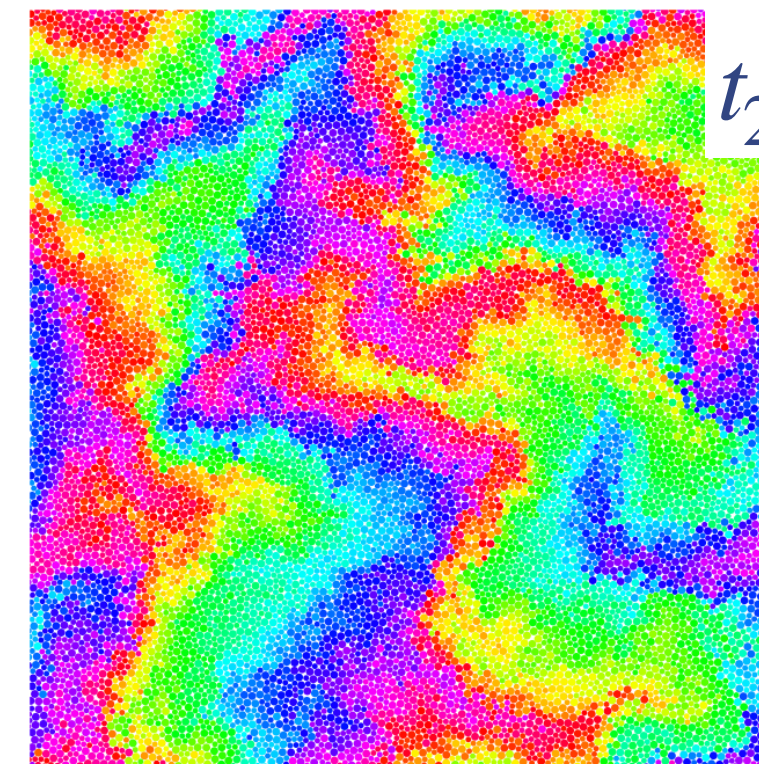
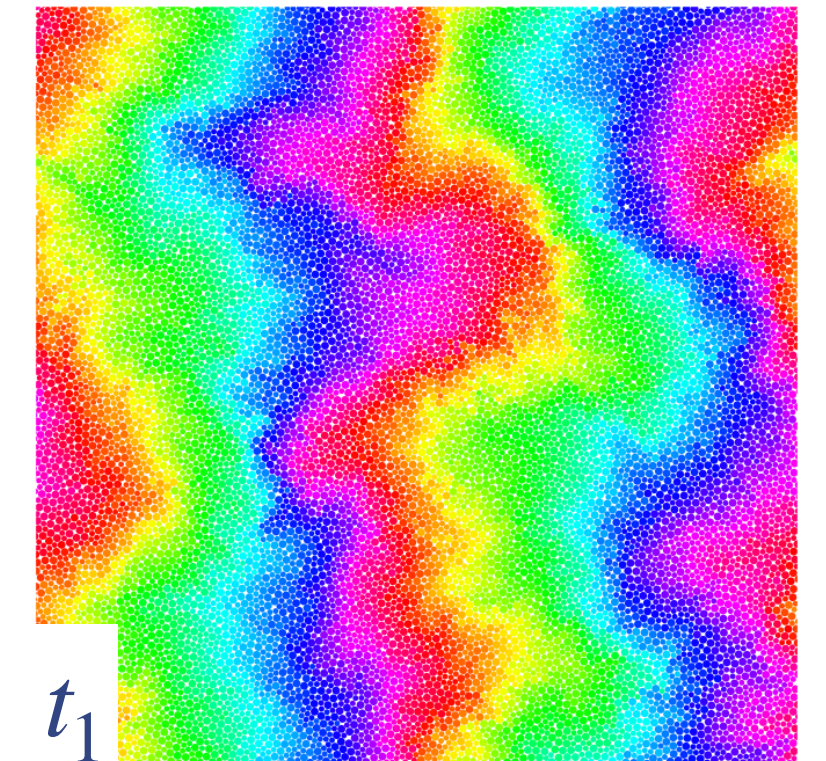
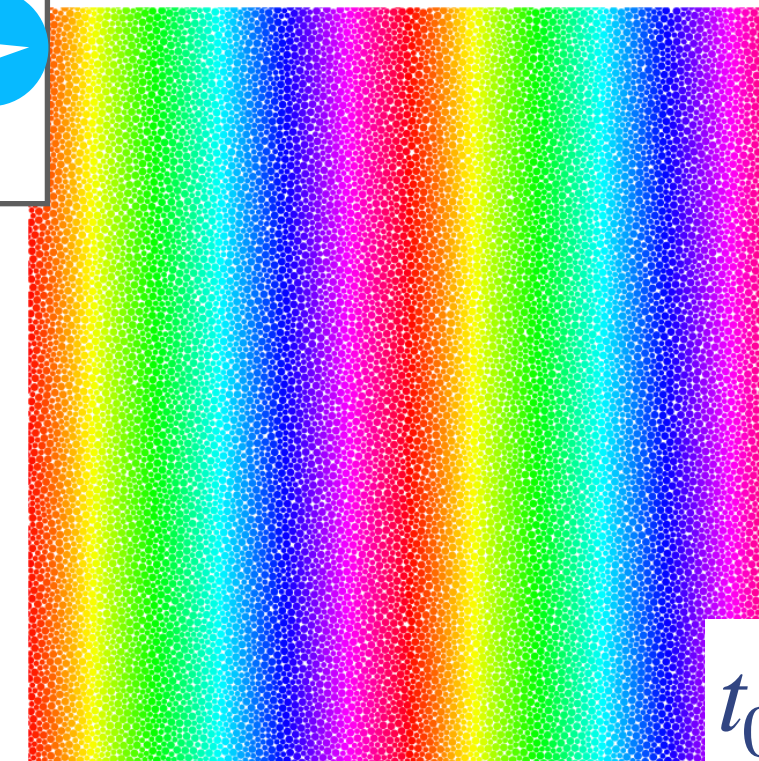
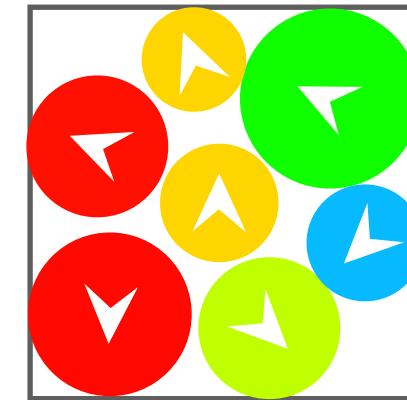
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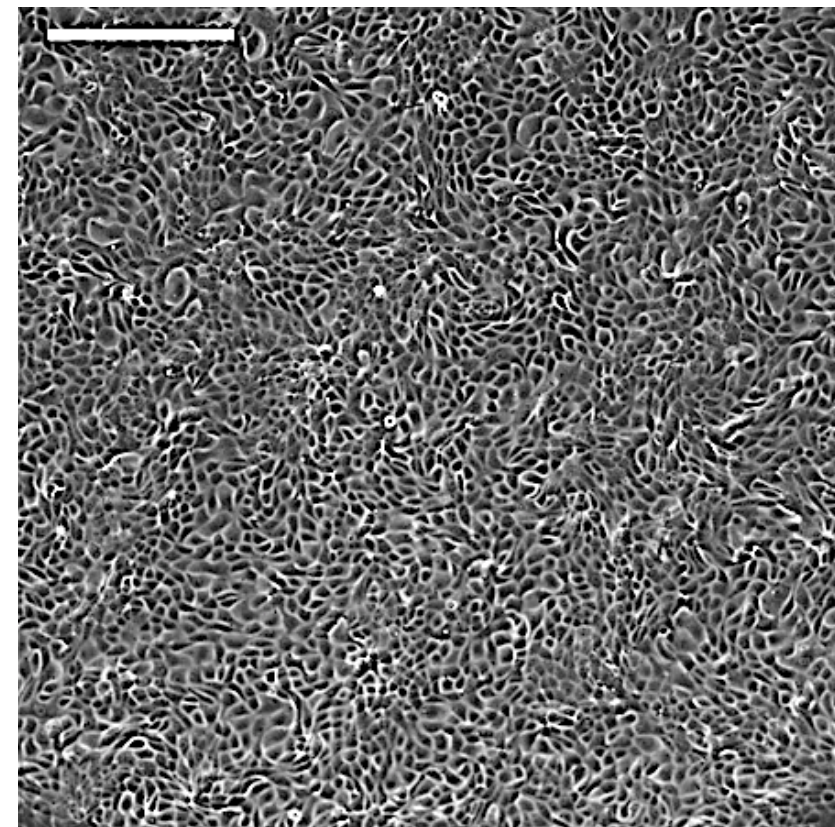
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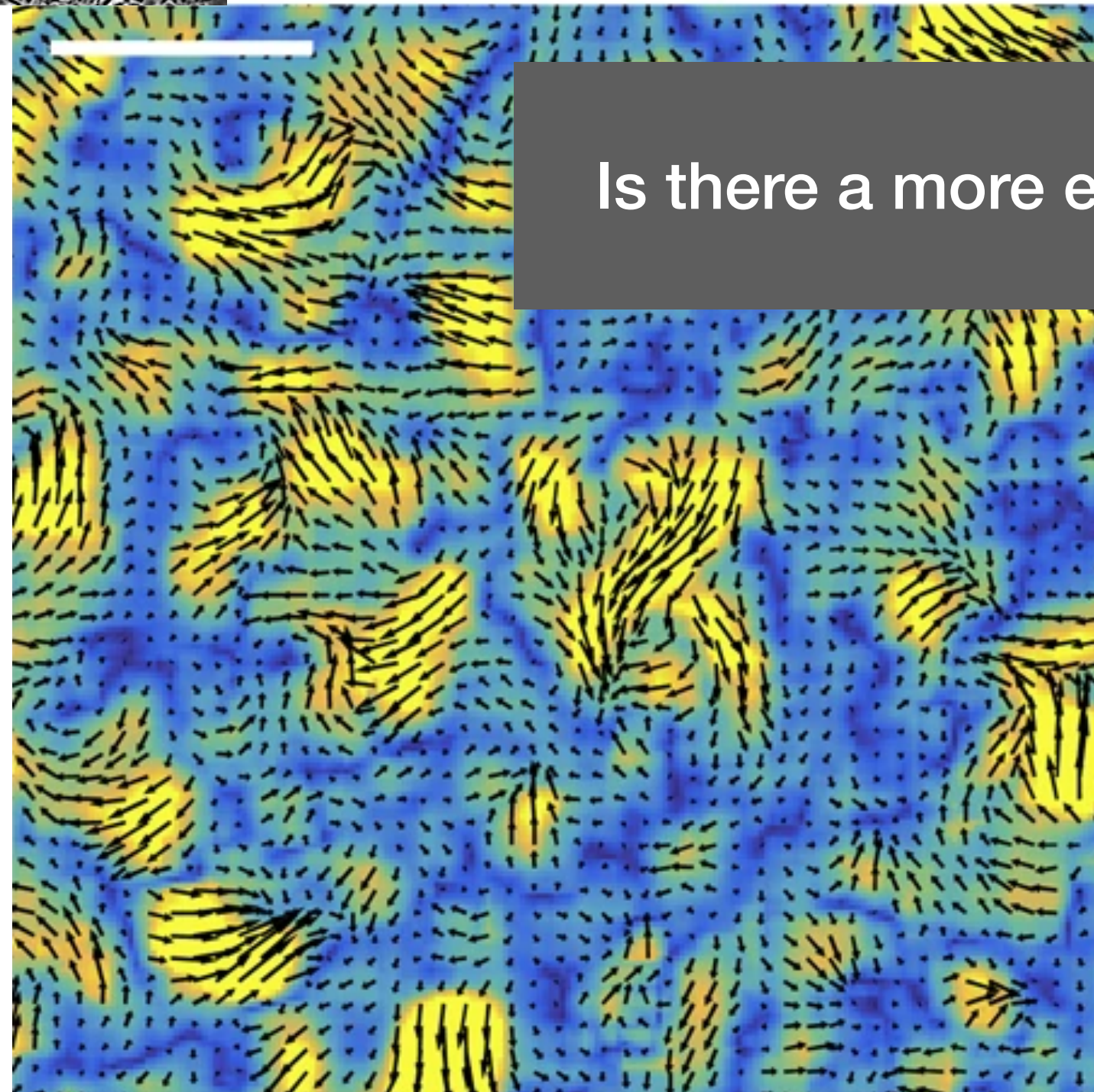
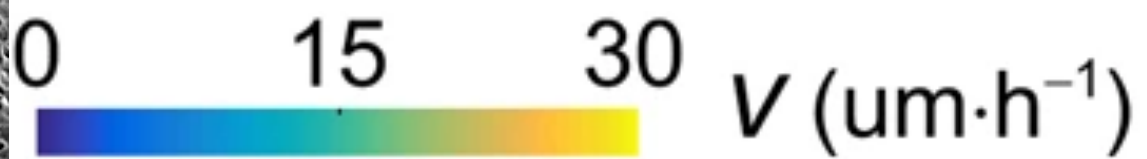
YE Keta, **J. Klamser**, RL Jack, L. Berthier, arXiv 2306.07172 (2023).



# Collective phenomena in active matter

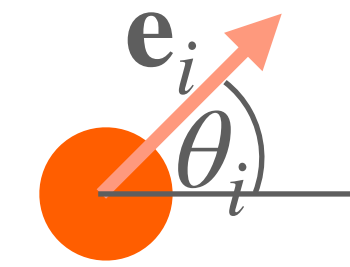


Complex chaotic flows  
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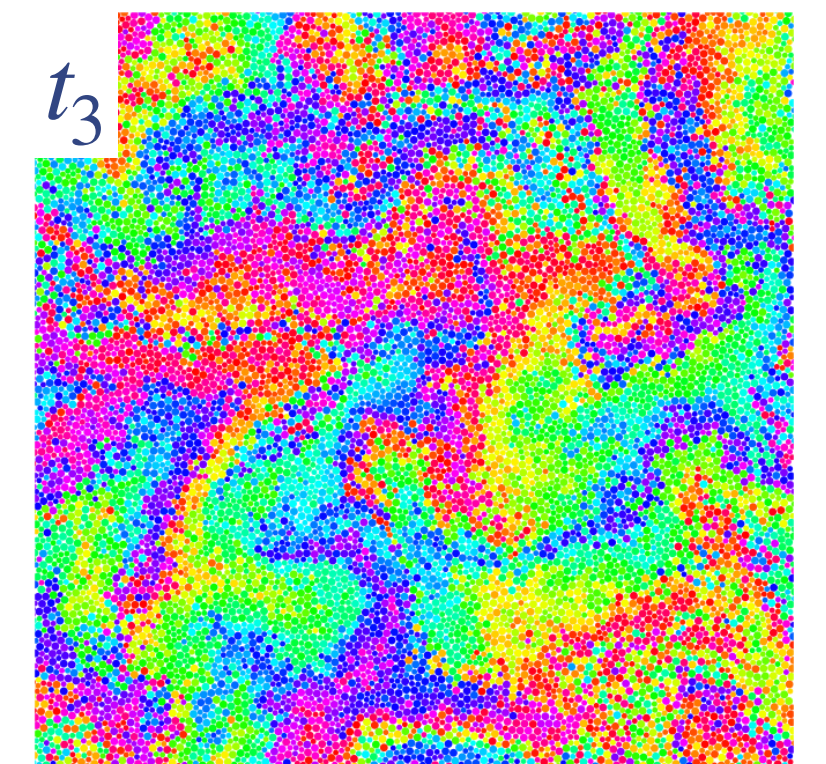
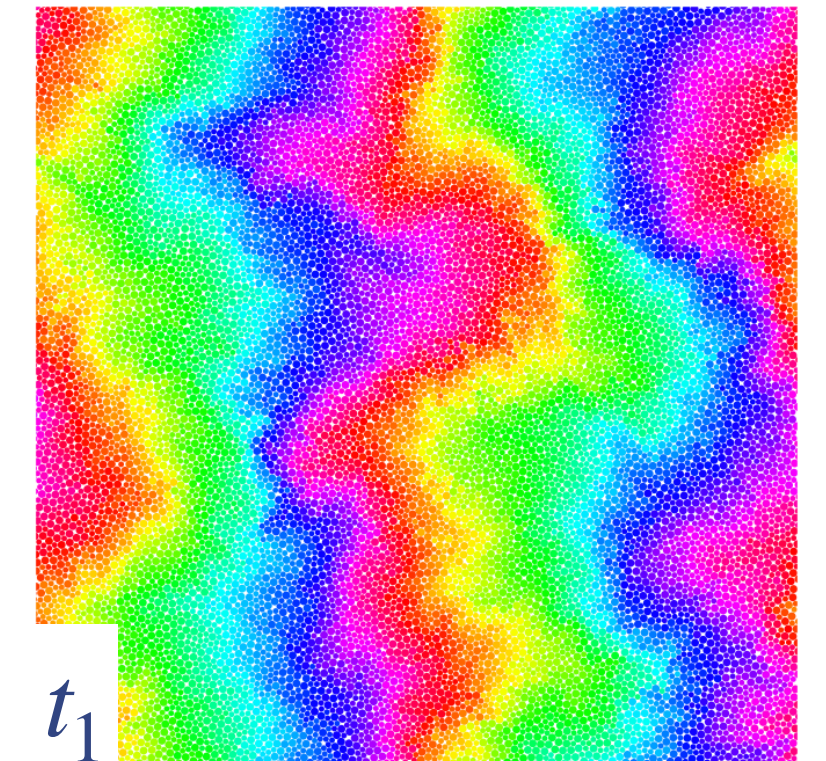
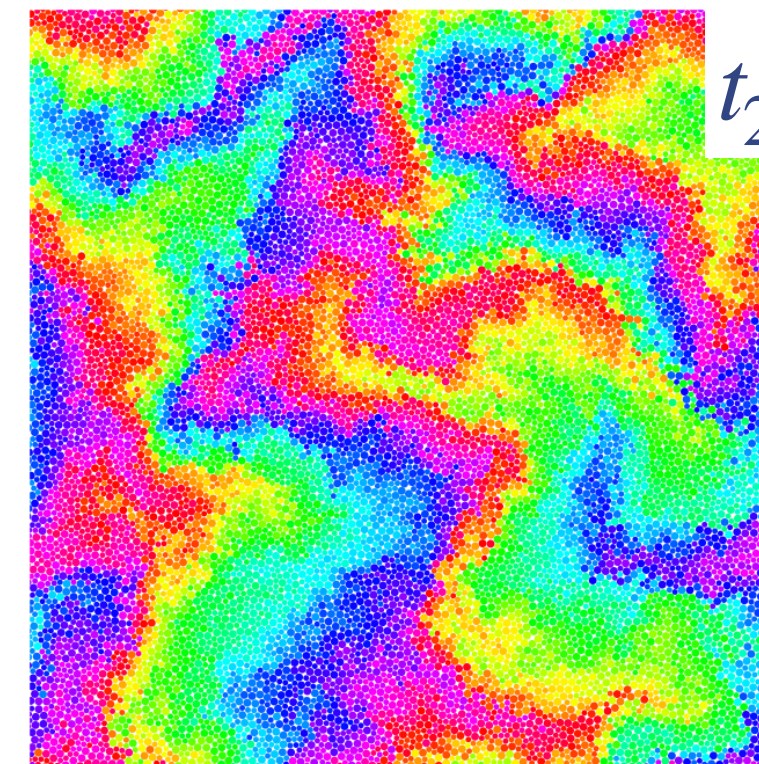
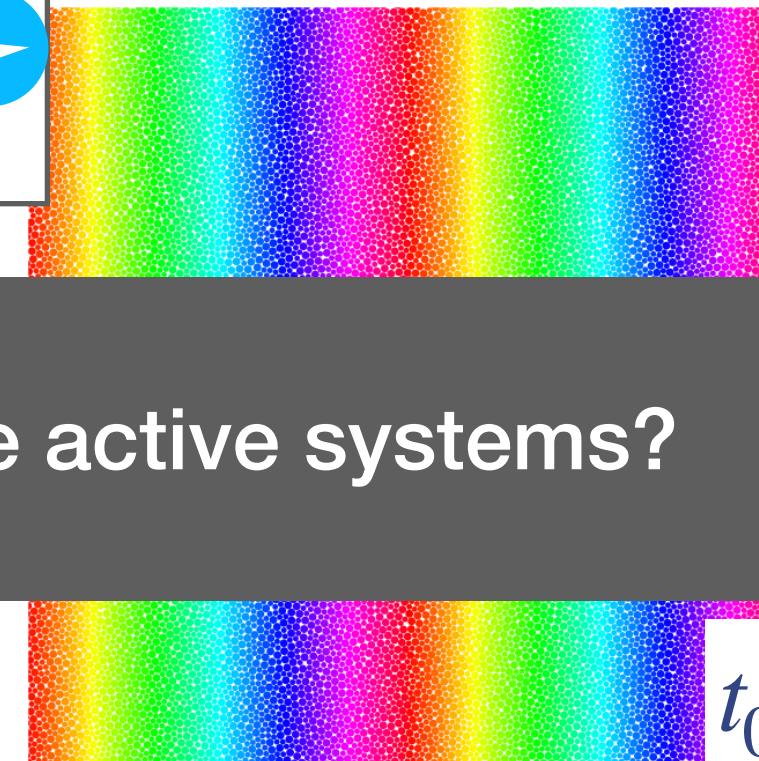
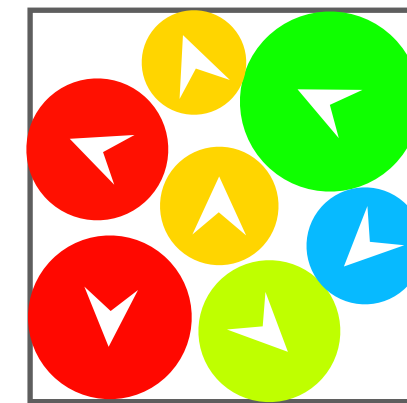
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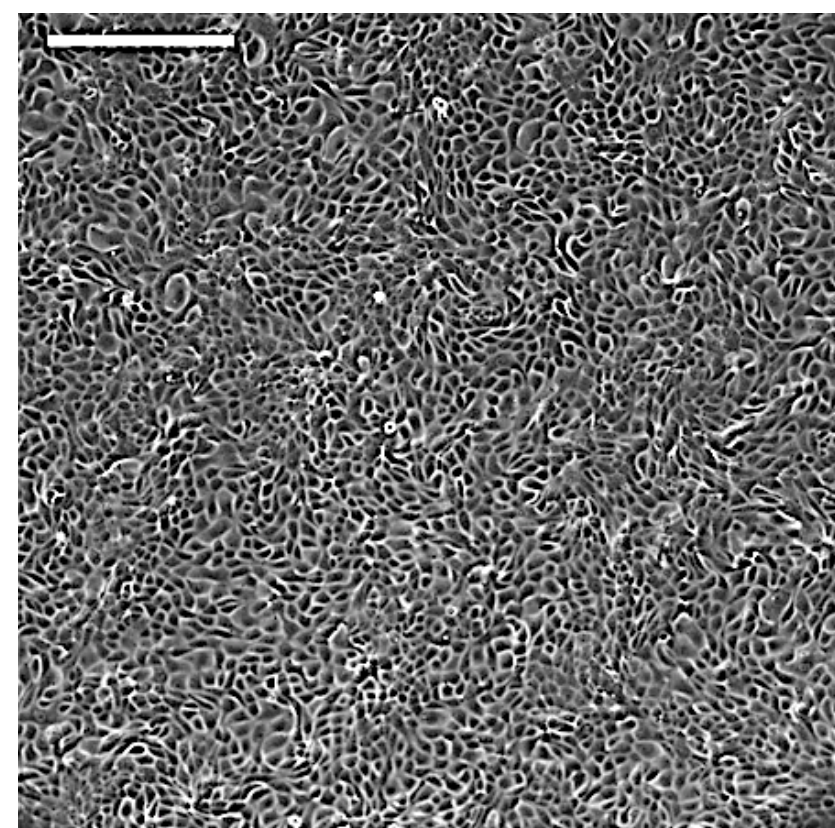


Is there a more efficient way to simulate active systems?

YE Keta, **J. Klamser**, RL Jack, L. Berthier, arXiv 2306.07172 (2023).



# Collective phenomena in active matter

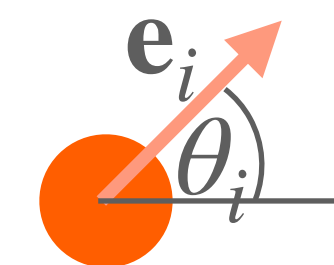


Complex chaotic flows  
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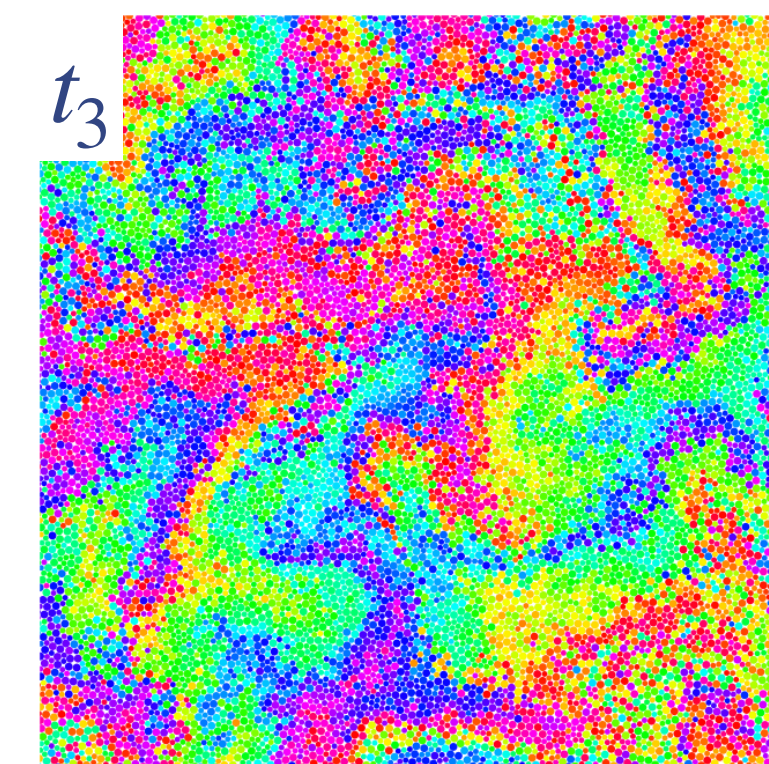
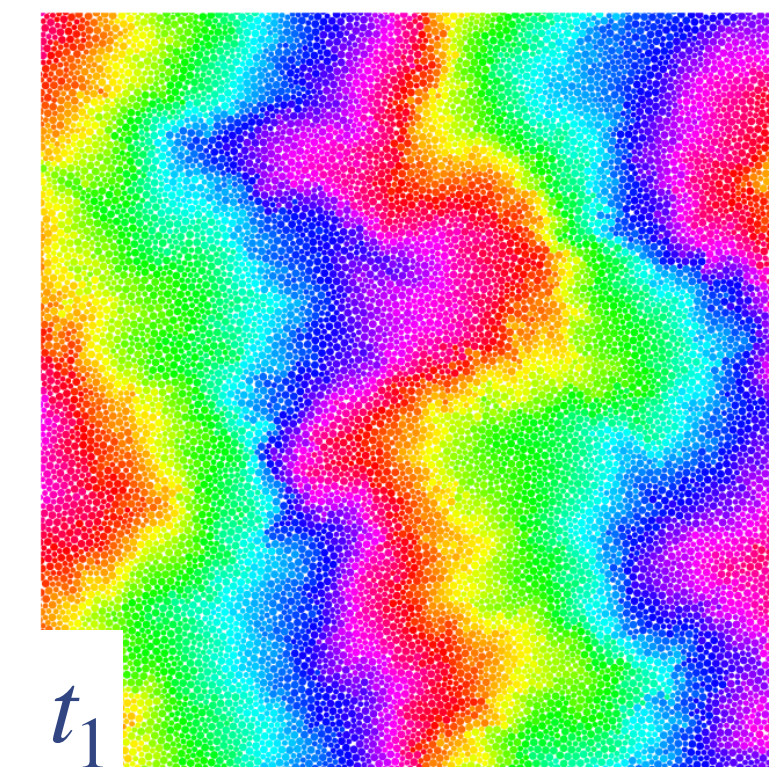
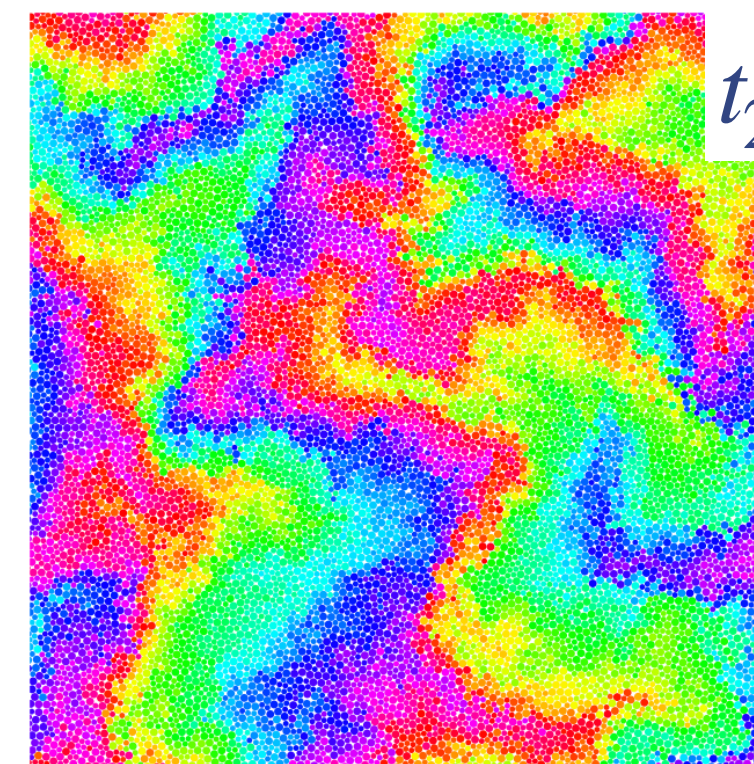
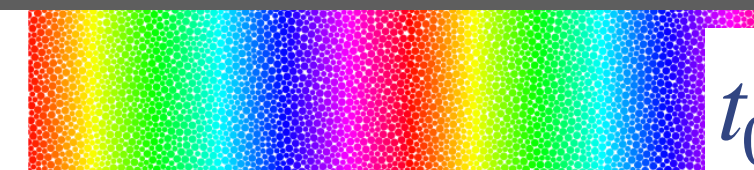
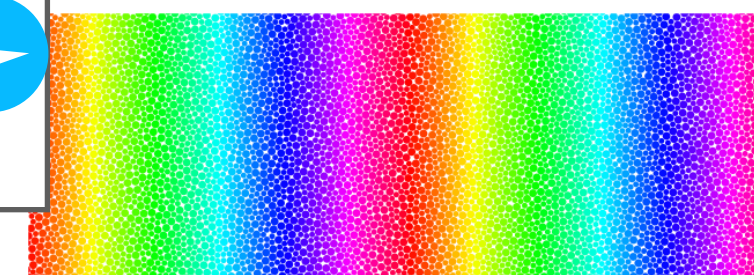
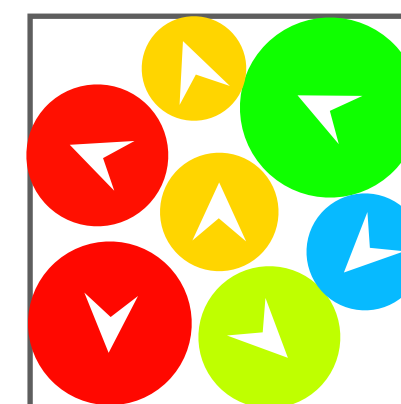
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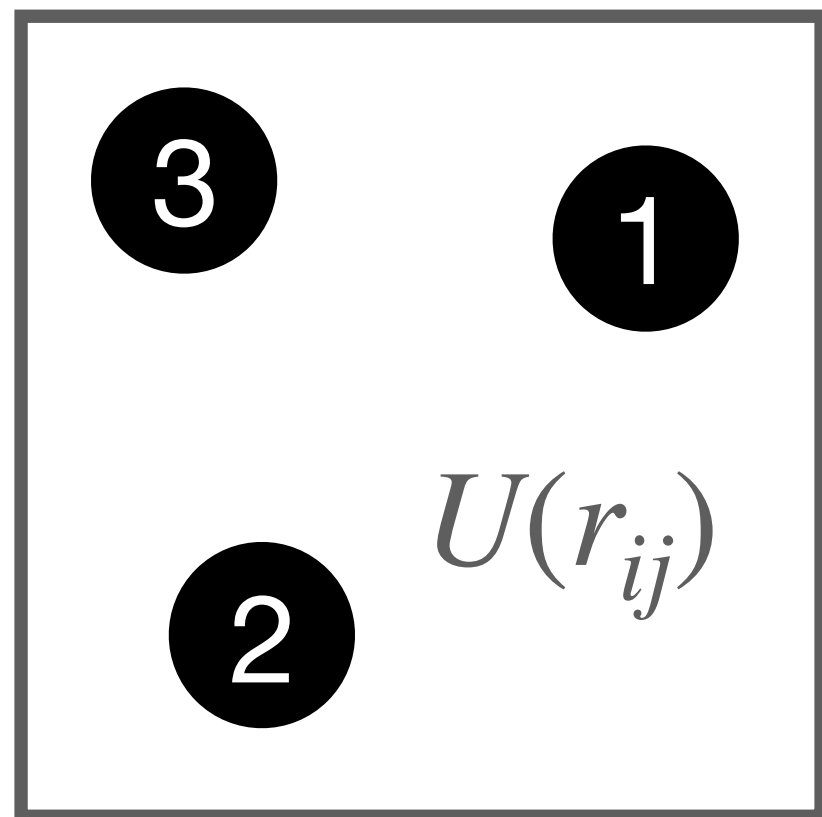
Is there a more efficient way to simulate active systems?

different

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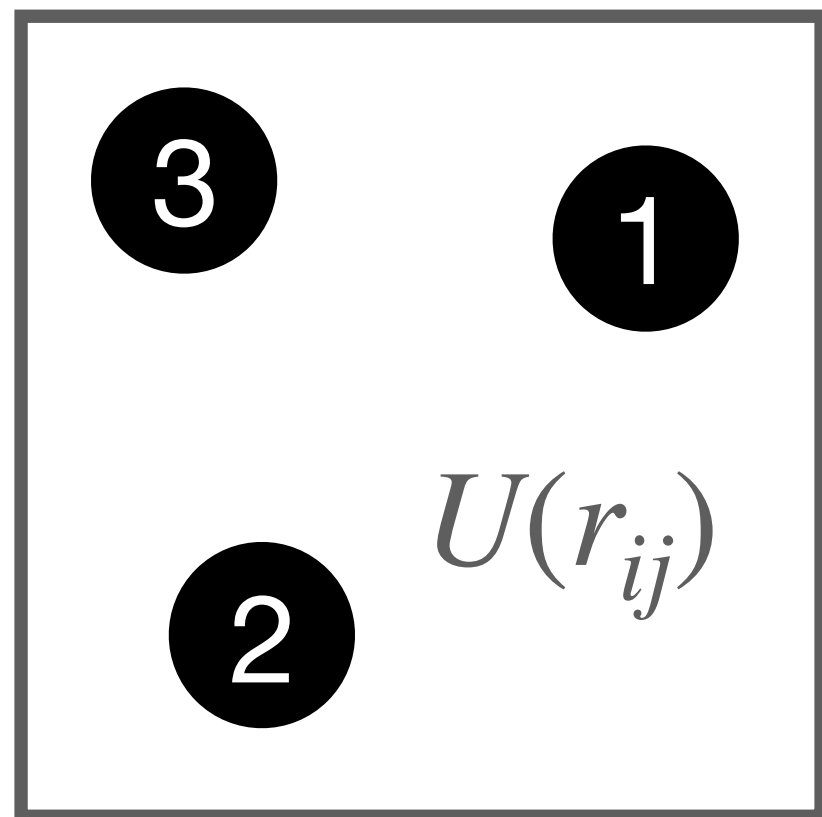


$$C = \{\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3\}$$





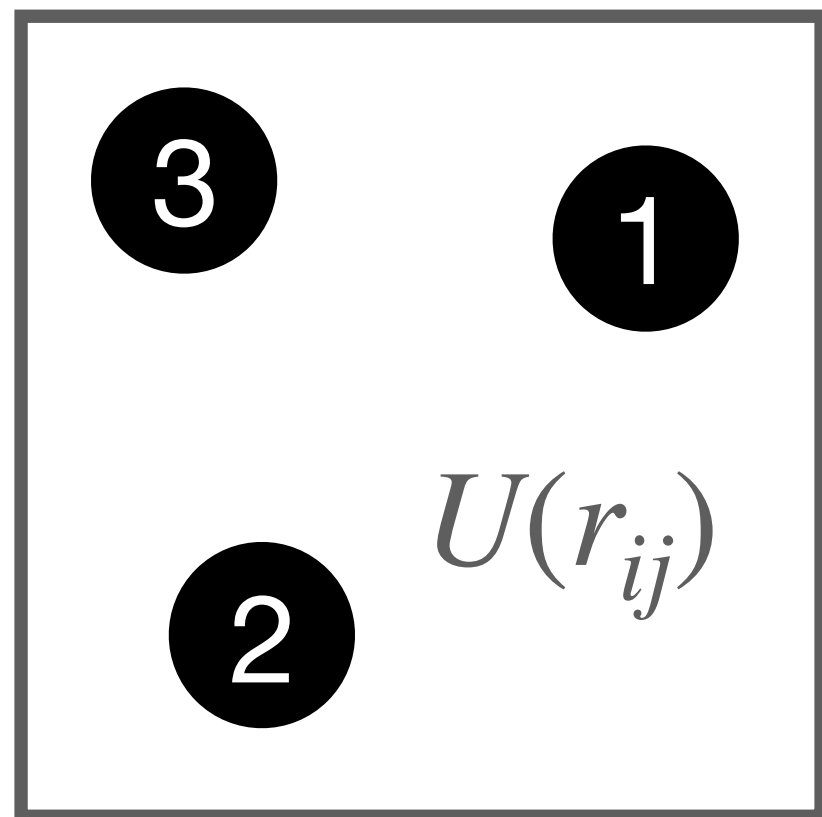
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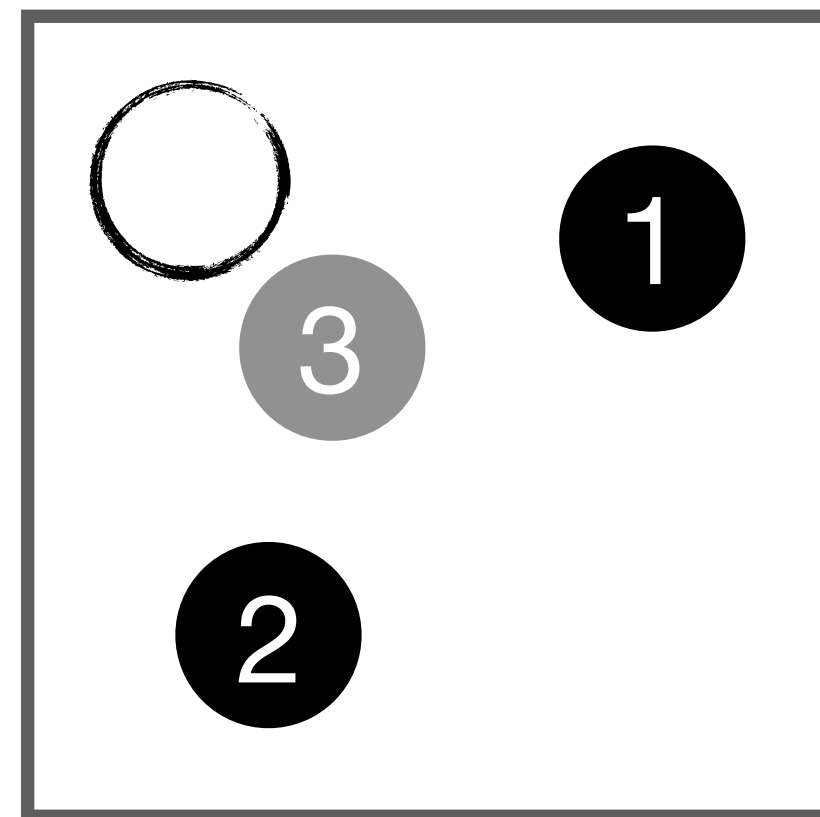
$$E(C) = \sum_i \sum_{j < i} U(r_{ij})$$



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$$C' = \{\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}'_3\}$$



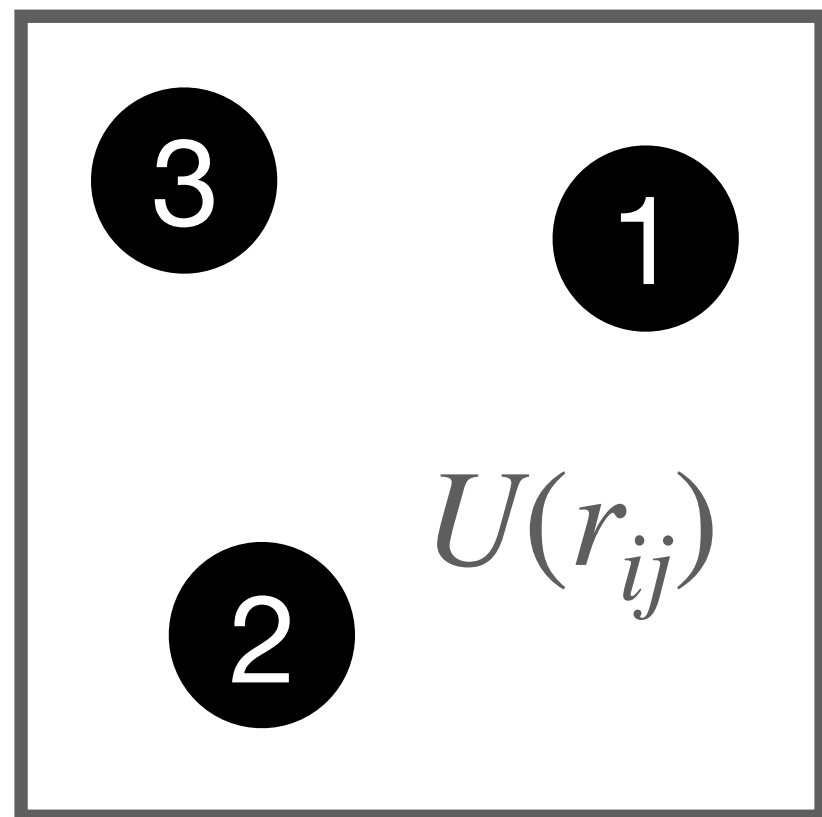
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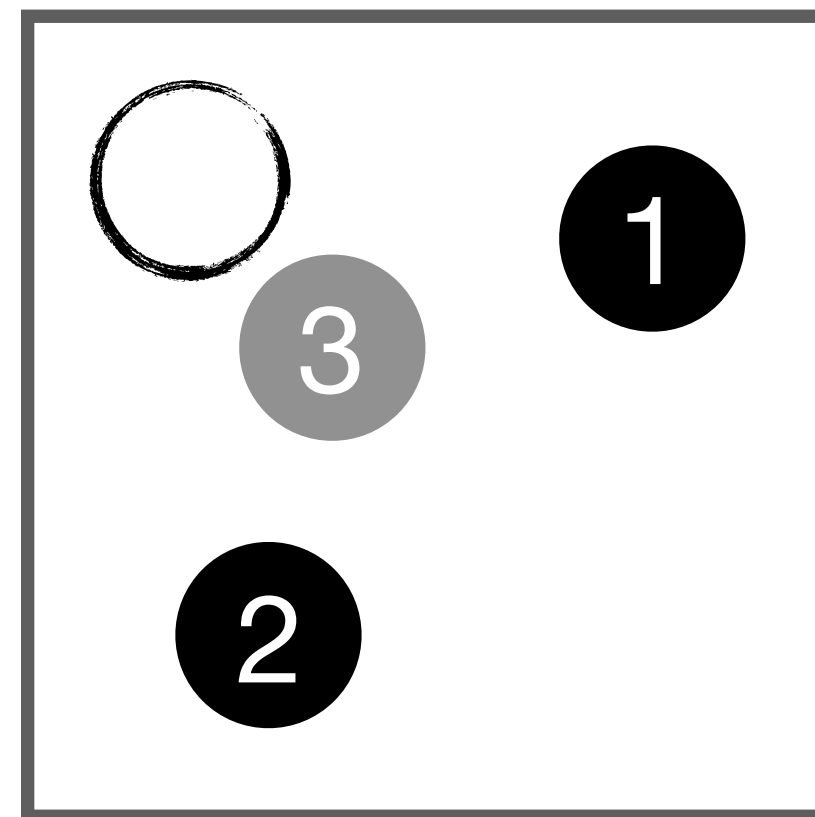


$$\mathcal{M}(C', C) = \min \left\{ 1, e^{-\beta[E(C') - E(C)]} \right\}$$

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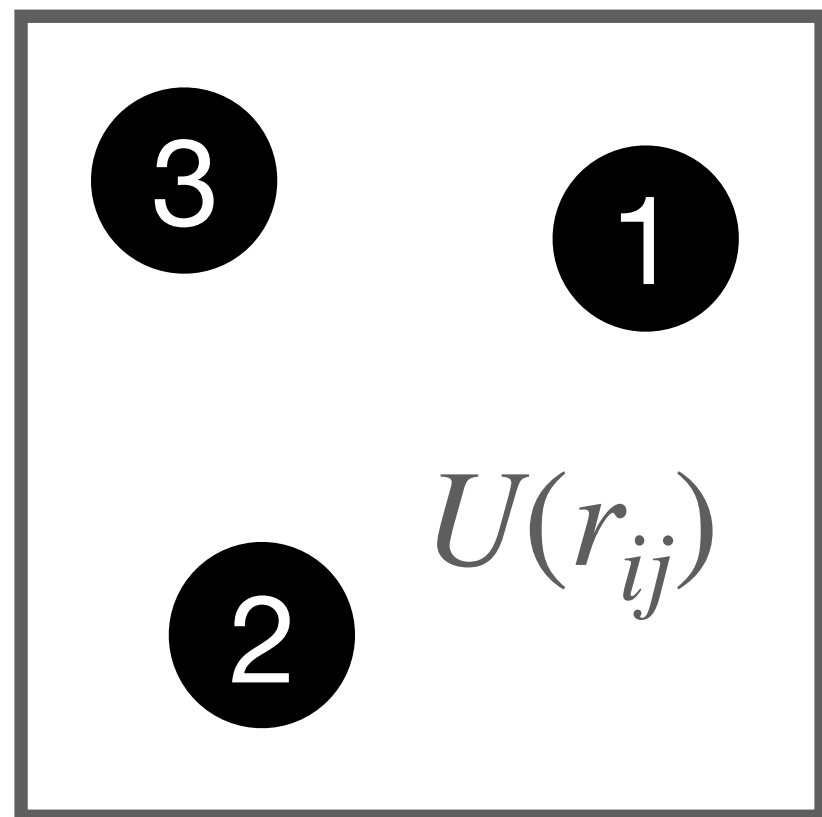
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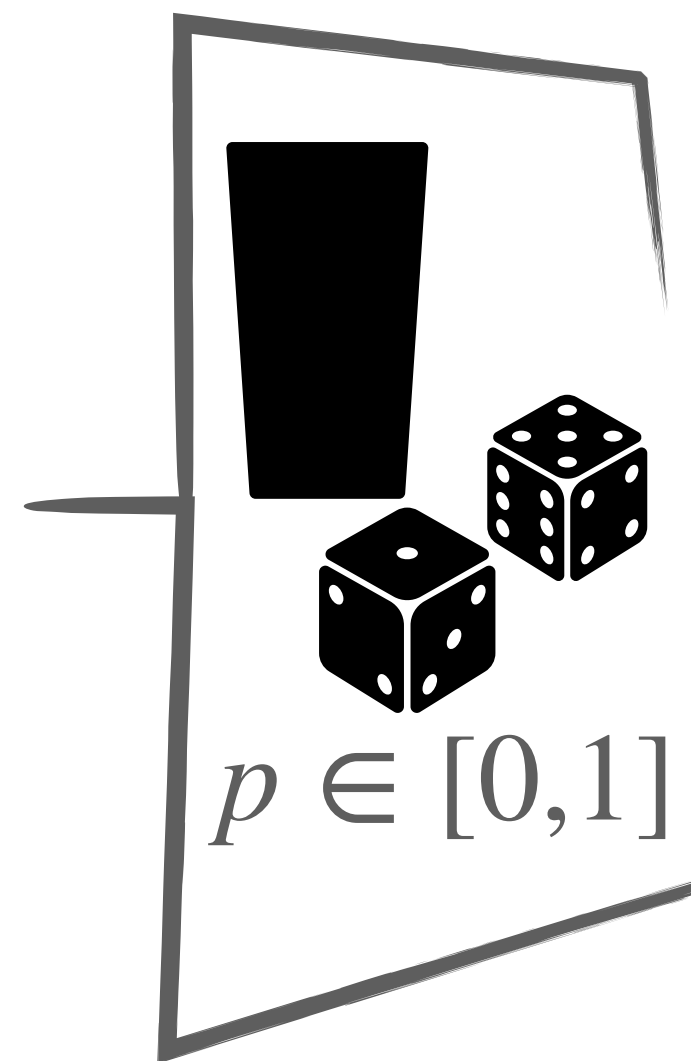
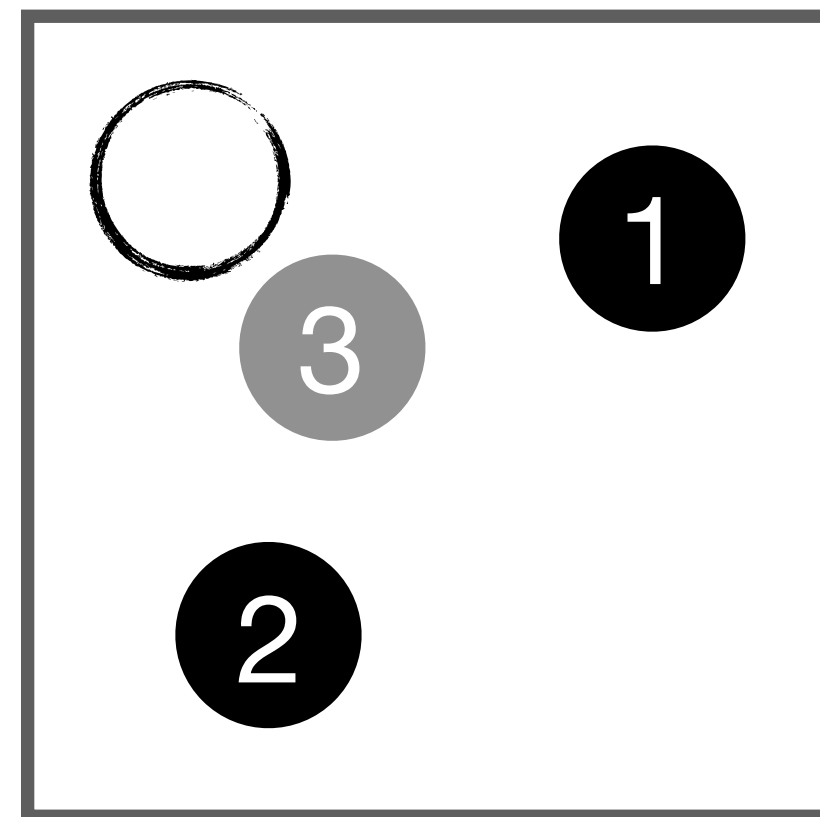


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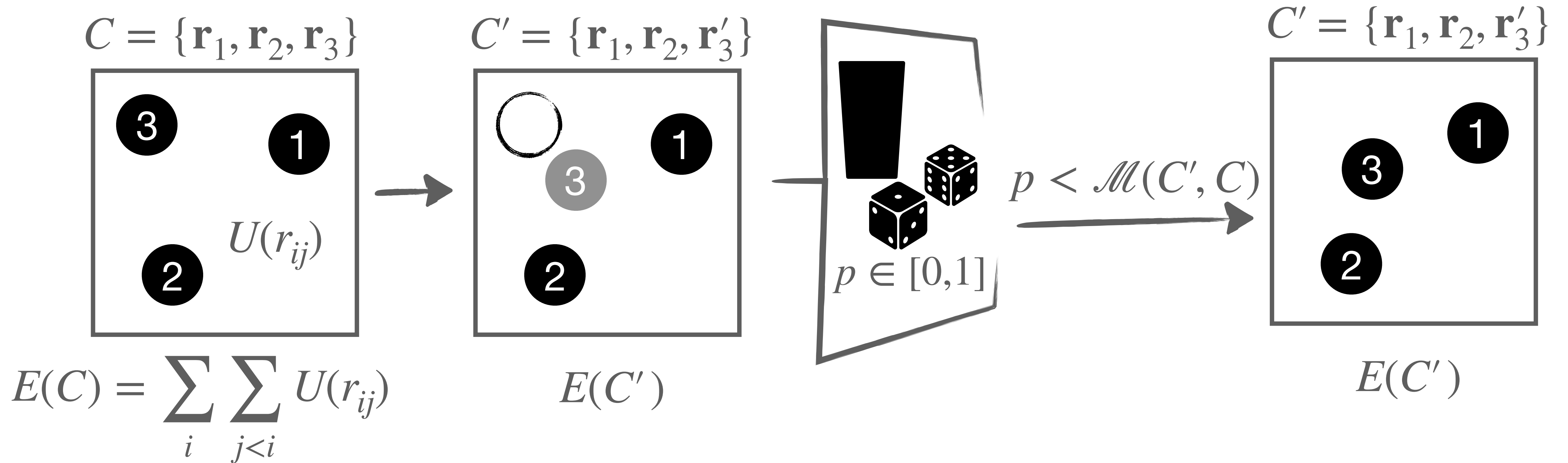


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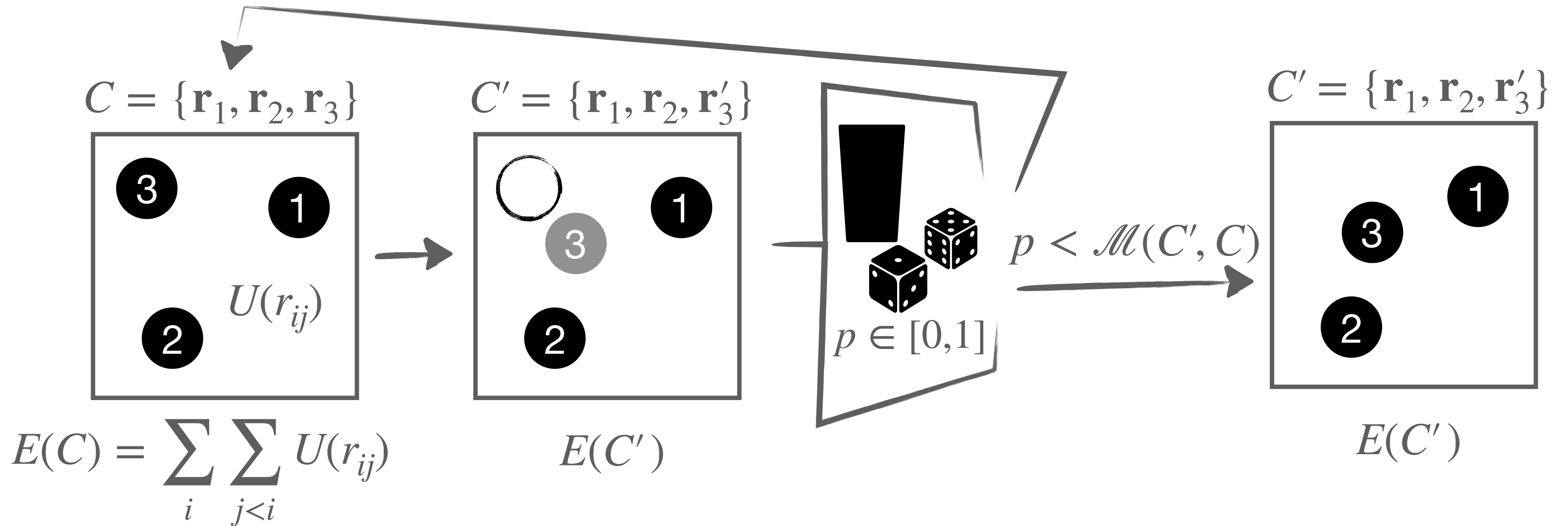


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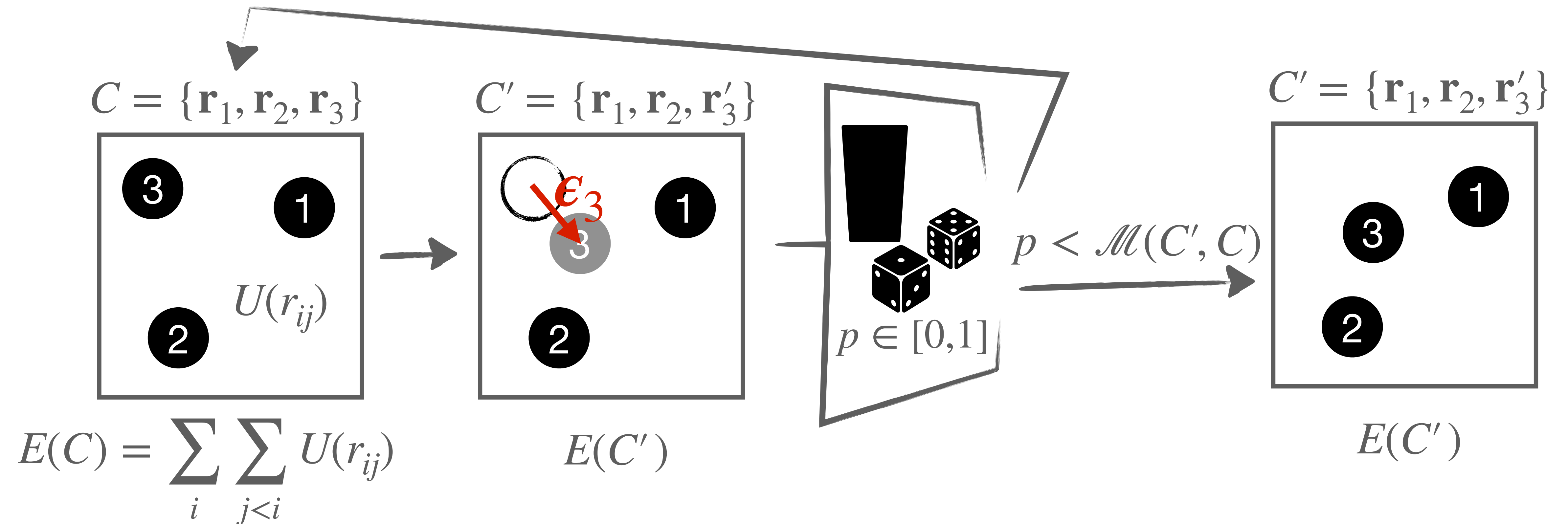


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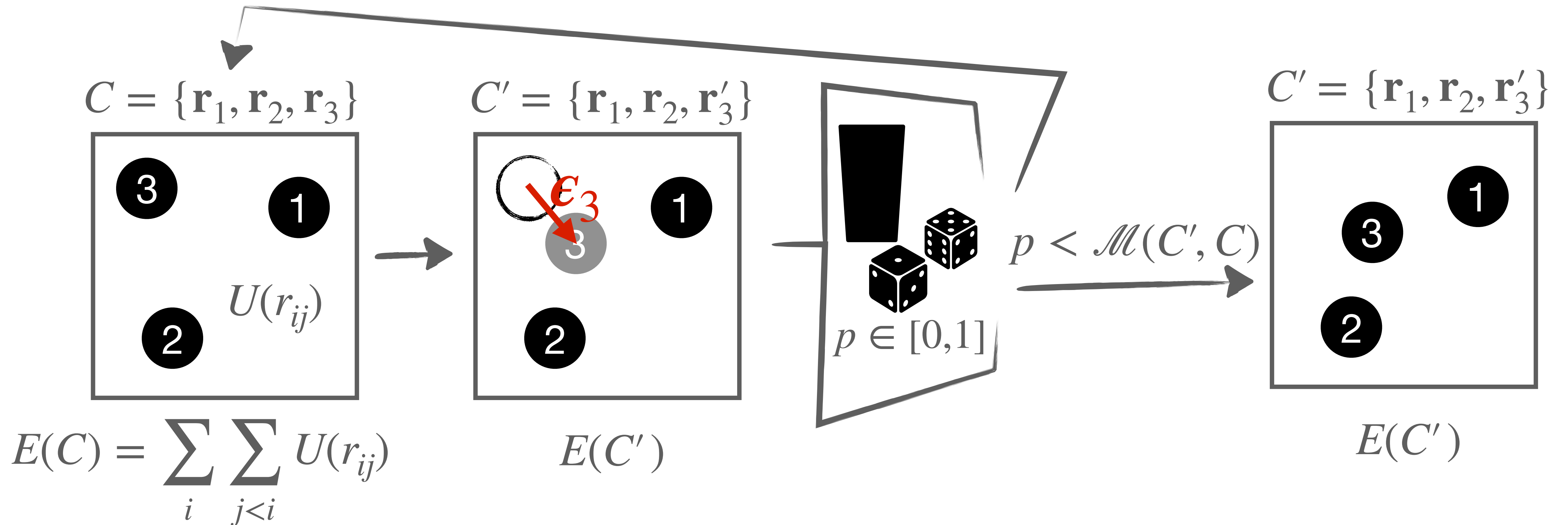
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Equilibrium Monte Carlo:  $G(\epsilon_i)$  = flat measure



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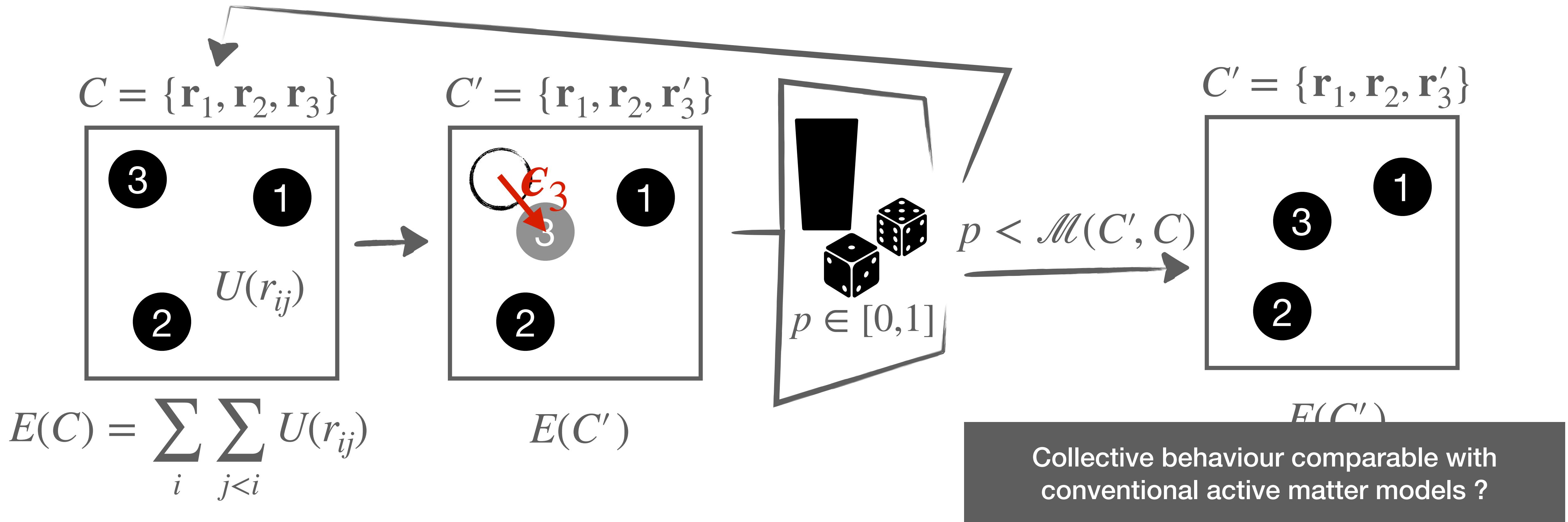


Equilibrium Monte Carlo:  $G(\epsilon_i)$  = flat measure

Active kinetic MC:  $g(\epsilon_i \rightarrow \epsilon'_i)$ , time correlated such that  $\epsilon'_i \simeq \epsilon_i$   
 $\lambda$  - persistence length



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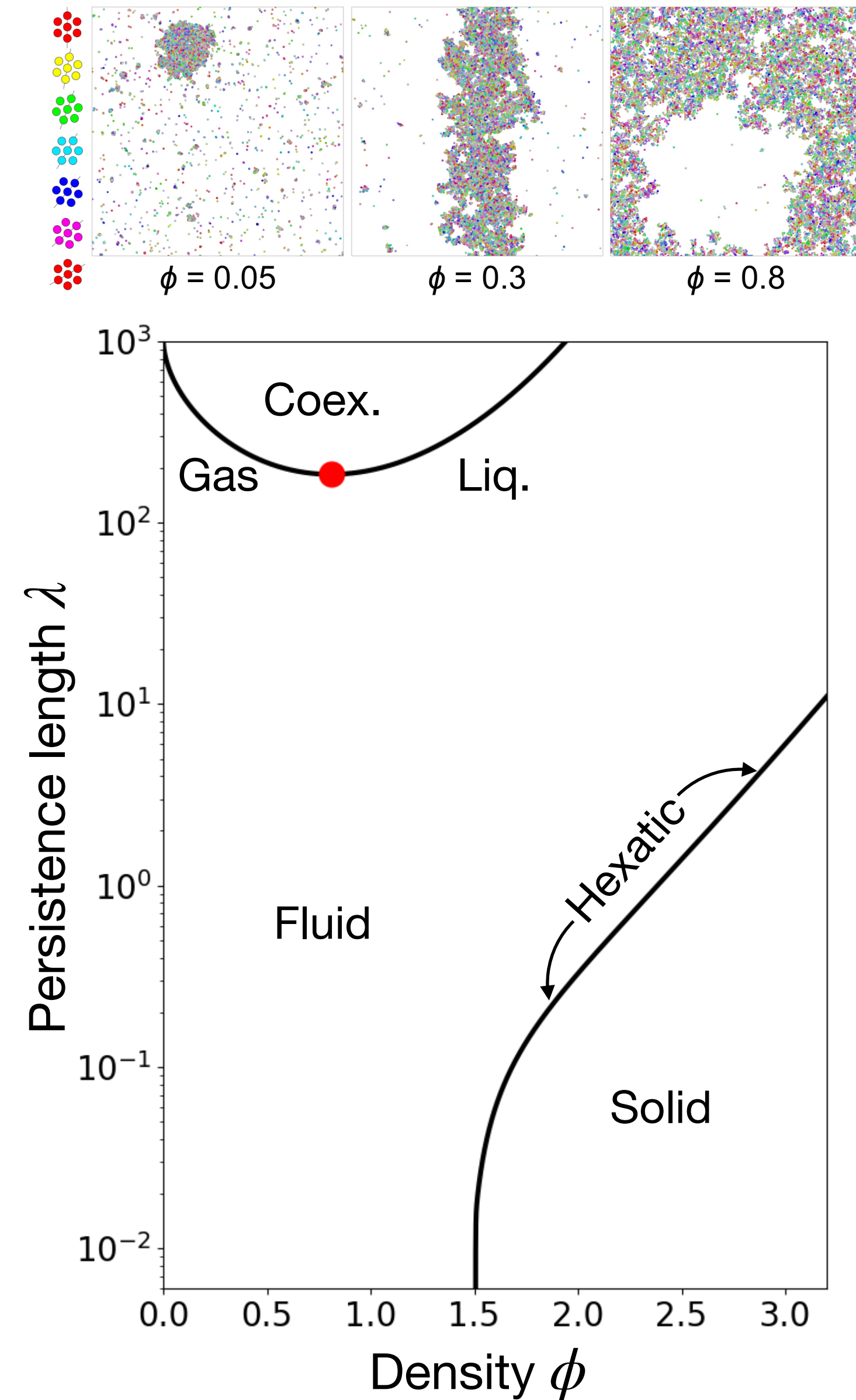
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Active kinetic MC:  $g(\epsilon_i \rightarrow \epsilon'_i)$ , time correlated such that  $\epsilon'_i \simeq \epsilon_i$   
 $\lambda$  - persistence length

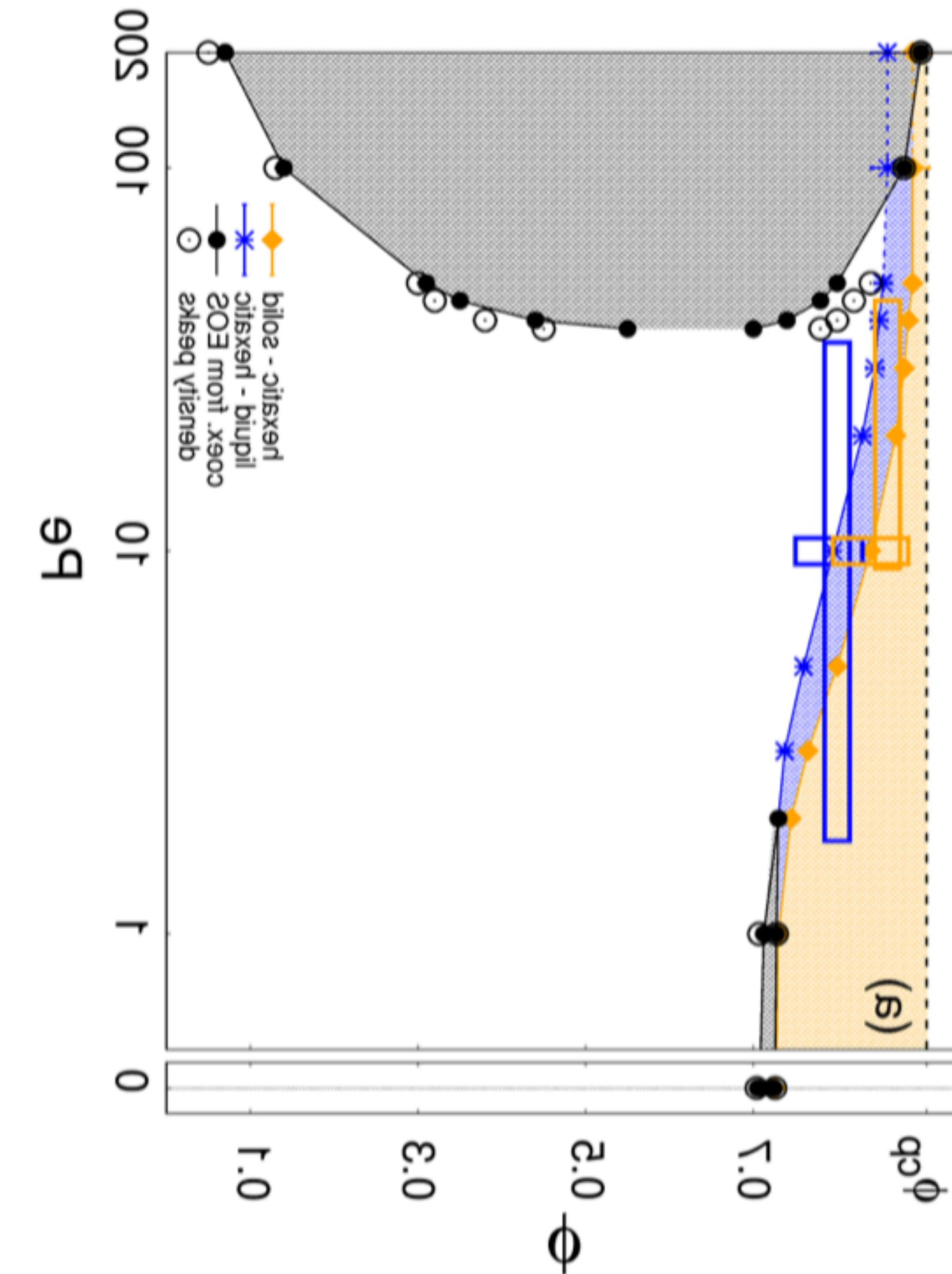
- D. Levis, L. Berthier, Phys. Rev. E (2014).
- L. Berthier, Phys Rev Lett (2014).
- **J. Klamser**, S. Kapfer, W. Krauth, Nat Commun (2018).
- **J. Klamser**, S. Kapfer, W. Krauth, J Chem Phys (2019).



# Active Kinetic Monte Carlo



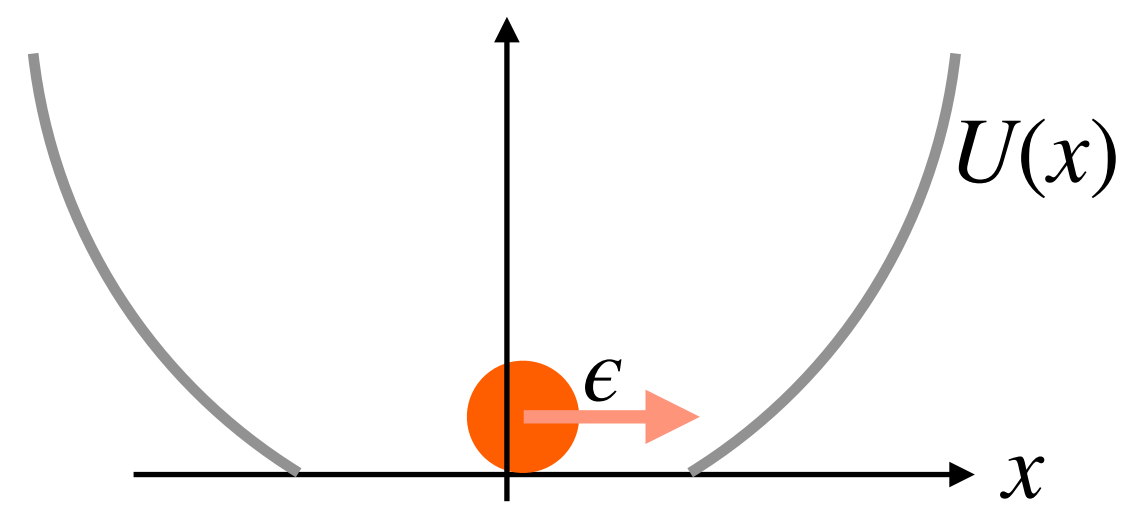
# Active Brownian Particles



P. Digregorio, D. Levis, A. Suma, L. F. Cugliandolo, G. Gonnella, I. Pagonabarraga, Phys Rev Lett (2018).

J. Klamser, S. Kapfer, W. Krauth, Nat Commun (2018).

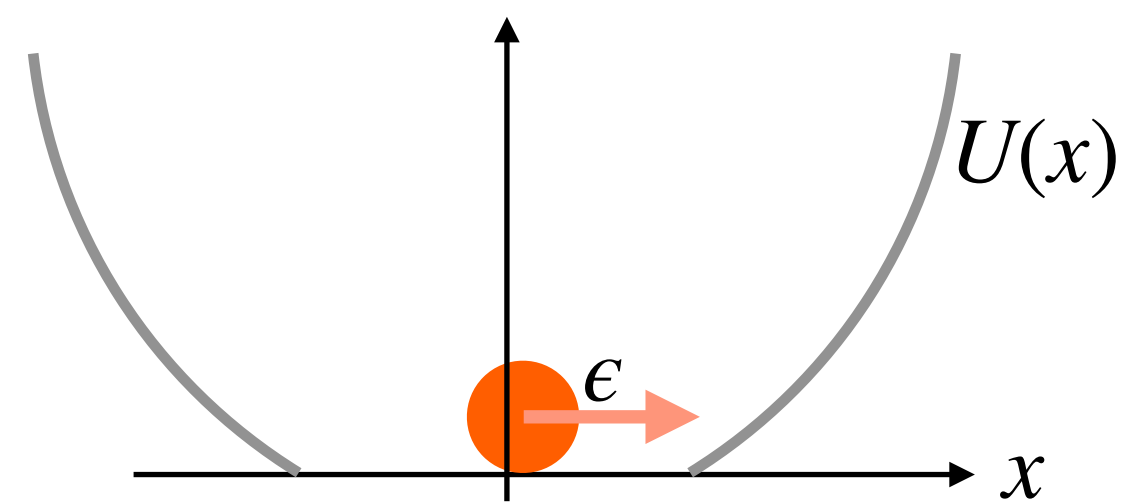




Continuous-time limit







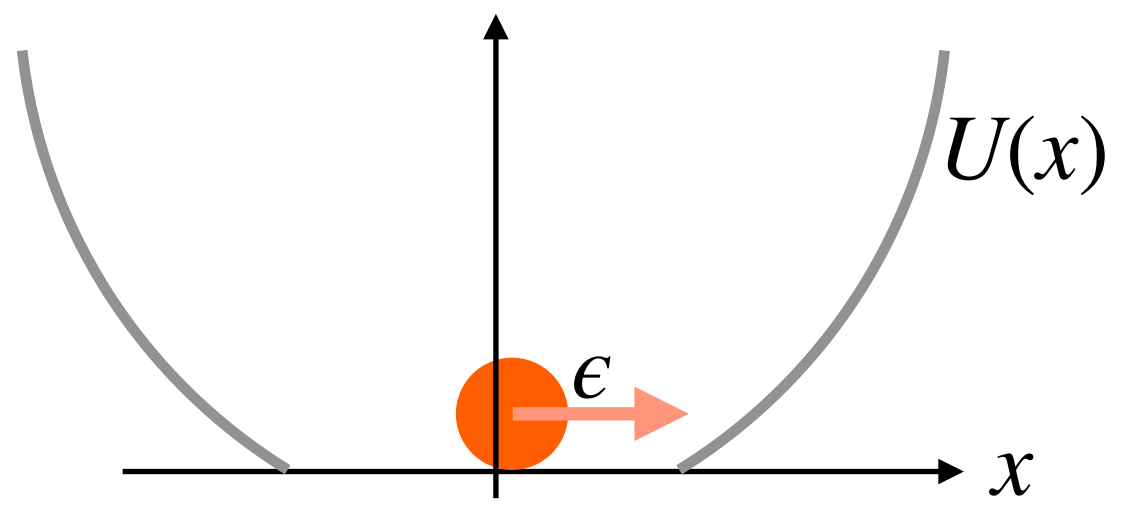
$$v = \frac{\epsilon}{dt} \quad t = ndt$$

Continuous-time limit

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Continuous-time limit

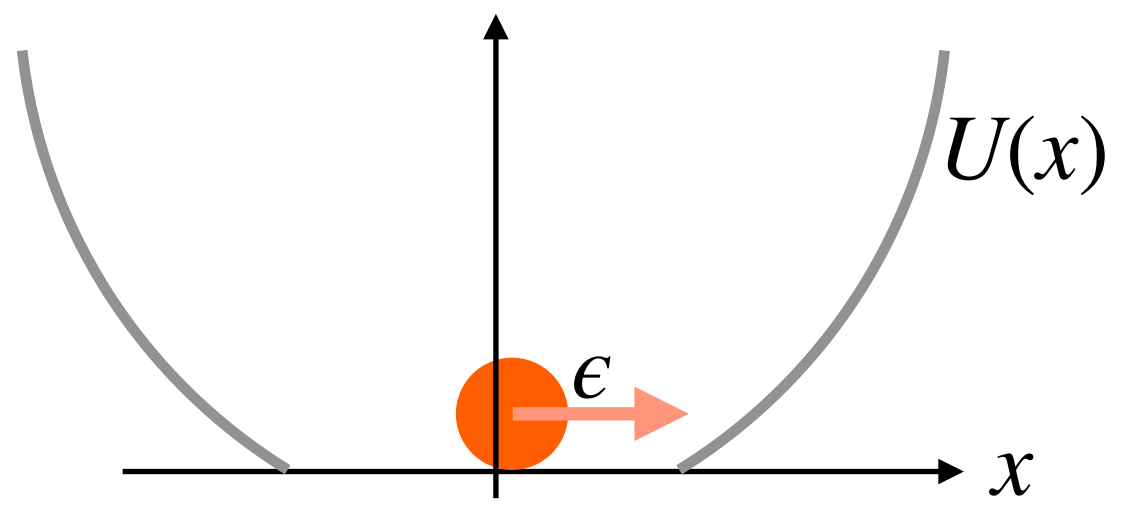
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$$v = \frac{\epsilon}{dt} \quad t = ndt$$


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$$dt \rightarrow 0 \quad \text{with} \quad v_0 = \text{const}$$





Continuous-time limit

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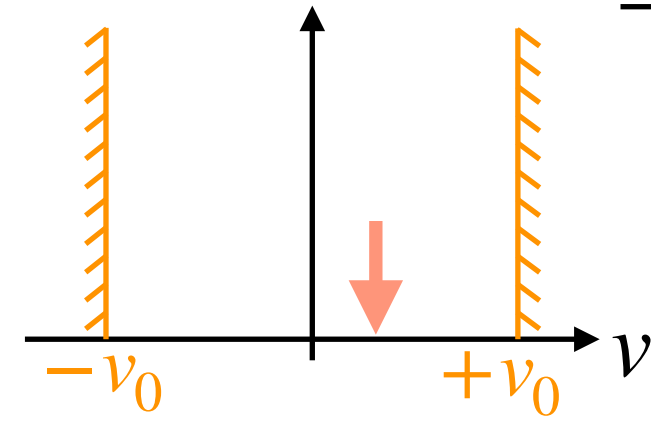
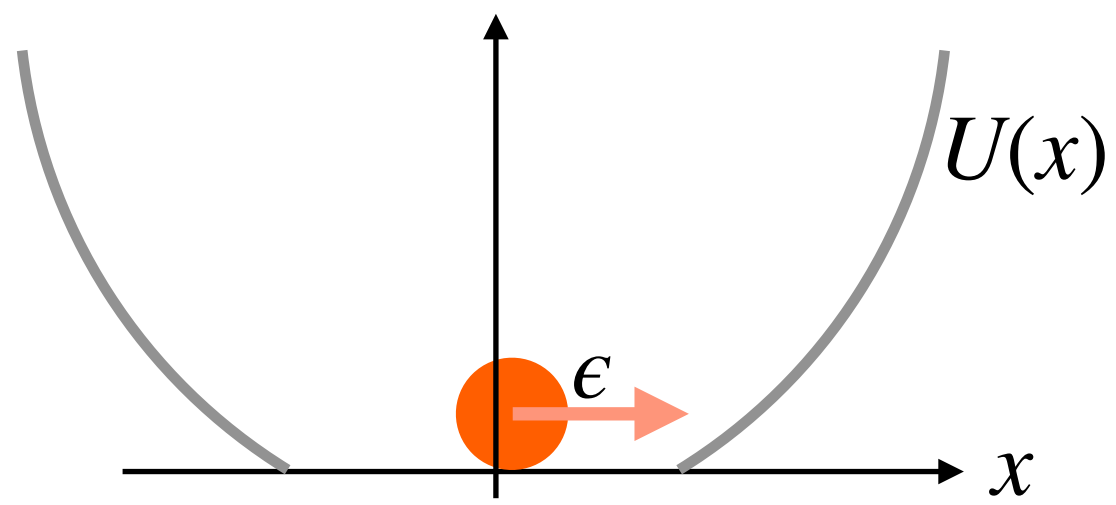
$$v = \frac{\epsilon}{dt} \quad t = ndt$$


---

$$dt \rightarrow 0 \quad \text{with} \quad v_0 = \text{const}$$

$$\tau = \text{const}$$





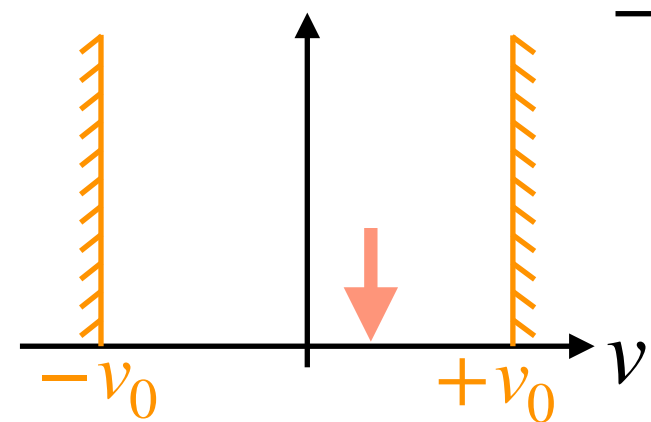
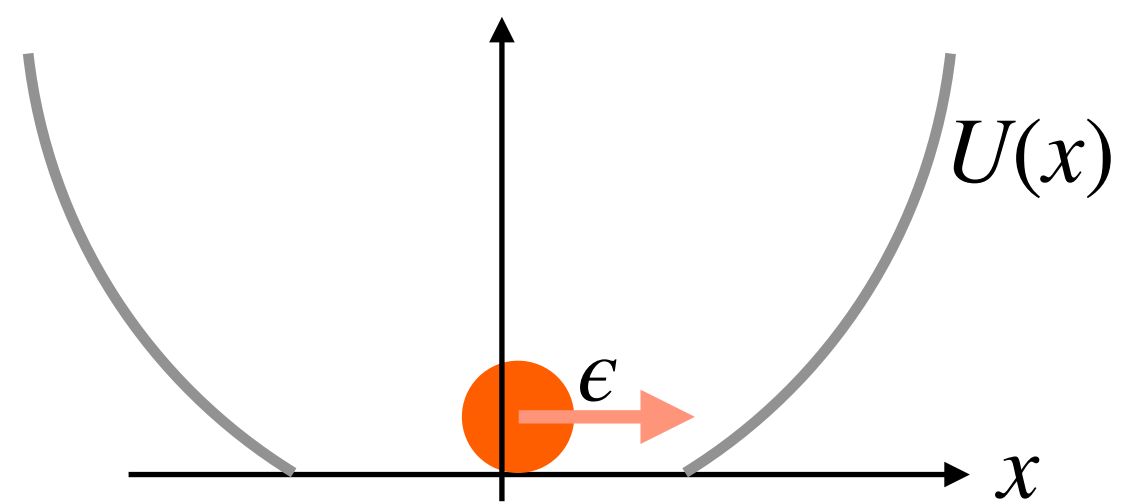
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Continuous-time limit

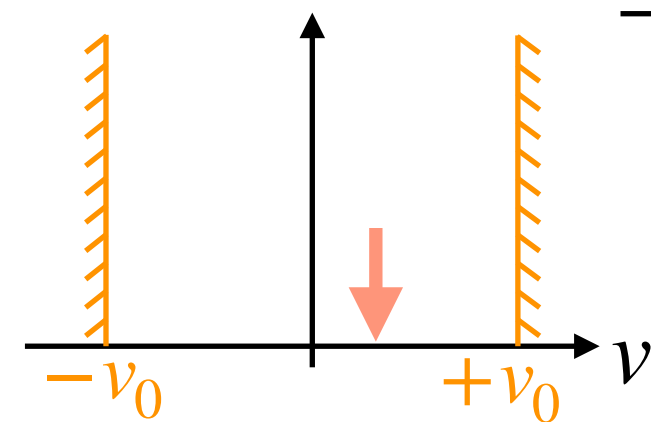
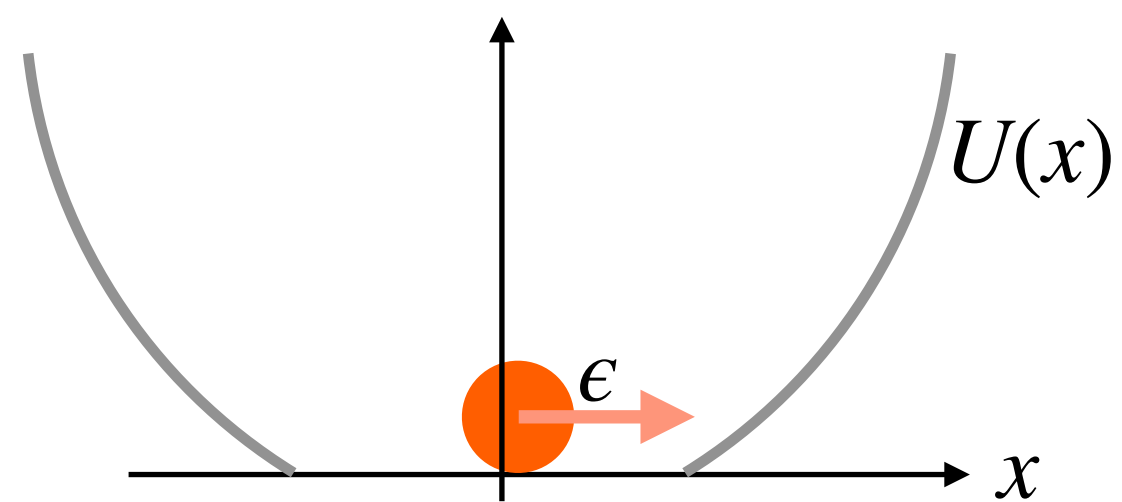
$$v = \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t}$$

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$$dt \rightarrow 0 \quad \text{with} \quad v_0 = \text{const}$$

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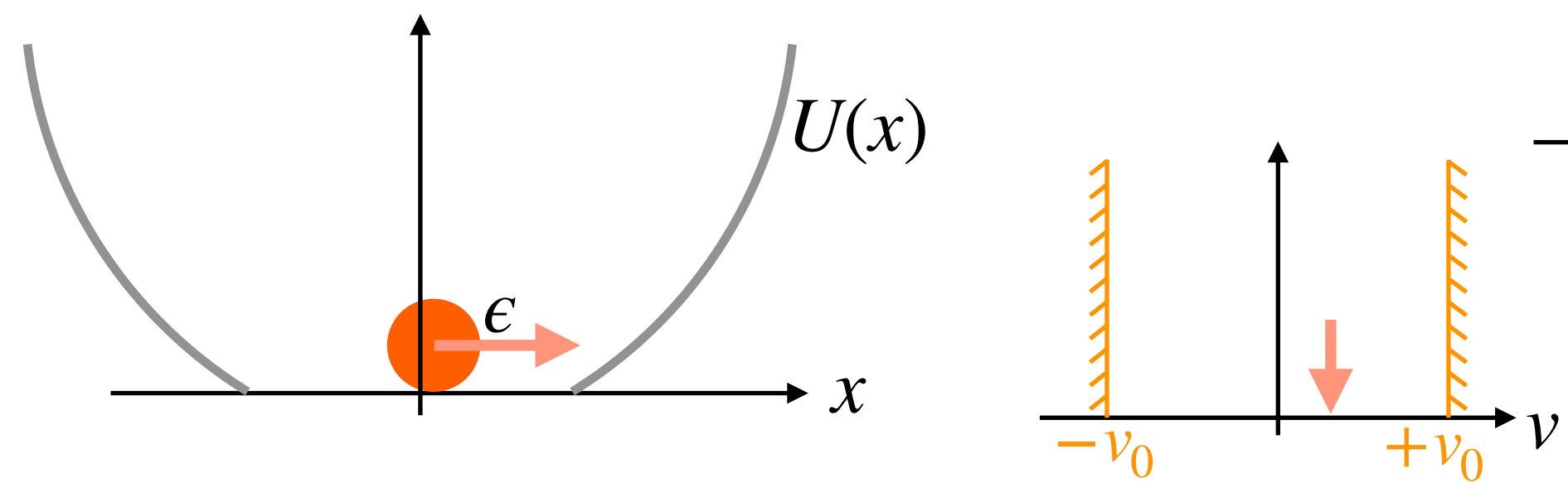
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Continuous-time limit

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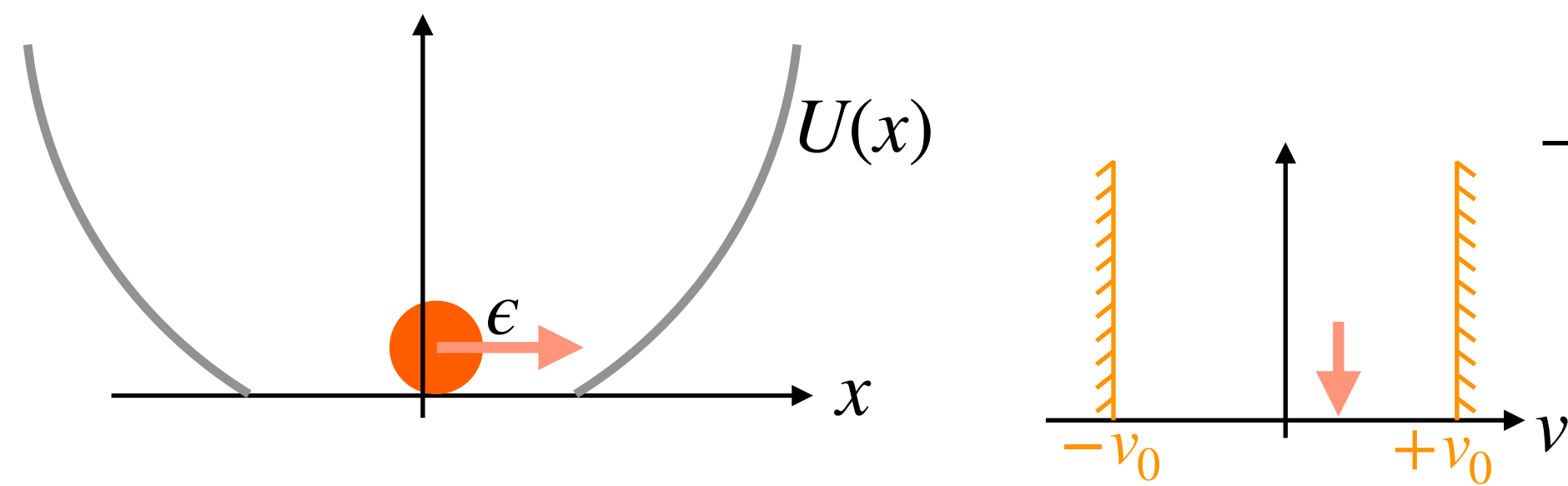

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$$dt \rightarrow 0 \quad \text{with} \quad v_0 = \text{const}$$

$$\tau = \text{const}$$

$$P_{n+1}(x, v) = \int_{-\infty}^{+\infty} dx' \int_{-v_0}^{v_0} dv' \, \textcolor{red}{g}(v | v') W(x | x', v' dt) P_n(x', v')$$





Continuous-time limit

$$v = \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t}$$


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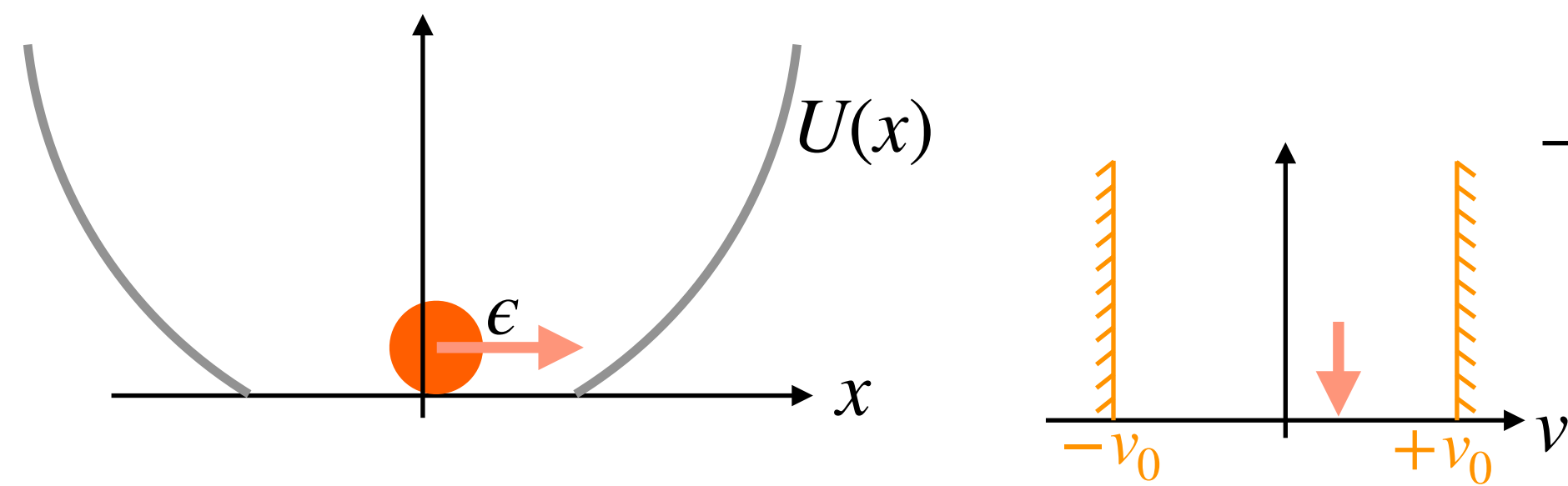

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Kramers-Moyal expansion

$$\frac{\partial P_t(x, v)}{\partial t} = D_v \frac{\partial^2}{\partial v^2} P_t(x, v) - \frac{\partial}{\partial x} v P_t(x, v)$$



Continuous-time limit

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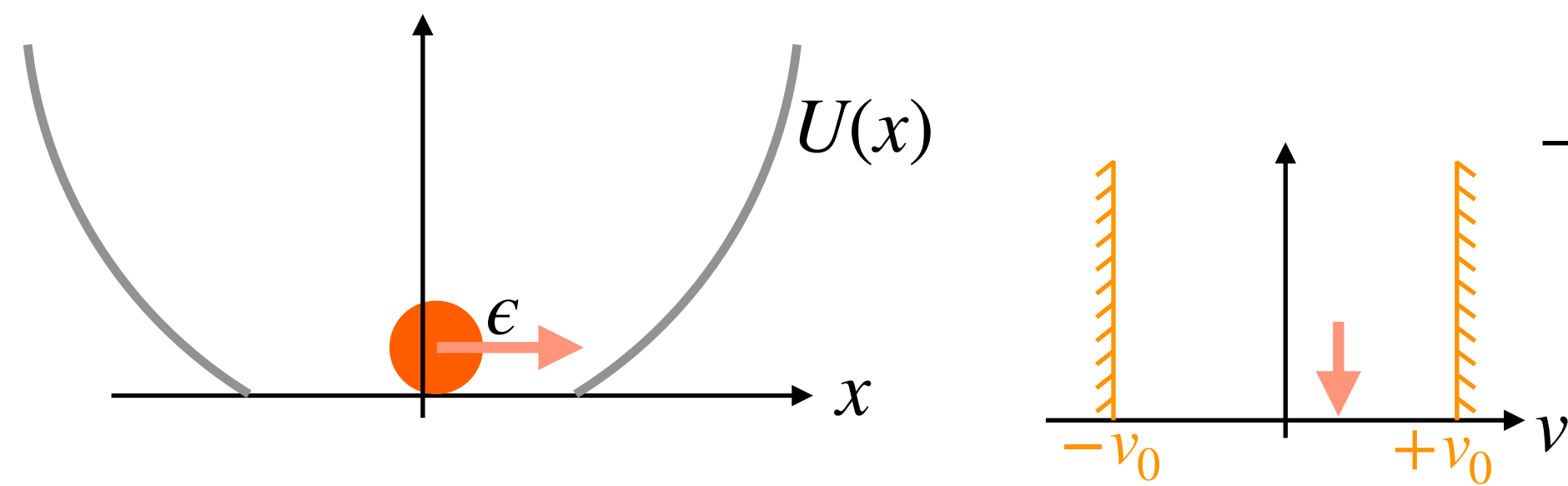
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$$\frac{dx}{dt} = v \quad ; \quad \frac{dv}{dt} = \sqrt{2D_v} \xi(t)$$





Continuous-time limit

$$v = \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t}$$


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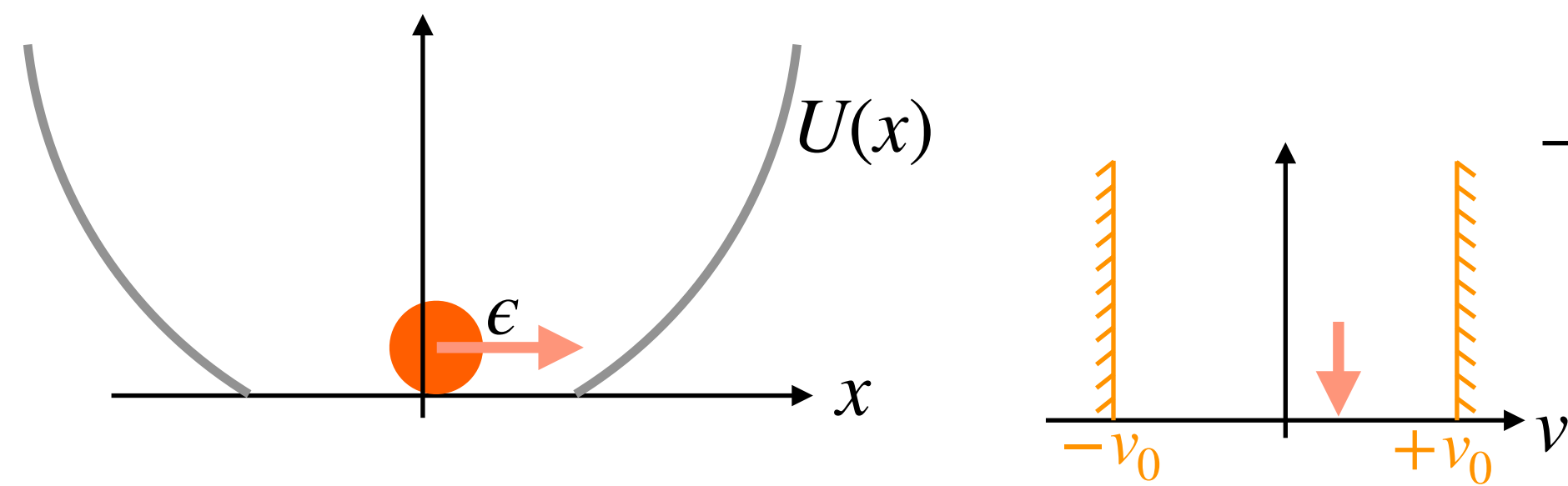
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$$\frac{dx}{dt} = v - \cancel{\mu U'(x)} ; \quad \frac{dv}{dt} = \sqrt{2D_v} \xi(t)$$



Continuous-time limit

$$v = \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t}$$


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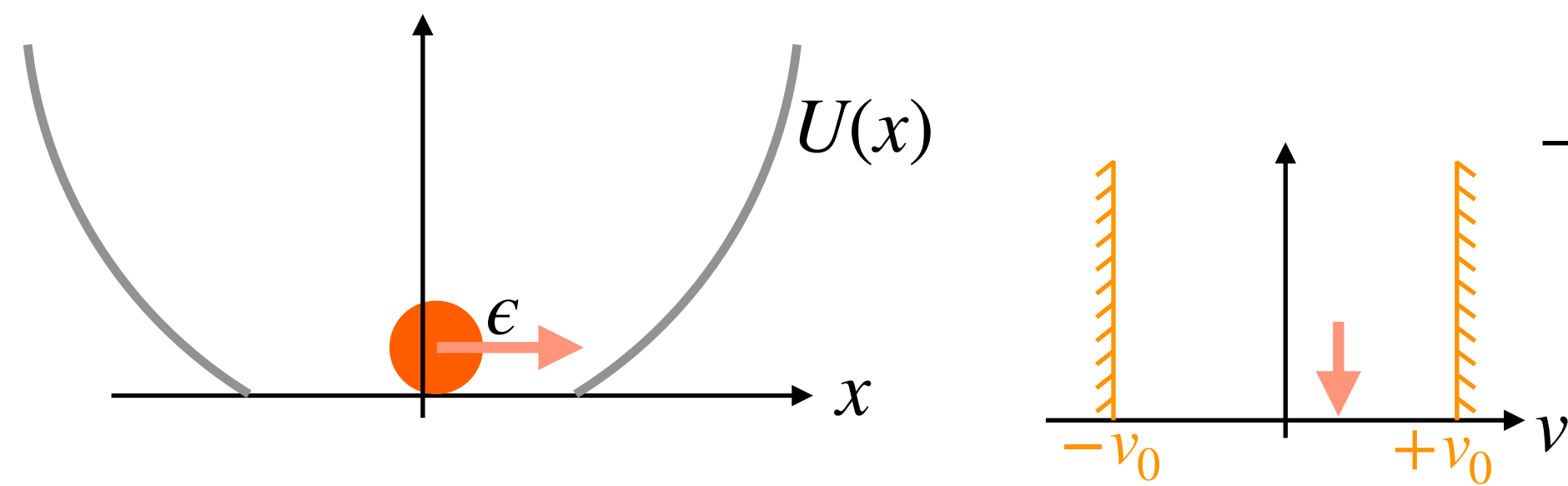
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$$\frac{dx}{dt} = v - \cancel{\mu U'(x)} ; \quad \frac{dv}{dt} = \sqrt{2D_v} \xi(t)$$

Correct active dynamics, but NO interactions.





Continuous-time limit

How do we also capture the confining potential?

$$v = \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t}$$


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$$dt \rightarrow 0 \quad \text{with} \quad v_0 = \text{const}$$

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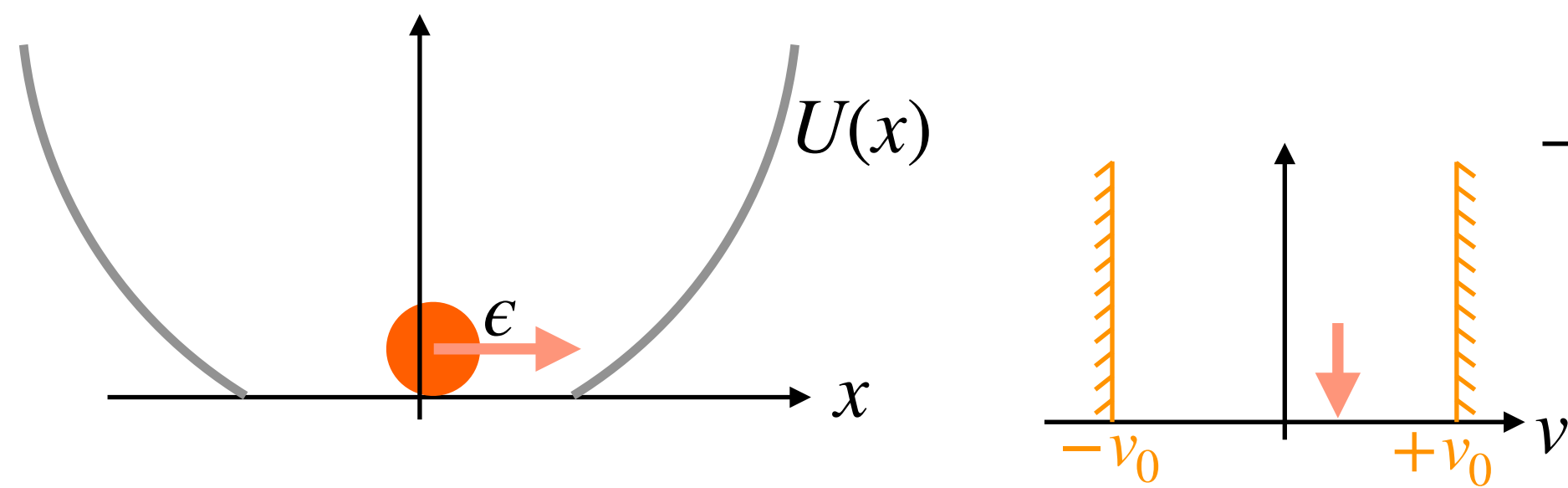
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Continuous-time limit

How do we also capture the confining potential?

$$\begin{aligned} v &= \frac{\epsilon}{dt} & t &= ndt & \langle \Delta v^2(t) \rangle &= \sqrt{2D_v t} \\ \hline dt &\rightarrow 0 & \text{with } v_0 &= \text{const} \\ & & \tau &= \text{const} \end{aligned}$$

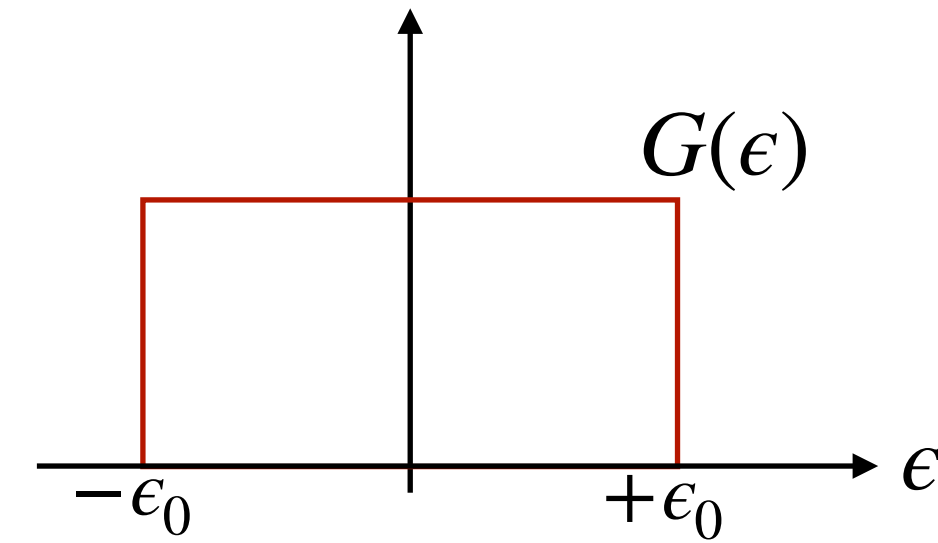
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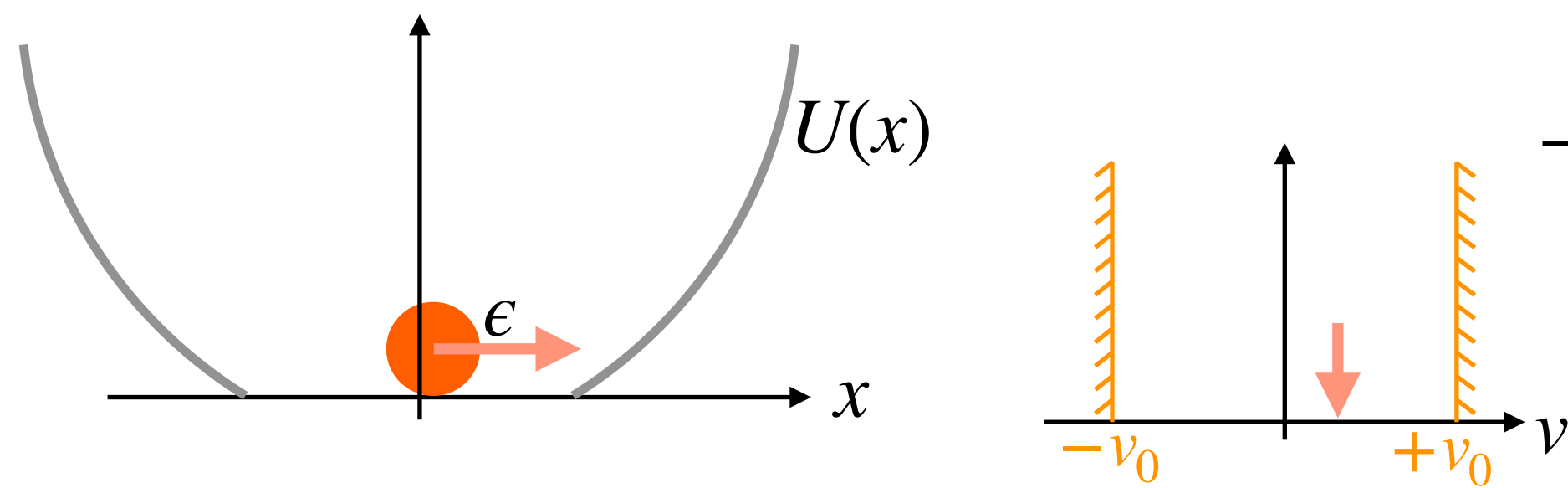
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$$P_{n+1}(x) = \int_{-\infty}^{+\infty} dx' \int_{-\epsilon_0}^{\epsilon_0} d\epsilon \, G(\epsilon) W(x | x', \epsilon) P_n(x')$$





Continuous-time limit

How do we also capture the confining potential?

$$\begin{aligned} v &= \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t} \\ \hline dt &\rightarrow 0 \quad \text{with} \quad v_0 = \text{const} \\ &\quad \tau = \text{const} \end{aligned}$$

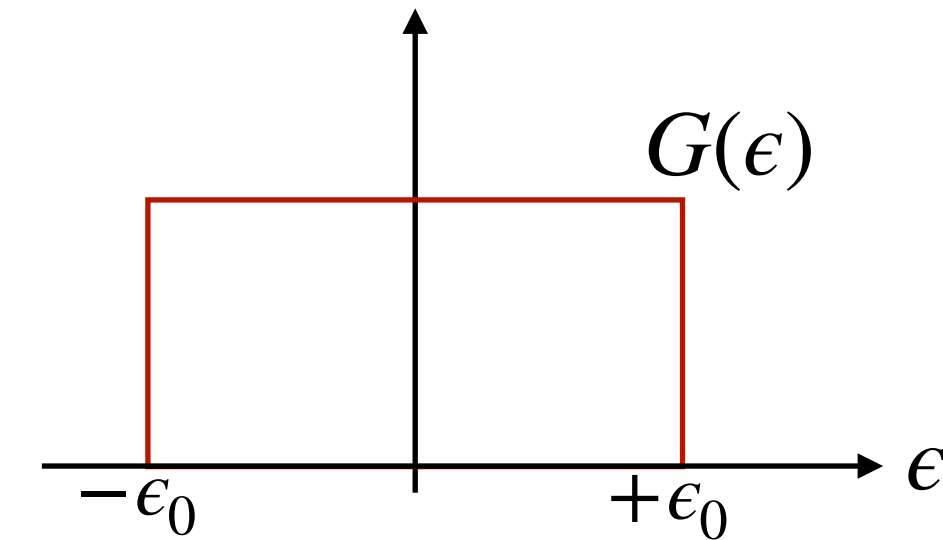
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$$\frac{\partial P_t(x, v)}{\partial t} = D_v \frac{\partial^2}{\partial v^2} P_t(x, v) - \frac{\partial}{\partial x} v P_t(x, v)$$

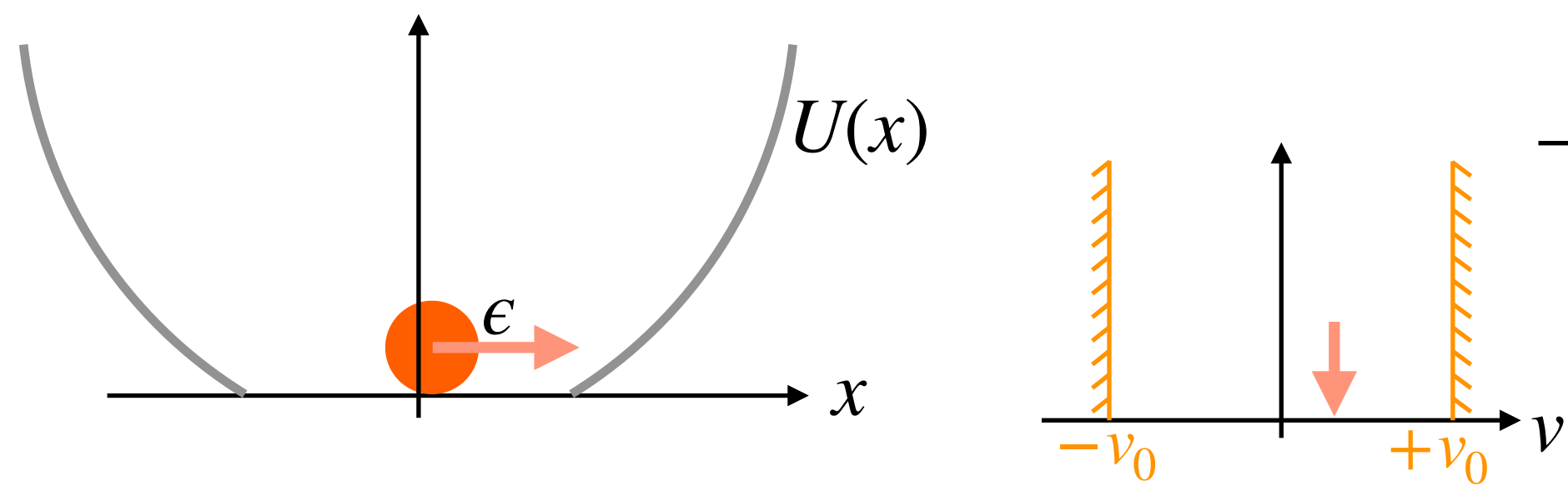
$$\frac{dx}{dt} = v - \cancel{\mu U'(x)} ; \quad \frac{dv}{dt} = \sqrt{2D_v} \xi(t)$$

Correct active dynamics, but NO interactions.



$$\langle \Delta \mathbf{r}(t)^2 \rangle = \sqrt{2D} t$$

$$P_{n+1}(x) = \int_{-\infty}^{+\infty} dx' \int_{-\epsilon_0}^{\epsilon_0} d\epsilon G(\epsilon) W(x | x', \epsilon) P_n(x')$$



Continuous-time limit

How do we also capture the confining potential?

$$v = \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t}$$


---


$$dt \rightarrow 0 \quad \text{with} \quad v_0 = \text{const} \quad \tau = \text{const}$$

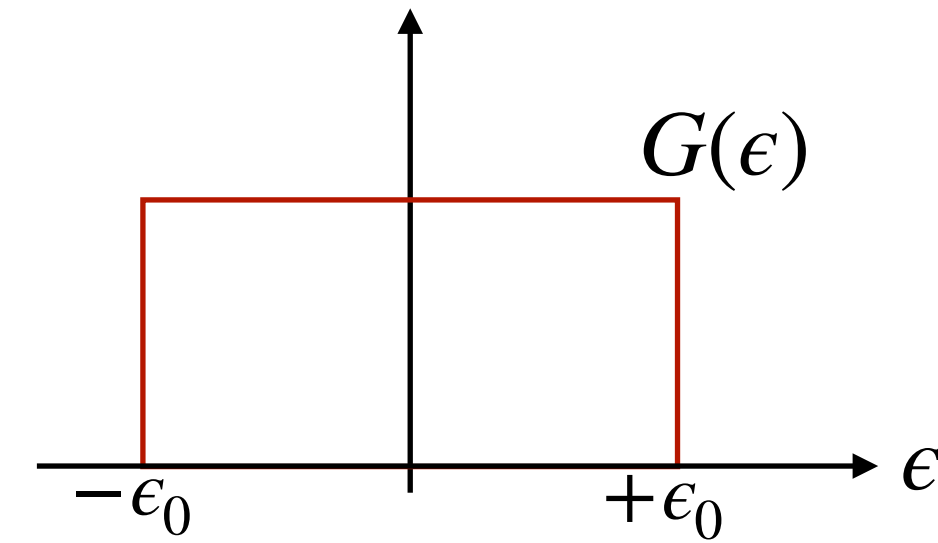
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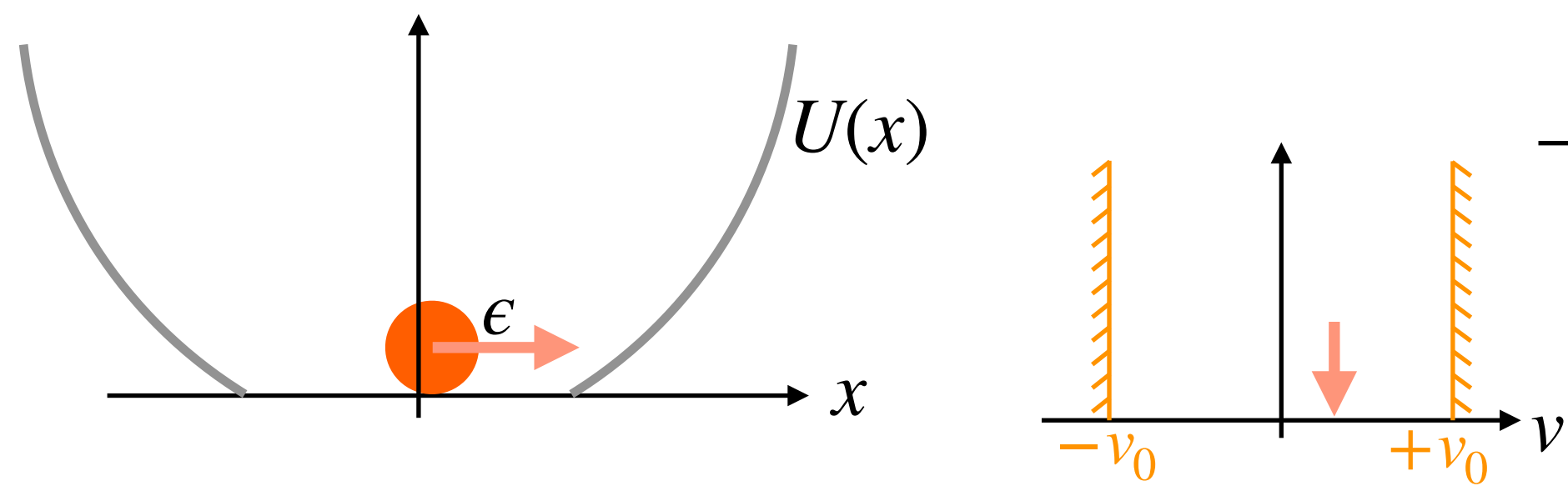
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Kramers-Moyal expansion

$$\frac{\partial P_t(x)}{\partial t} = D \frac{\partial^2}{\partial x^2} P_t(x) - \mu \frac{\partial}{\partial x} U'(x) P_t(x)$$





$$v = \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t}$$


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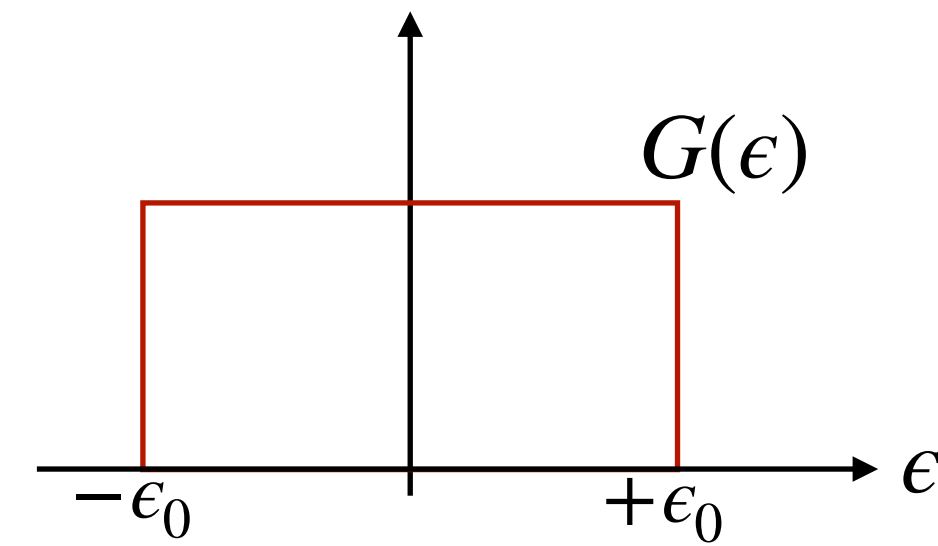
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Correct active dynamics, but NO interactions.

How do we also capture the confining potential?



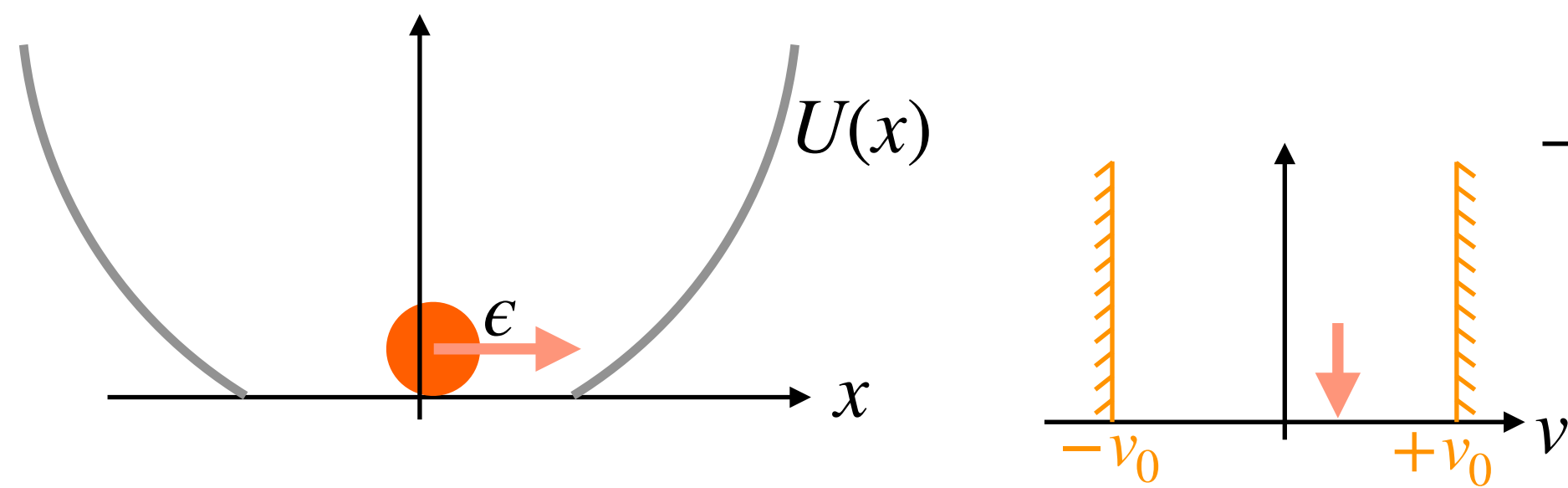
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$$\frac{\partial P_t(x)}{\partial t} = D \frac{\partial^2}{\partial x^2} P_t(x) - \mu \frac{\partial}{\partial x} U'(x) P_t(x)$$

$$\frac{dx}{dt} = -\mu U'(x) + \sqrt{2D} \eta(t), \quad \beta = \mu/D$$



Continuous-time limit

How do we also capture the confining potential?

$$v = \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t}$$


---


$$dt \rightarrow 0 \quad \text{with} \quad v_0 = \text{const} \quad \tau = \text{const}$$

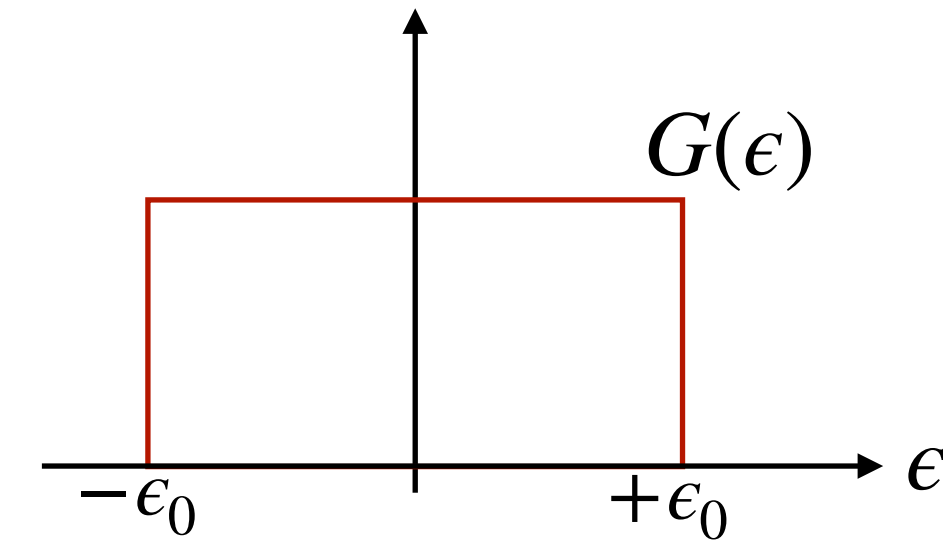
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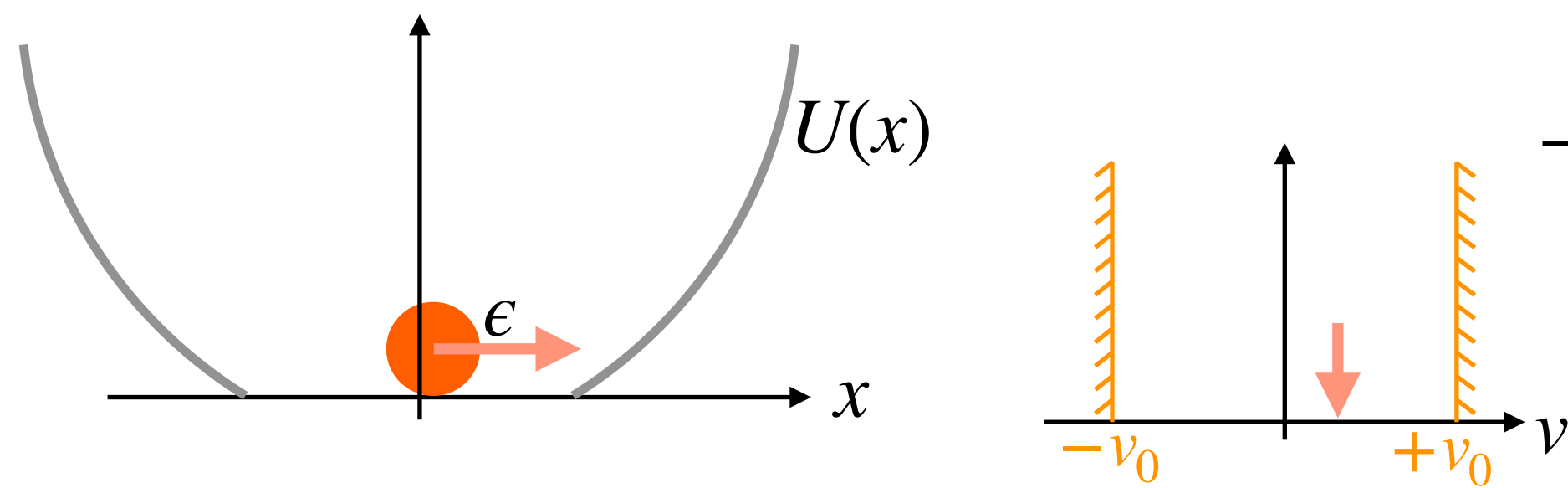
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Interactions (and thermal noise).





$$v = \frac{\epsilon}{dt} \quad t = ndt \quad \langle \Delta v^2(t) \rangle = \sqrt{2D_v t}$$


---


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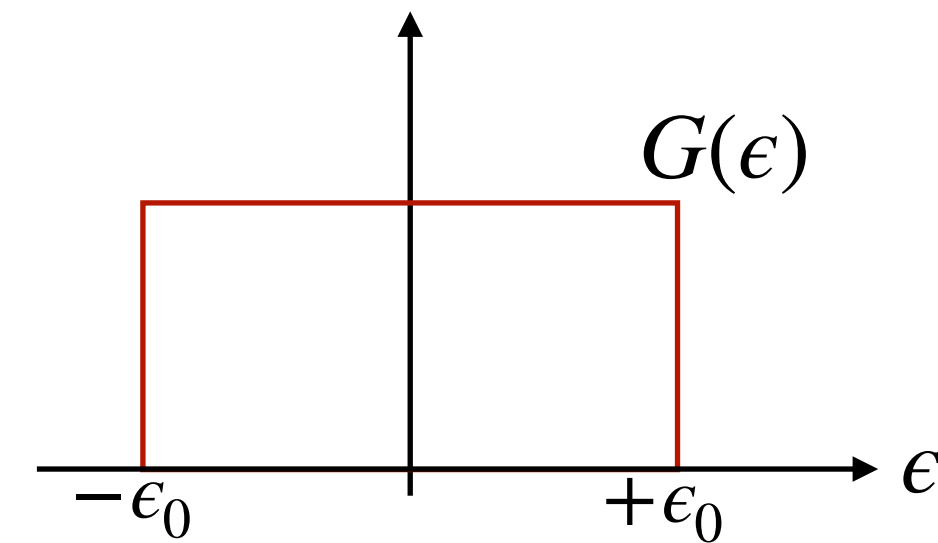
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Correct active dynamics, but NO interactions.

How do we also capture the confining potential?



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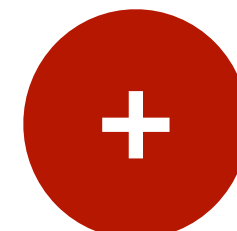
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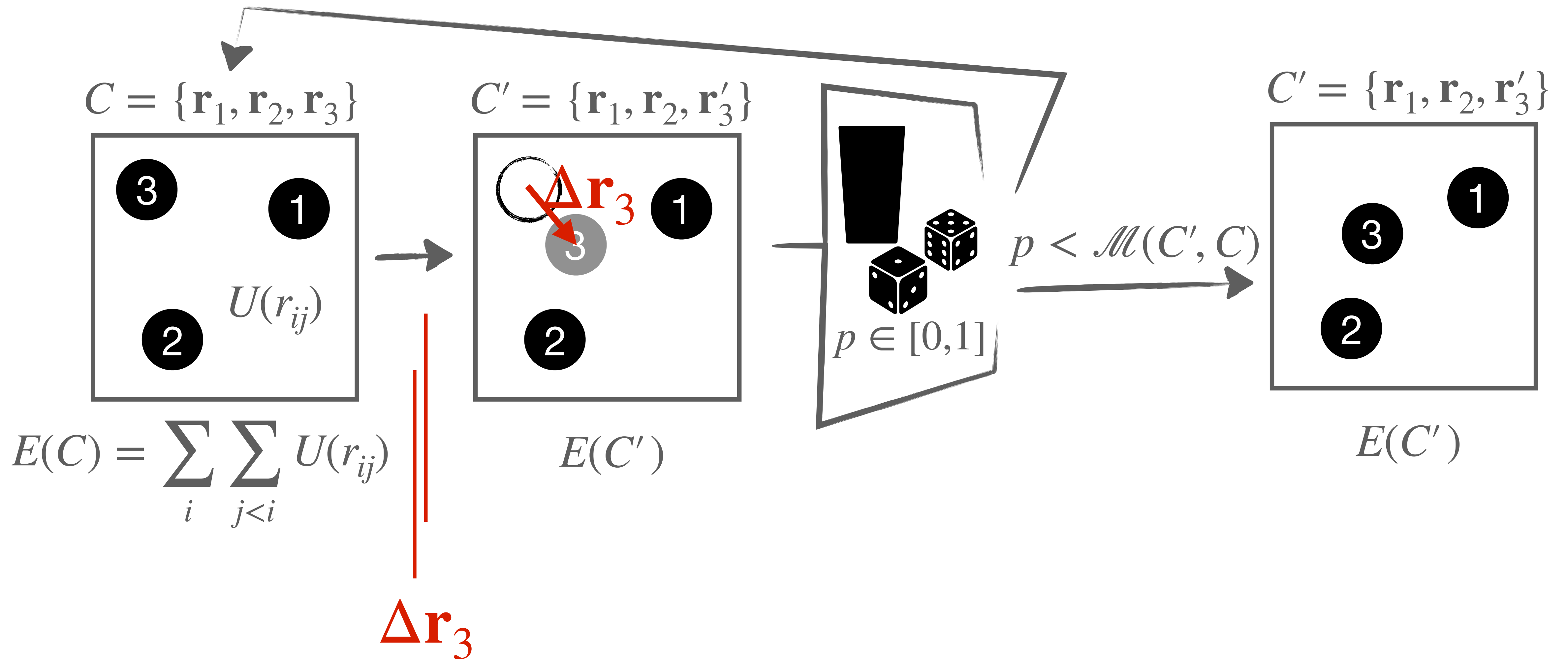
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Interactions (and thermal noise).

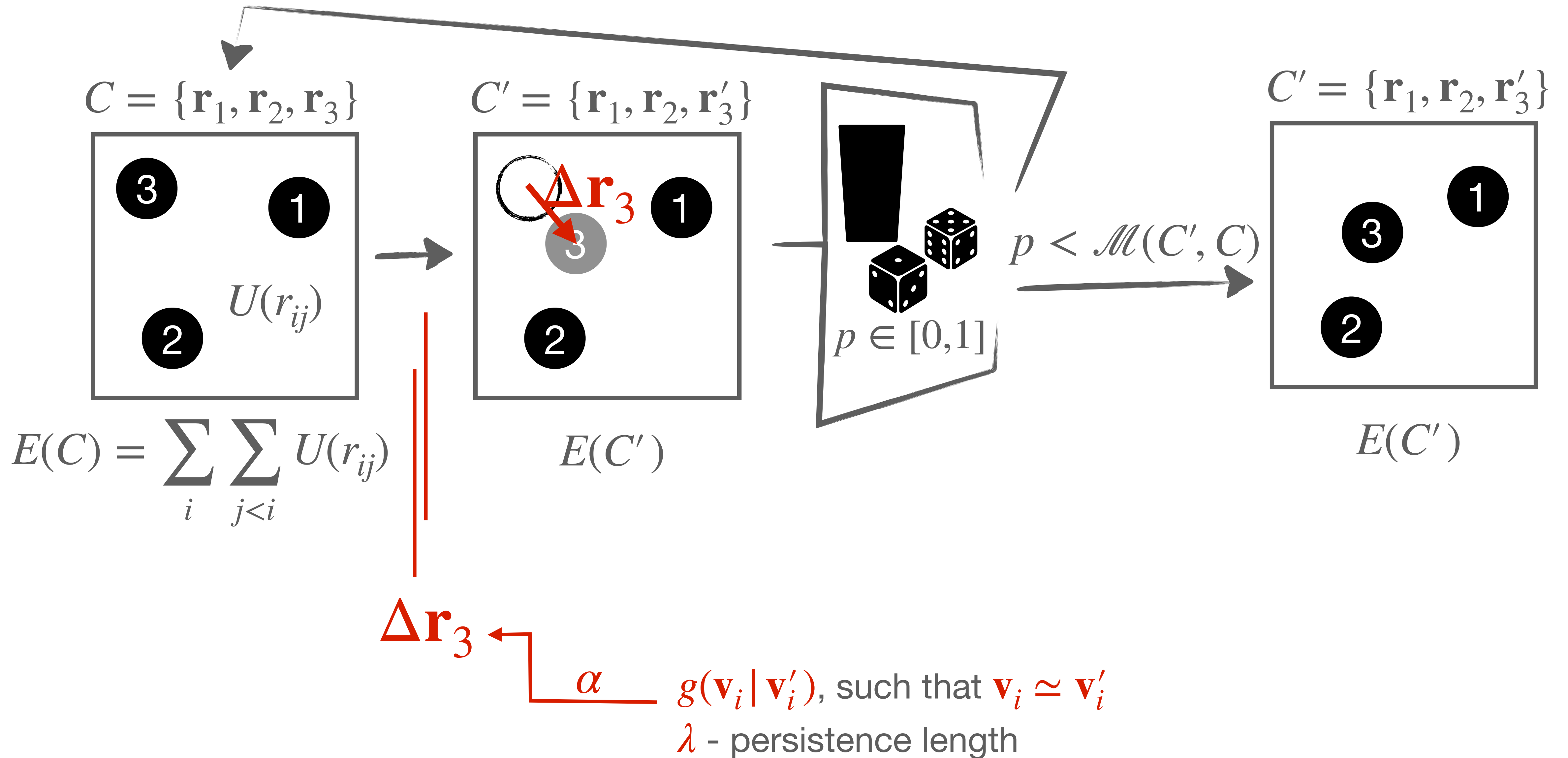


$$\mathcal{M}(C', C) = \min \left\{ 1, e^{-\beta[E(C') - E(C)]} \right\}$$

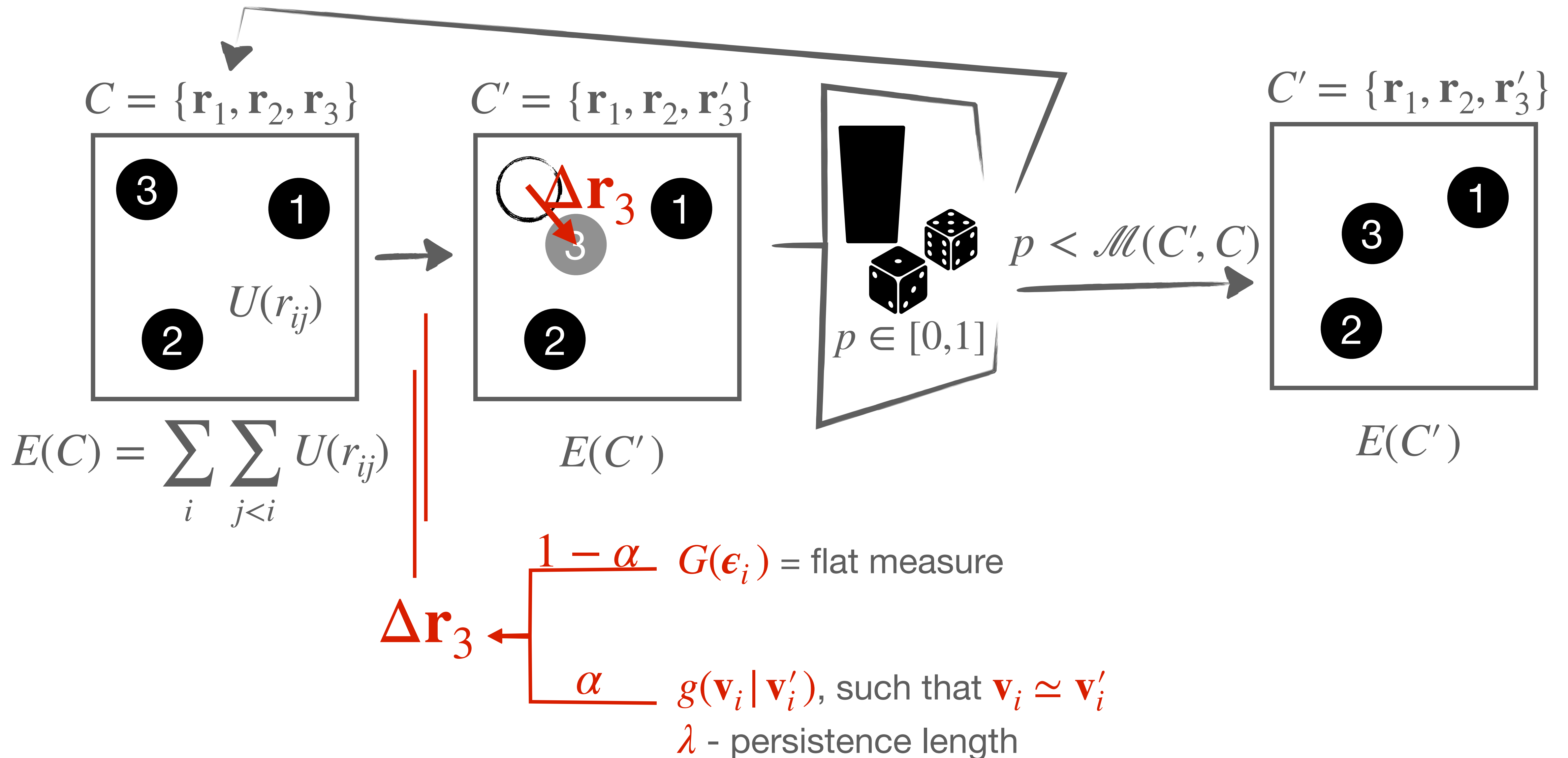




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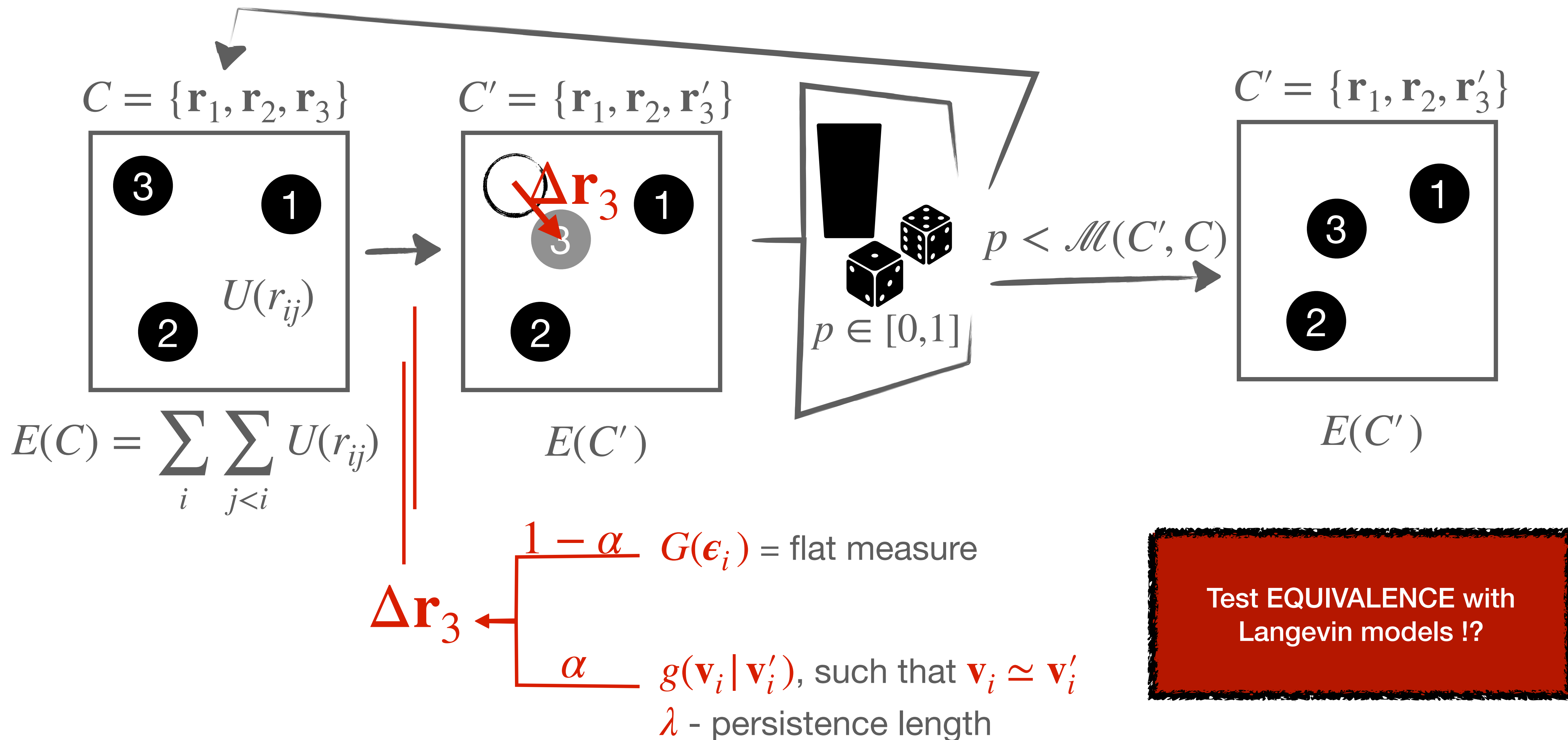


$$\mathcal{M}(C', C) = \min \left\{ 1, e^{-\beta[E(C') - E(C)]} \right\}$$





$$\mathcal{M}(C', C) = \min \left\{ 1, e^{-\beta[E(C') - E(C)]} \right\}$$



## How to test?

---

$$\frac{d\mathbf{r}_i}{dt} = \mathbf{v}_i(t) + \mu \sum_{j \neq i} F(\mathbf{r}_{ij}) + \sqrt{2D} \boldsymbol{\eta}_i(t)$$



## How to test?

---

$$\frac{d\mathbf{r}_i}{dt} = \mathbf{v}_i(t) + \mu \sum_{j \neq i} F(\mathbf{r}_{ij}) + \sqrt{2D} \boldsymbol{\eta}_i(t)$$

### Active Brownian Particles

$$\mathbf{v}_i = v_0 \mathbf{n}(\theta_i); \quad \frac{d\theta_i}{dt} = \sqrt{2D_r} \xi_i$$

### Active Ornstein-Uhlenbeck Process

$$\tau \frac{d\mathbf{v}_i}{dt} = -\mathbf{v}_i + \sqrt{2D_v} \boldsymbol{\xi}_i$$

### Run-And-Tumble Particles

$$\mathbf{v}_i = v_0 \mathbf{n}(\theta_i)$$

$\theta_i$  sampled uniformly at finite rate

### Active Random-Acceleration Process

$$\frac{d\mathbf{v}_i}{dt} = \sqrt{2D_v} \boldsymbol{\xi}_i \quad \text{\& reflecting boundaries at } |\mathbf{v}| = v_0$$

## How to test?

$$\frac{d\mathbf{r}_i}{dt} = \mathbf{v}_i(t) + \mu \sum_{j \neq i} F(\mathbf{r}_{ij}) + \sqrt{2D} \boldsymbol{\eta}_i(t)$$

### Active Brownian Particles

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### Active Random-Acceleration Process

$$\frac{d\mathbf{v}_i}{dt} = \sqrt{2D_v} \boldsymbol{\xi}_i \quad \text{\& reflecting boundaries at } |\mathbf{v}| = v_0$$

Thermodynamic pressure:

$$\mathcal{P}_{ideal} = \rho_{bulk} (k_B T + \text{const})$$

A. P. Solon et al Nature Phys (2015).



# How to test?

$$\frac{d\mathbf{r}_i}{dt} = \mathbf{v}_i(t) + \mu \sum_{j \neq i} \mathbf{F}(\mathbf{r}_{ij}) + \sqrt{2D} \boldsymbol{\eta}_i(t)$$

## Active Brownian Particles

$$\mathbf{v}_i = v_0 \mathbf{n}(\theta_i); \quad \frac{d\theta_i}{dt} = \sqrt{2D_r} \xi_i$$

## Active Ornstein-Uhlenbeck Process

$$\tau \frac{d\mathbf{v}_i}{dt} = -\mathbf{v}_i + \sqrt{2D_v} \xi_i$$

## Run-And-Tumble Particles

$$\mathbf{v}_i = v_0 \mathbf{n}(\theta_i)$$

$\theta_i$  sampled uniformly at finite rate

## Active Random-Acceleration Process

$$\frac{d\mathbf{v}_i}{dt} = \sqrt{2D_v} \xi_i \quad \text{\& reflecting boundaries at } |\mathbf{v}| = v_0$$

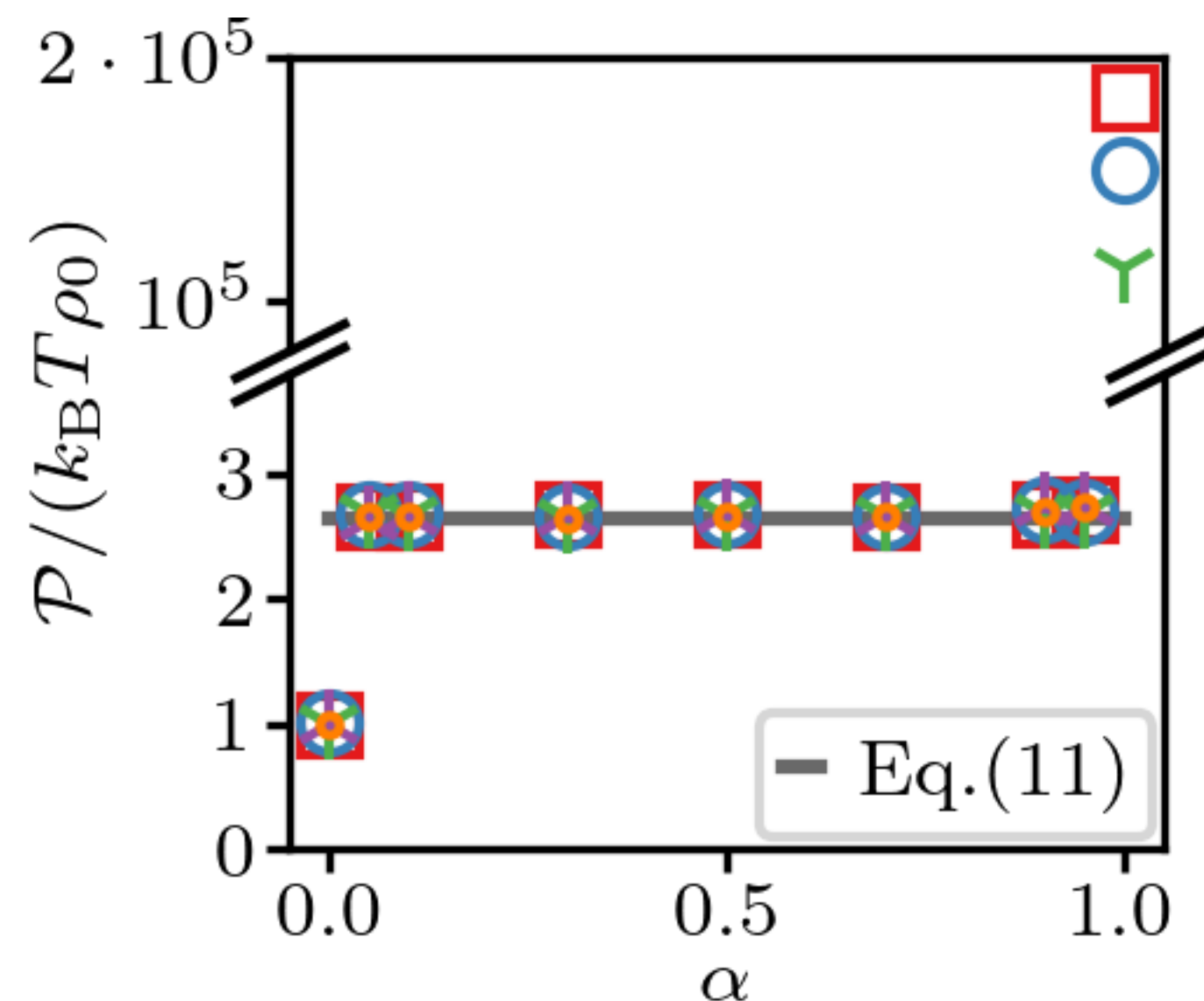
Thermodynamic pressure:

$$\mathcal{P}_{ideal} = \rho_{bulk} (k_B T + \text{const})$$

A. P. Solon et al Nature Phys (2015).

$$P_{n+1}(\mathbf{r}, \mathbf{v}) = \int d\mathbf{r}' \int d\mathbf{v}' M(\mathbf{r}, \mathbf{v} | \mathbf{r}', \mathbf{v}') P_n(\mathbf{r}', \mathbf{v}')$$

$$M(\mathbf{r}, \mathbf{v} | \mathbf{r}', \mathbf{v}') = g(\mathbf{v} | \mathbf{v}') \left[ \alpha W(\mathbf{r} | \mathbf{r}', \mathbf{v}' dt) + (1 - \alpha) \int G(\epsilon) W(\mathbf{r} | \mathbf{r}', \epsilon) \right]$$



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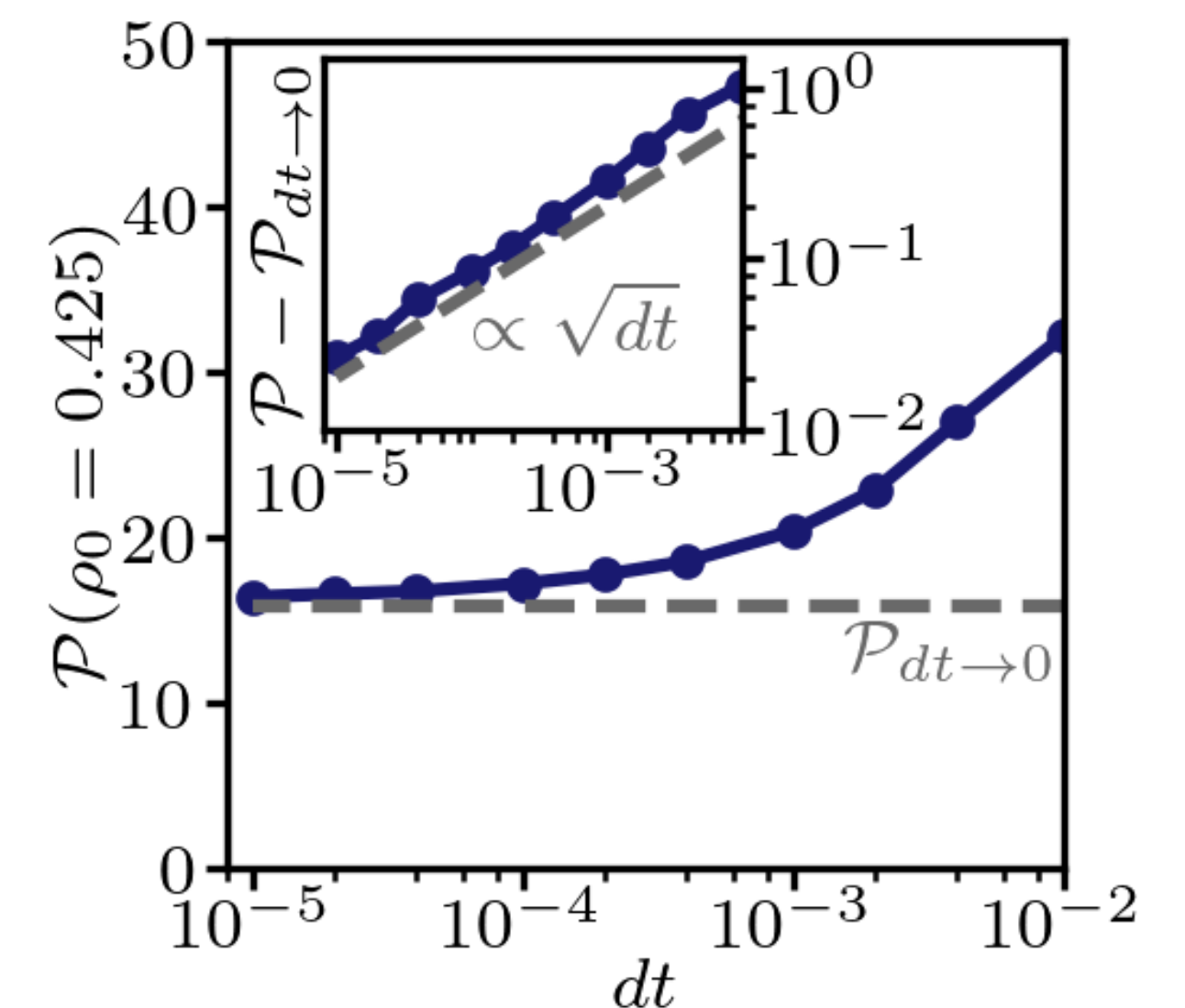
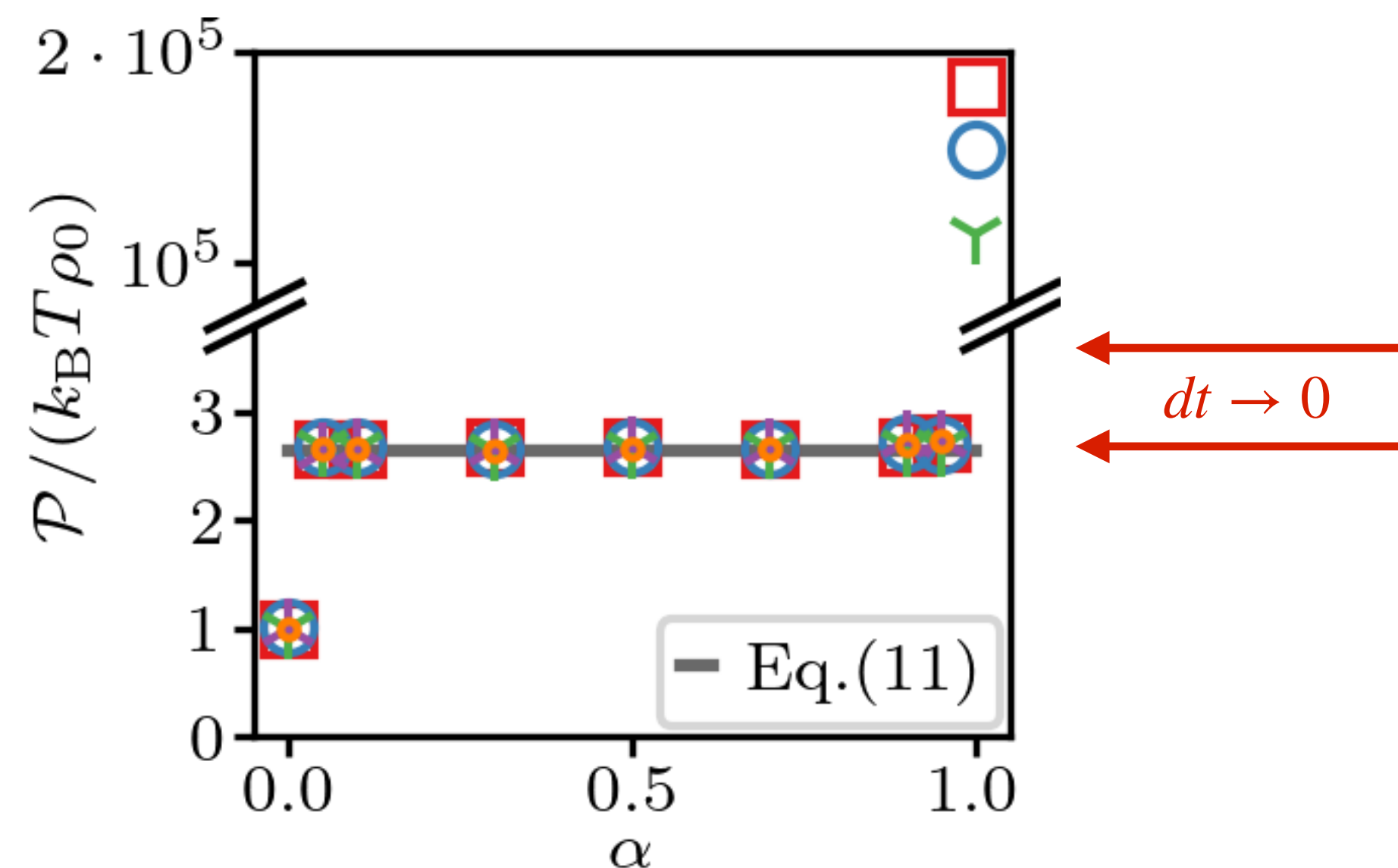
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# Conclusion

## Collective behaviour of persistent particles with active Monte Carlo

- D. Levis, L. Berthier, Phys. Rev. E (2014).
- L. Berthier, Phys Rev Lett (2014).
- **J. Klamser**, S. Kapfer, W. Krauth, Nat Commun (2018).
- **J. Klamser**, S. Kapfer, W. Krauth, J Chem Phys (2019).

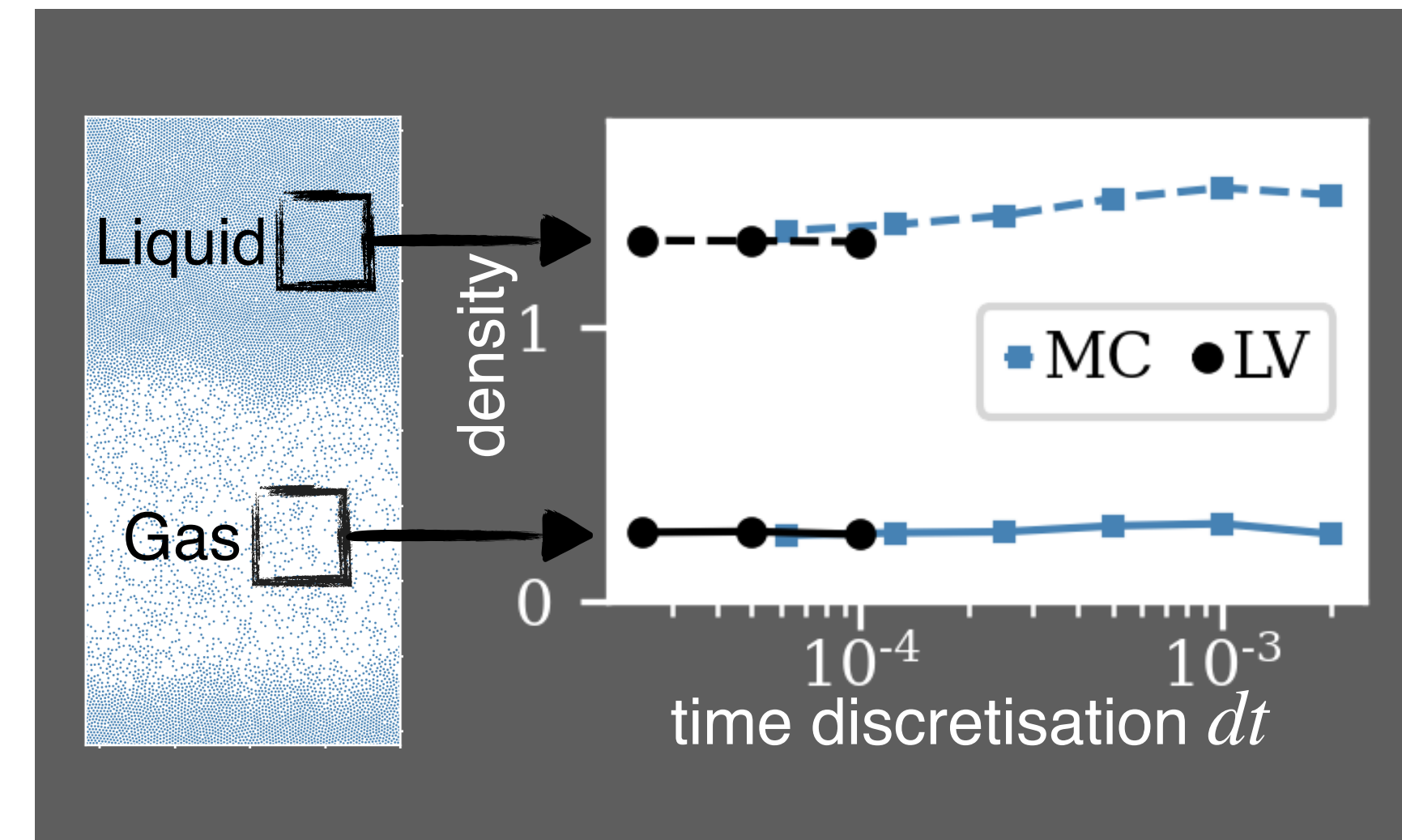
Please, approach anyway!



## Connection to standard models via continuous-time limit

Kinetic Monte Carlo Algorithms for Active Matter Systems

**Juliane U. Klamser**, Olivier Dauchot, Julien Tailleur  
Phys. Rev. Lett. **127** (2021), 150602



## Other MC approaches

Kinetic Event-Chain Algorithm for Active Matter  
T. A. Kampmann, T. Sathiyanesan and J. Kierfeld  
arXiv (2021)